

Quantum Monte Carlo Studies of Mott Metal-Insulator Transitions within Dynamical Mean-field Theory

Nils Blümer

Outline

Motivation: Mott transition in V_2O_3

Introduction: Hubbard Model, DMFT, QMC

Mott transition in frustrated 1-band Hubbard model

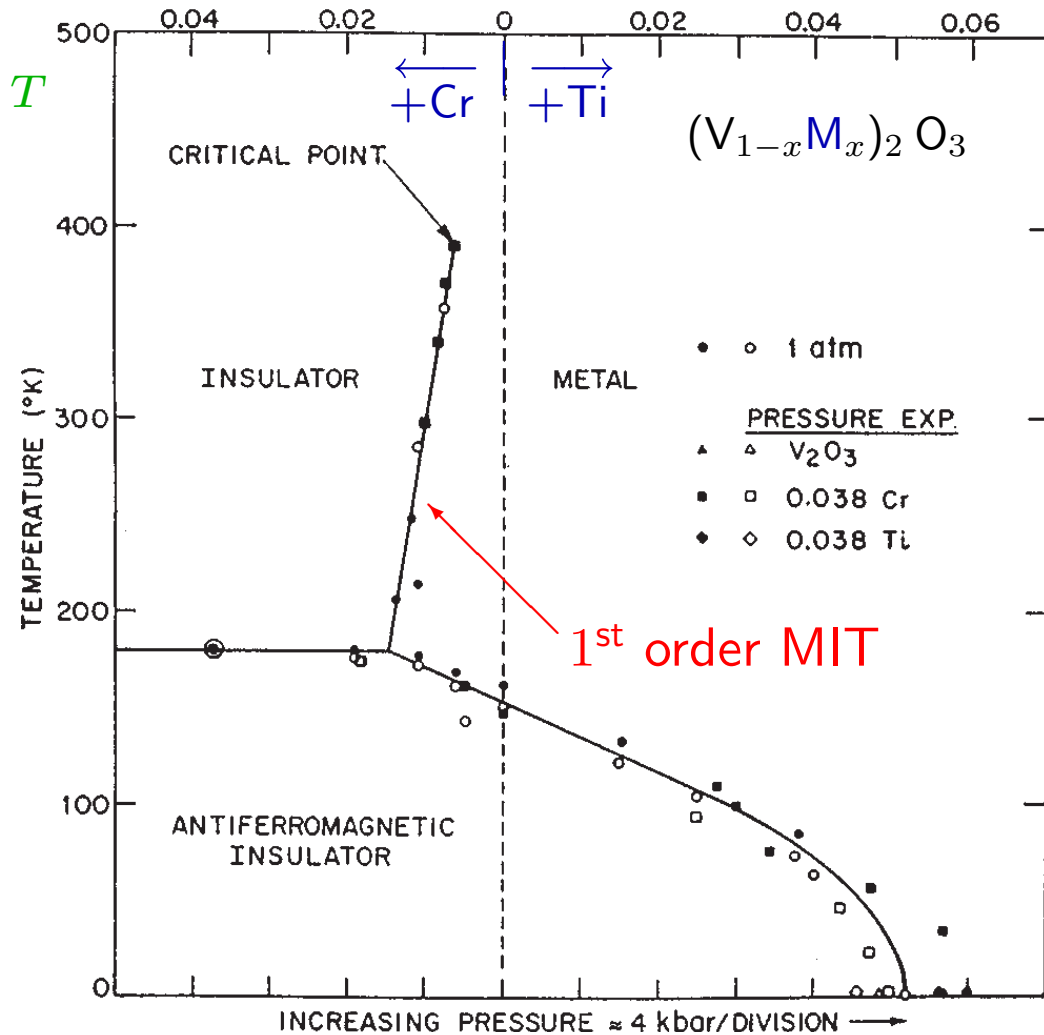
High-precision ground state estimates from QMC

Orbital-selective Mott transition in 2-band Hubbard model

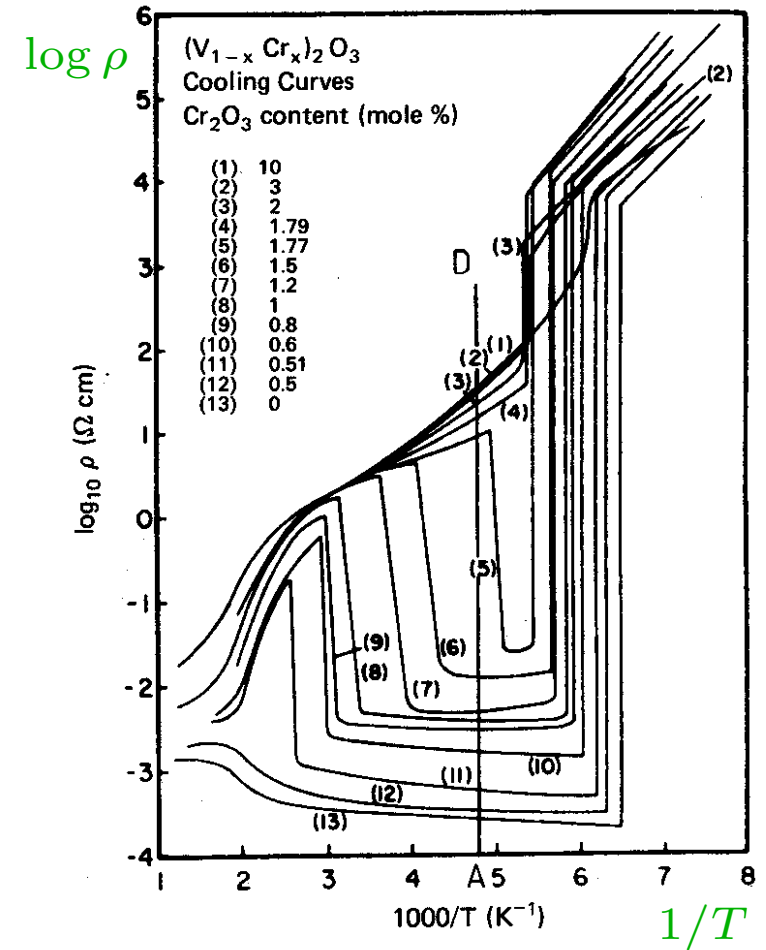
Summary and Outlook

Motivation

Prototype correlated system: V_2O_3

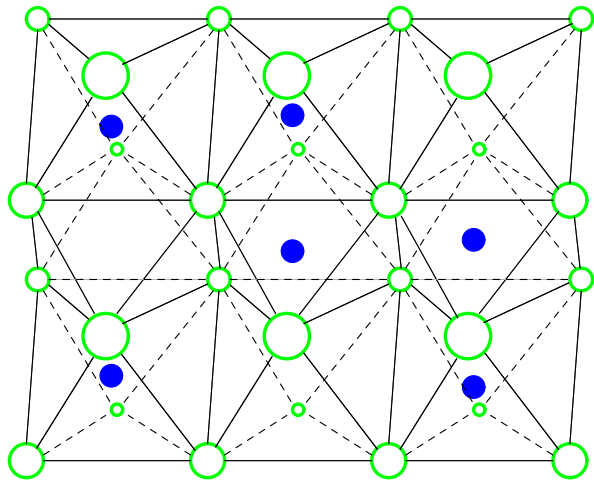


McWhan et al, 1971



Kawamoto et al, 1980

MIT without long-range order
 resistivity ρ increases by factor 10^3
 shift in lattice parameters



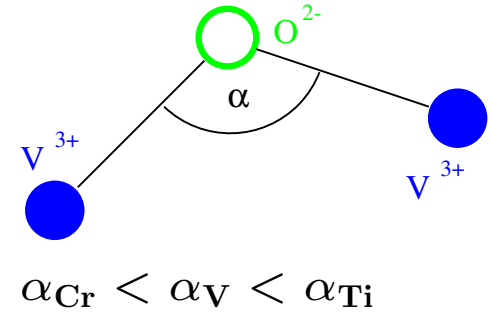
Corundum structure

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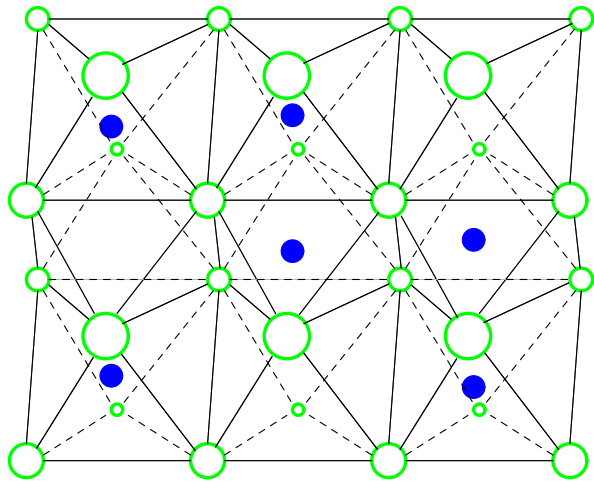
- hcp O^{2-} lattice
- V^{3+} fill 2/3 of octahedra

doping with Ti, Cr:

- (nearly) isovalent
- distorts lattice \rightarrow changes overlap
- drives MIT (like pressure)



Paramagnetic, bandwidth-controlled metal-insulator transition in V_2O_3 \rightarrow microscopic model?



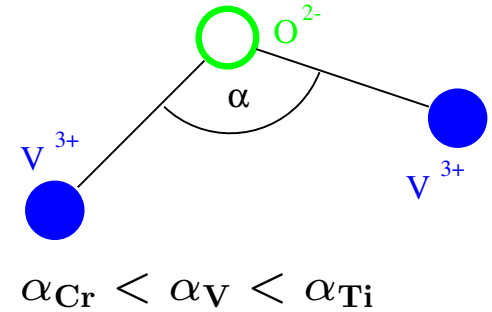
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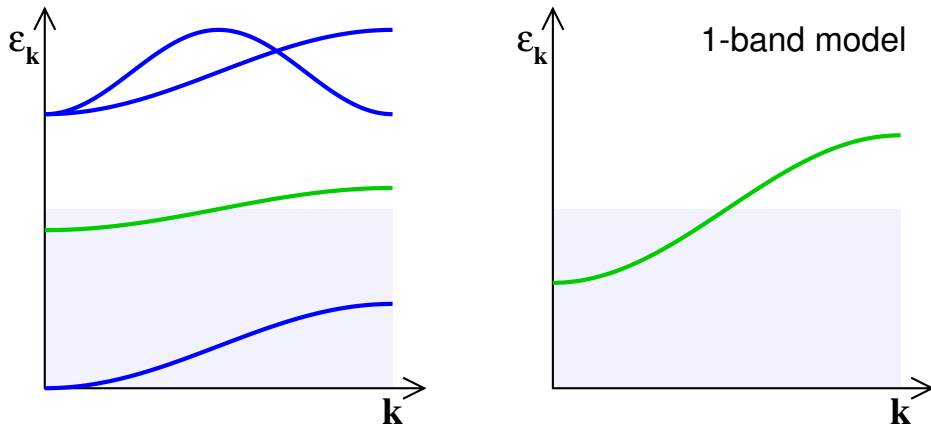
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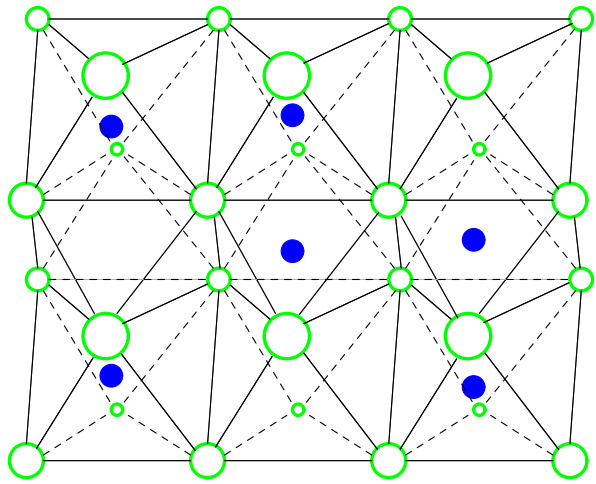
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Bloch states near Fermi energy,





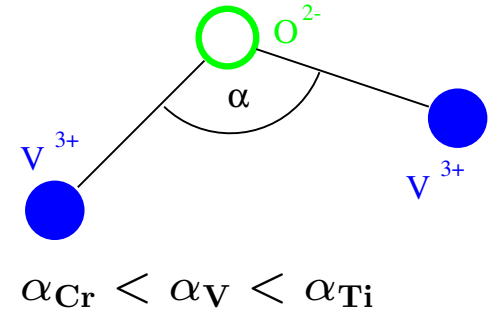
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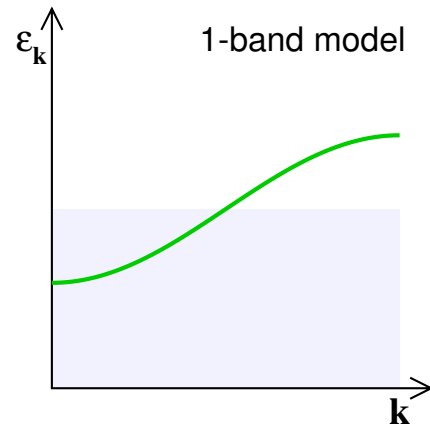
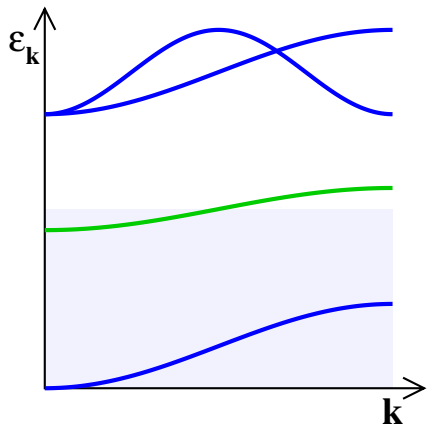
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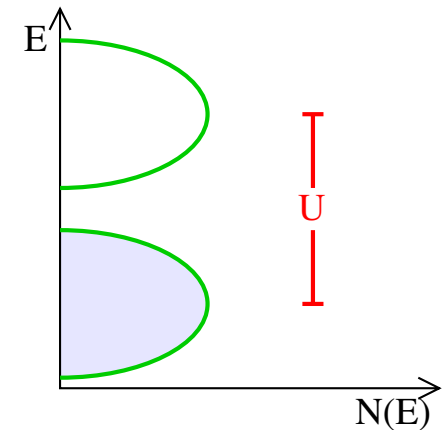
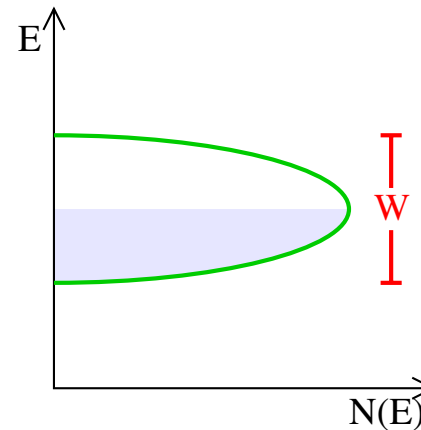


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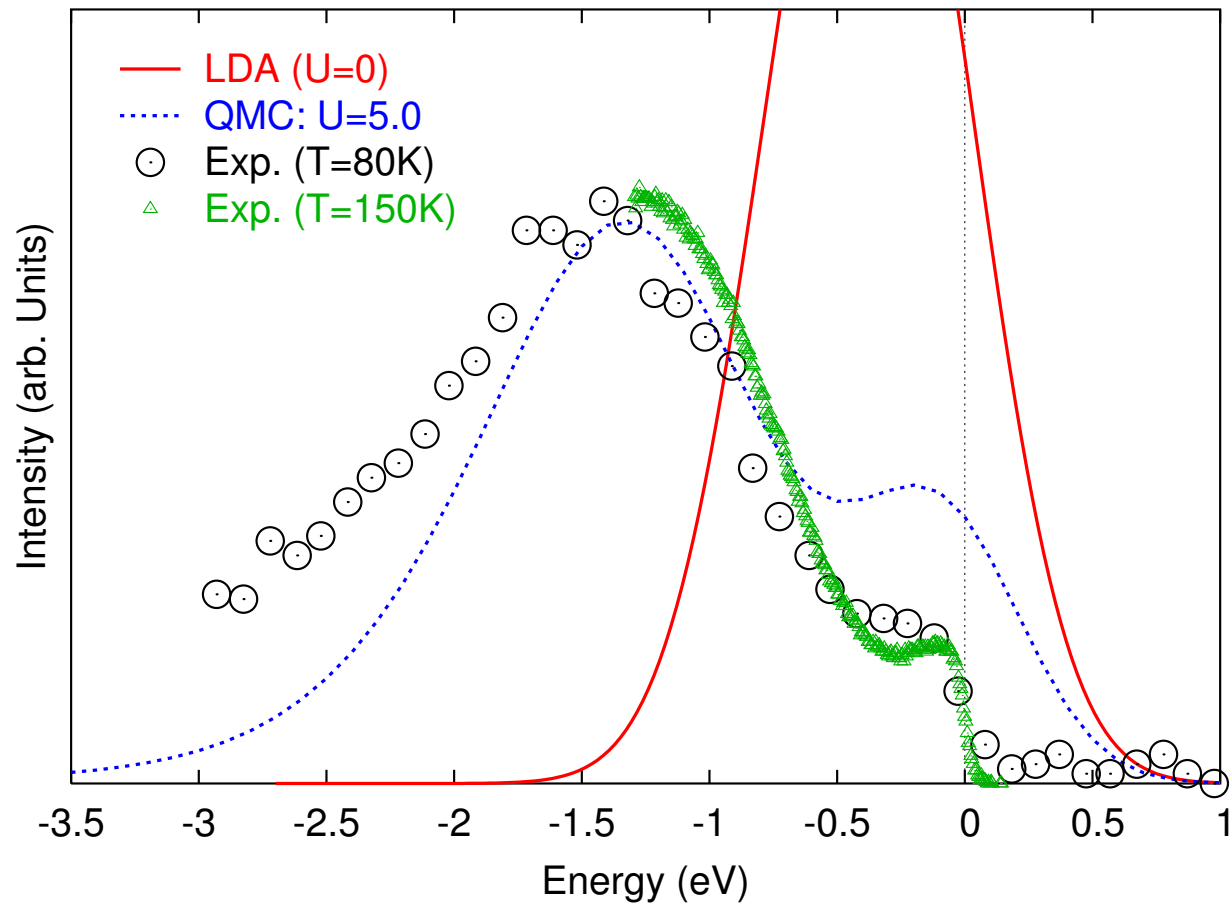


band-splitting by Coulomb correlations ($\sim U$)



Paramagnetic Mott transition not captured by LDA band structure calculations!

System near Mott transition: $\text{La}_{1-x}\text{Sr}_x\text{TiO}_3$ ($x=0.06$) – photoemission spectra



[Nekrasov, Held, NB, Poteryaev, Anisimov, Vollhardt (2000)]

LDA fails to capture low-energy features

LDA+DMFT(QMC): Reasonable accuracy, drastic improvement over LDA

Introduction: Hubbard model, DMFT, QMC

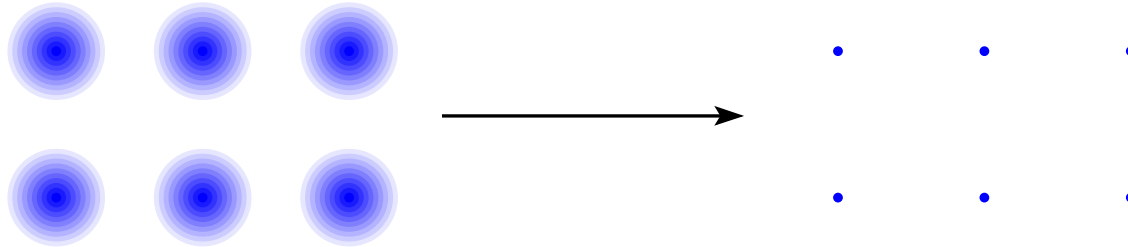
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Born-Oppenheimer approximation ↓

$$H = \sum_{i=1}^{N_e} \frac{\mathbf{p}_i^2}{2m} + \sum_i V(\mathbf{r}_i) + \sum_{i < j} \frac{e^2}{|\mathbf{r}_i - \mathbf{r}_j|}$$



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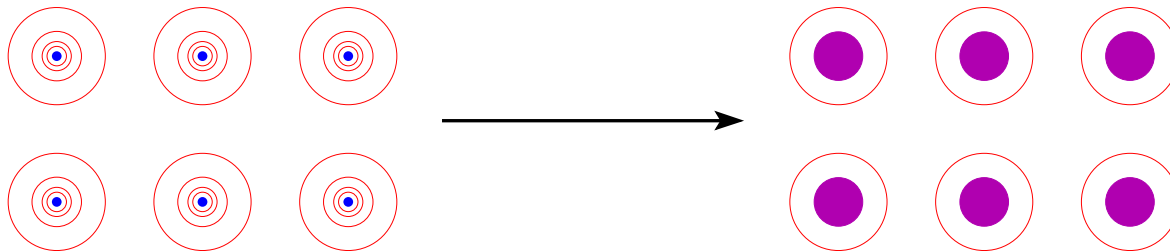
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reduction to valence electrons ↓

$$H = \sum_{i=1}^{N_v} \frac{\mathbf{p}_i^2}{2m} + \sum_{i=1}^{N_v} V^{\text{ion}}(\mathbf{r}_i) + \sum_{i=1}^{N_v-1} \sum_{j=i+1}^{N_v} V^{ee}(\mathbf{r}_i, \mathbf{r}_j)$$



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occupation number formalism, Wannier orbitals ↓

$$\hat{H} = \sum_{i\nu j\sigma} t_{ij}^\nu \hat{c}_{i\nu\sigma}^\dagger \hat{c}_{j\nu\sigma} + \frac{1}{2} \sum_{\nu\nu'\mu\mu'} \sum_{ijmn} \sum_{\sigma\sigma'} v_{ijmn}^{\nu\nu'\mu\mu'} \hat{c}_{i\nu\sigma}^\dagger \hat{c}_{j\nu'\sigma'}^\dagger \hat{c}_{n\mu'\sigma'} \hat{c}_{m\mu\sigma}$$

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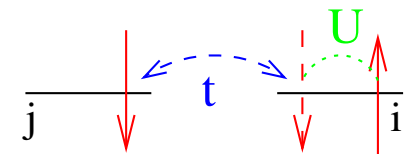
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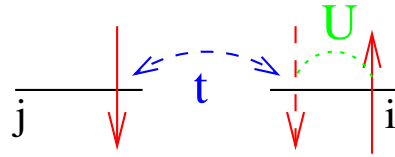
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Hubbard model

$$\hat{H} = \sum_{(i,j),\sigma} t_{ij} (\hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \text{h.c.}) + U \sum_i \hat{n}_{i\uparrow} \hat{n}_{i\downarrow}$$



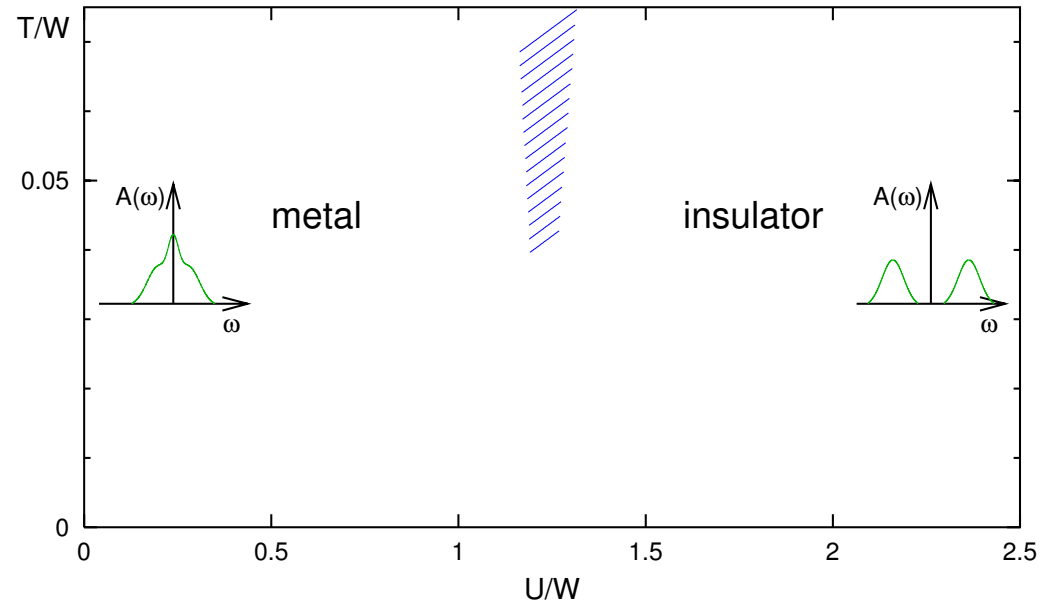
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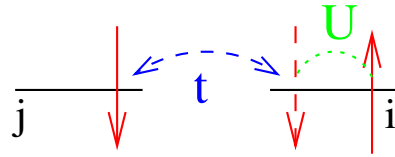
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Minimal model for correlated electrons

MIT/crossover at $U/W \approx 1$ (and half filling)



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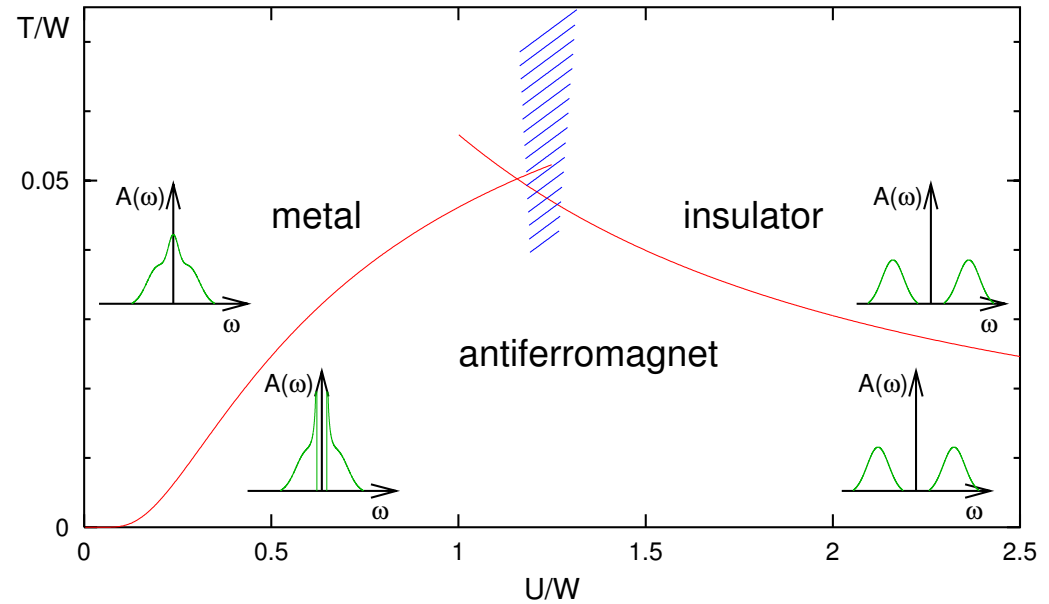


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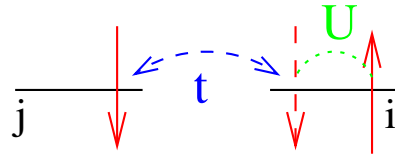
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But: generically antiferromagnetism at low T



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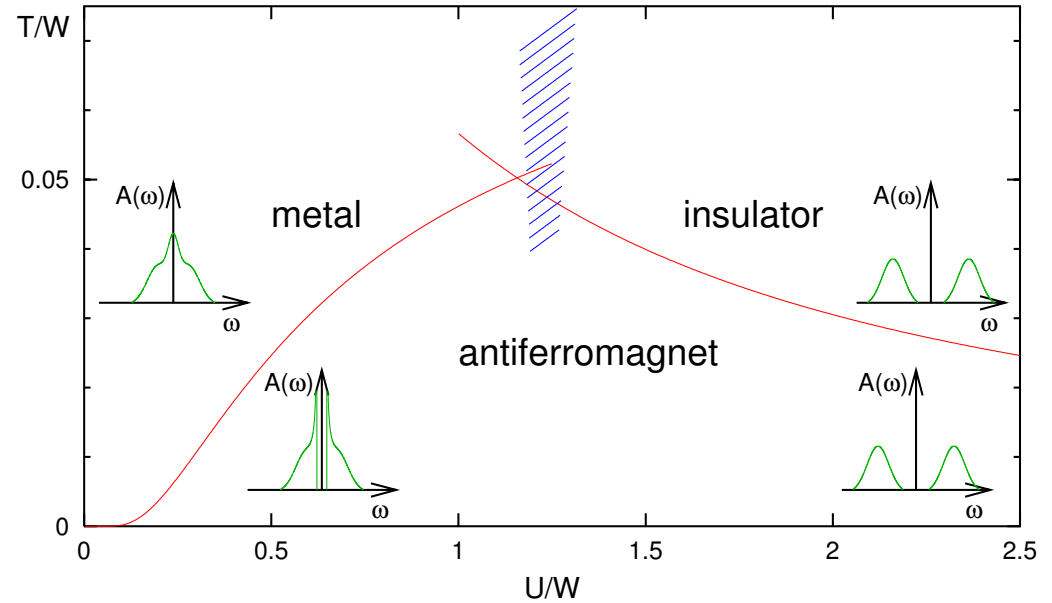


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How to solve (frustrated) Hubbard model in regime of $W \sim U$?

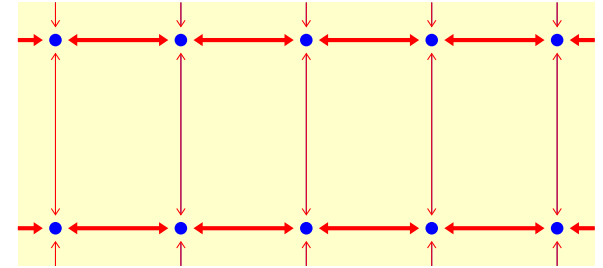
Approaches for Hubbard-type models

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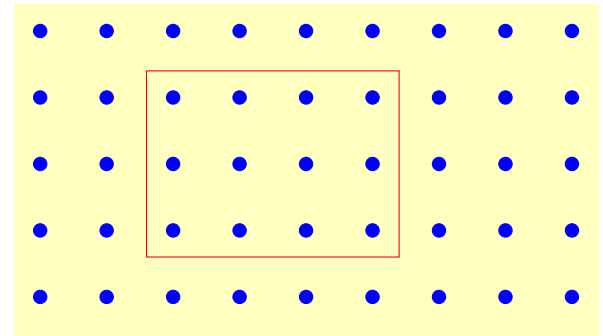
Perturbation theory, e.g.

- $U \rightarrow 0$: Hartree-Fock (**uncorrelated**)
- $t/U \rightarrow 0$: half filling ($n = 1$) \rightsquigarrow Heisenberg model
- $T \rightarrow \infty, n \rightarrow 0$
- ($V_{\text{ion}} \rightarrow 0 \rightsquigarrow$ jellium model \rightsquigarrow LDA)

$d = 1$: Bethe ansatz, DMRG



finite clusters: ED, QMC



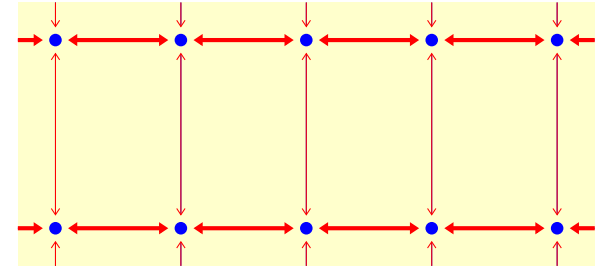
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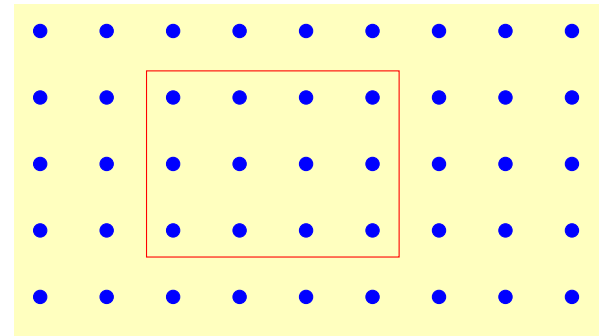
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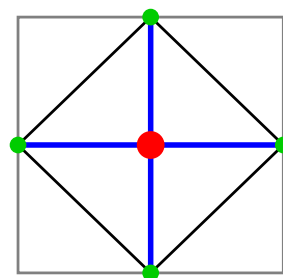


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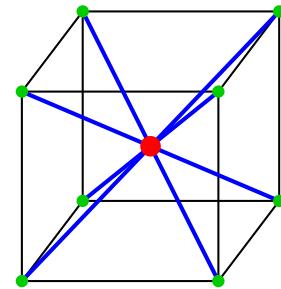


Dynamical mean-field theory (DMFT): local self-energy $\Sigma(\mathbf{k}, \omega) \equiv \Sigma(\omega)$

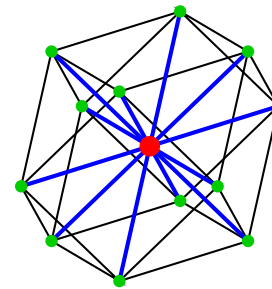
- + non-perturbative \rightsquigarrow valid at MIT
- + dynamical on-site correlations preserved
- + in thermodynamic limit
- + exact for $Z \rightarrow \infty$



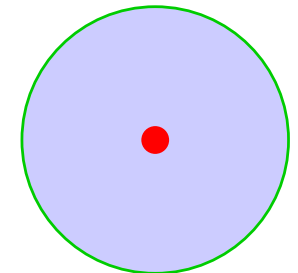
$d=2: Z = 4$



bcc: $Z = 8$



fcc: $Z = 12$



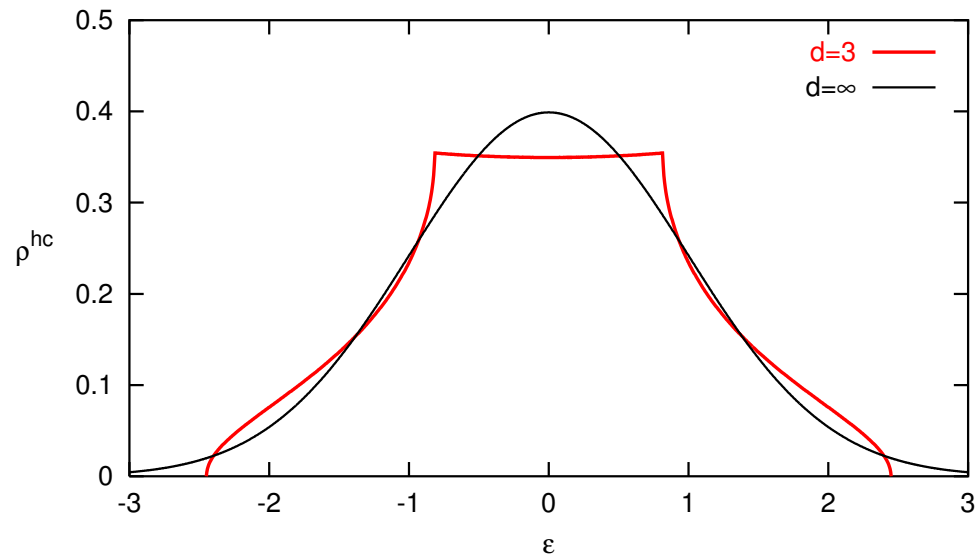
DMFT: $Z = \infty$

Brief History of DMFT

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Müller-Hartmann: local self-energy: $\Sigma(\mathbf{q}, \omega) \rightarrow \Sigma(\omega)$

Brandt, Mielsch: exact solution of Falicov-Kimball model



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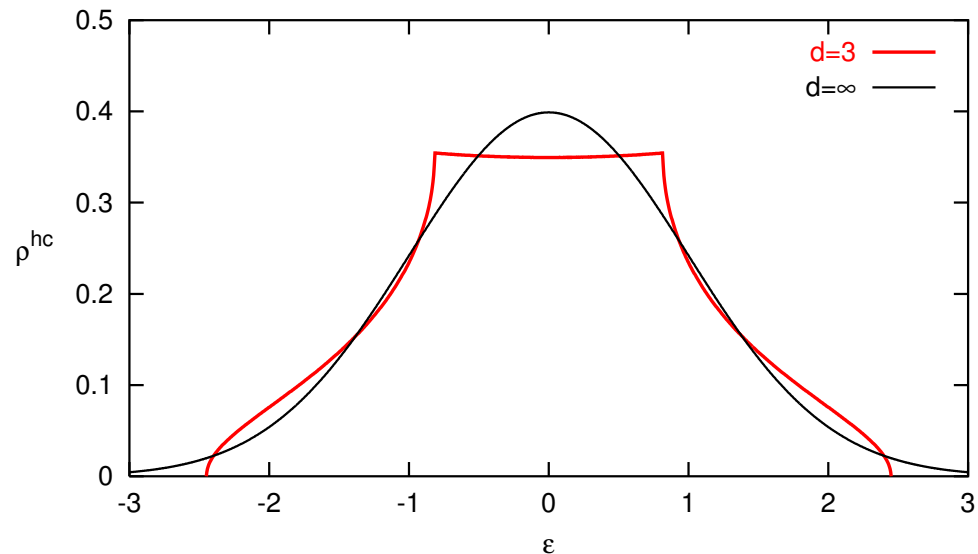
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Jarrell: Quantum Monte-Carlo \rightsquigarrow antiferromagnetism, Mott-Hubbard behavior



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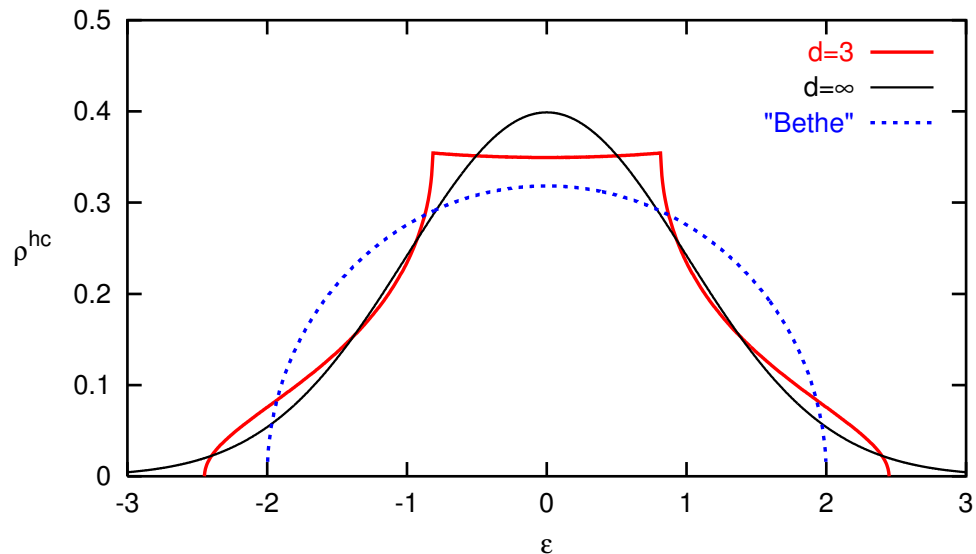
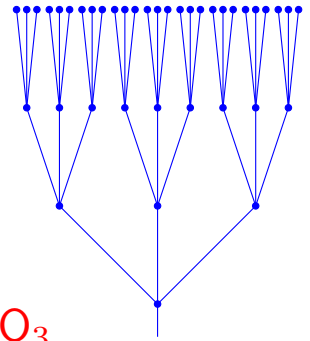
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Vollhardt group: ferro/metamagnetism (NN exchange, asymmetric DOS, multi-band)



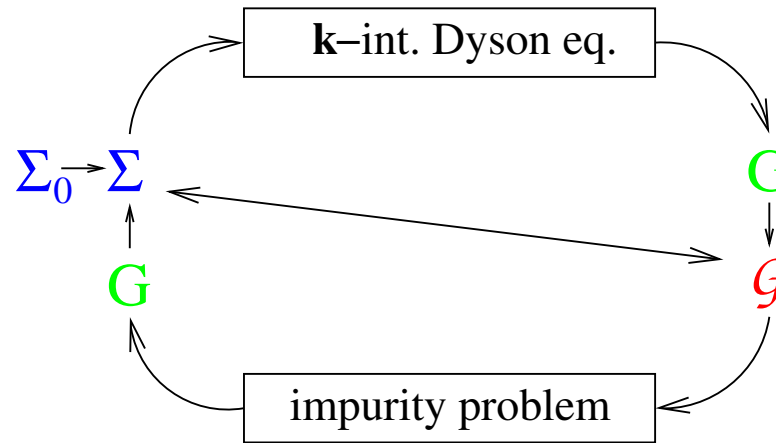
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Jarrell group: dynamical cluster approximation (DCA) \rightsquigarrow d -wave superconductivity

Iterative solution of DMFT equations

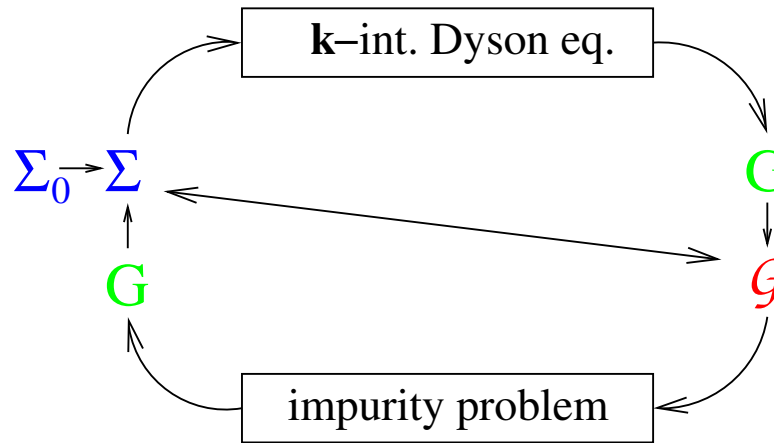


← DOS / t_{ij} / $\epsilon_{\nu k}$

Fourier transformations

← local interactions

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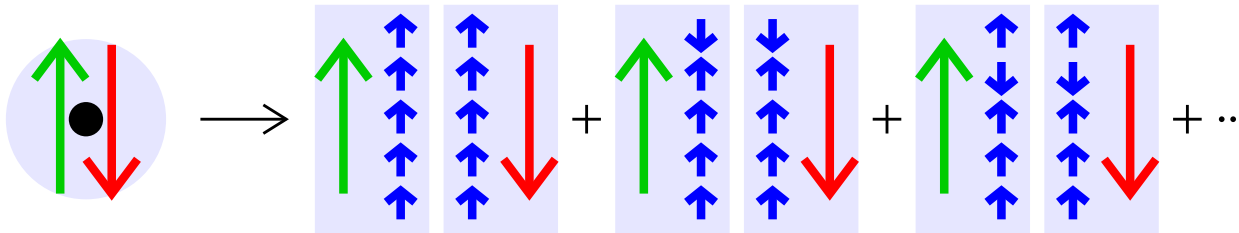


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QMC: discretization $\beta = \Lambda \Delta\tau$, Trotter decoupling, discrete Hubbard-Stratonovich transformation



Wick theorem:

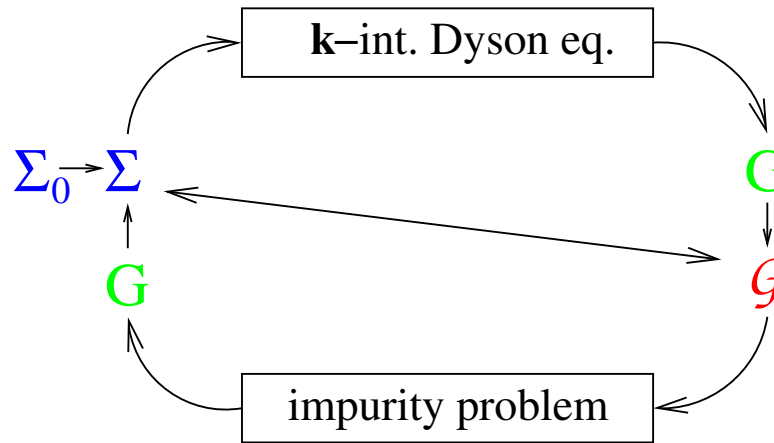
$$G = \frac{\sum M \det\{M\}}{\sum \det\{M\}}$$

Metropolis MC importance sampling over auxiliary Ising field, 2^Λ configurations, $50 \lesssim \Lambda \lesssim 400$

+ numerically exact

– effort scales as T^{-3}

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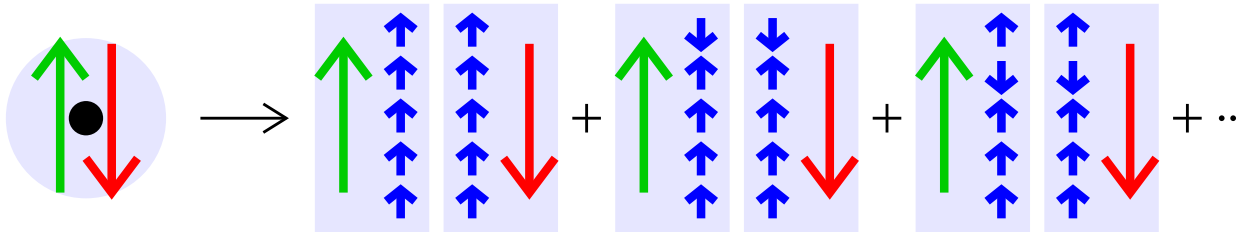


← DOS / t_{ij} / $\epsilon_{\nu\mathbf{k}}$

Fourier transformations

← local interactions

QMC: discretization $\beta = \Lambda \Delta\tau$, Trotter decoupling, discrete Hubbard-Stratonovich transformation

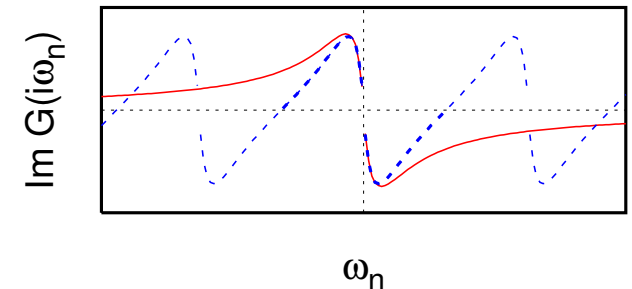
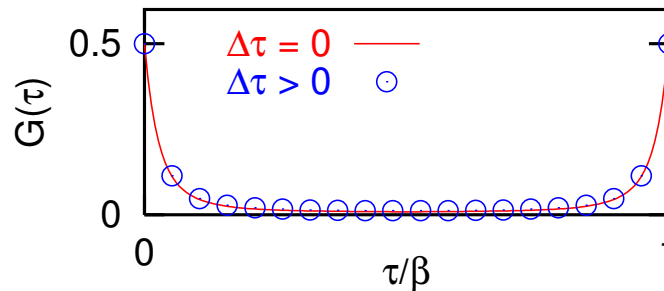


Wick theorem:

$$G = \frac{\sum M \det\{M\}}{\sum \det\{M\}}$$

Metropolis MC importance sampling over auxiliary Ising field, 2^Λ configurations, $50 \lesssim \Lambda \lesssim 400$

- + numerically exact
- effort scales as T^{-3}
- no info for $\omega \gtrsim \omega_{\text{Nyquist}}$



Mott transition in frustrated 1-band Hubbard model

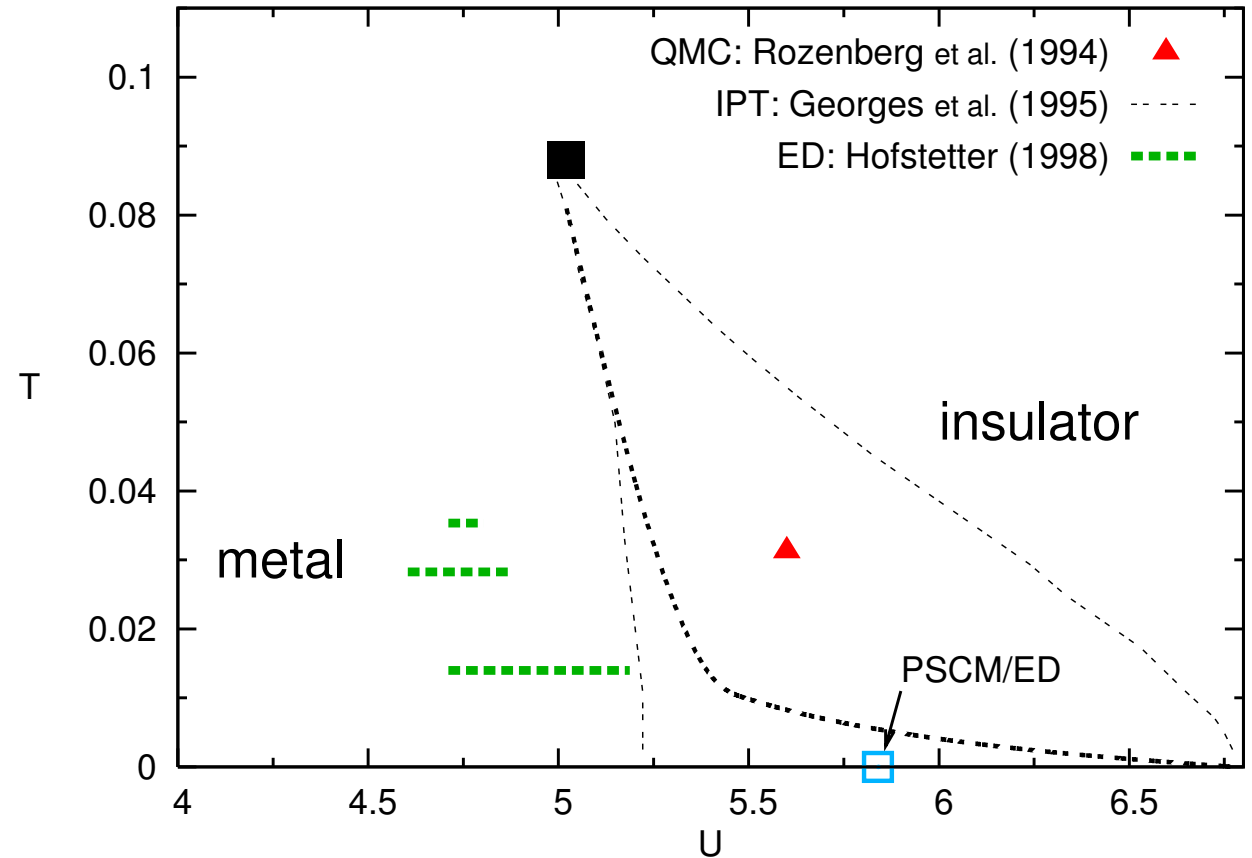
Controversy in 1999:

1st order MIT (“Bethe” DOS)?

Hysteresis in DMFT cycle?

Coexistence of metal + insulator?

- IPT, ED: **yes!**



Mott transition in frustrated 1-band Hubbard model

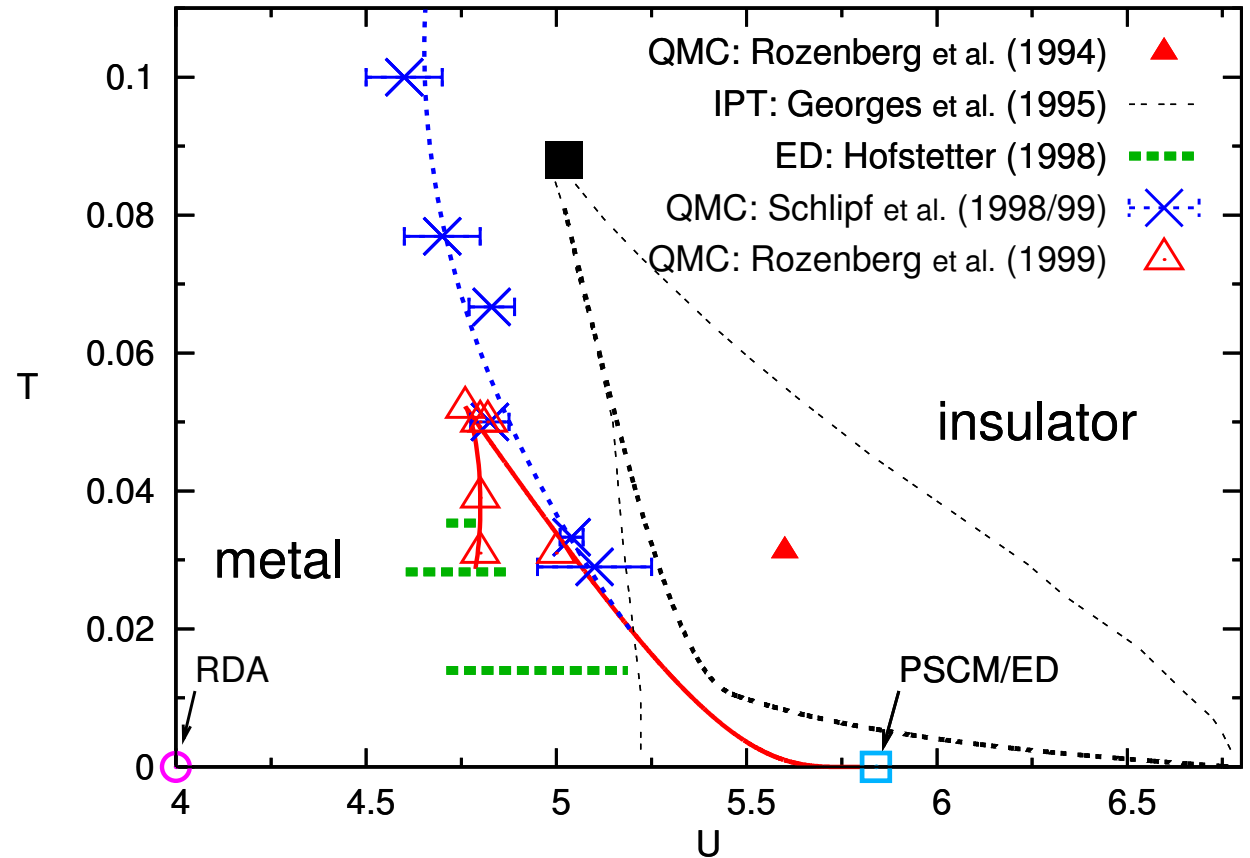
Controversy in 1999:

1st order MIT (“Bethe” DOS)?

Hysteresis in DMFT cycle?

Coexistence of metal + insulator?

- IPT, ED: **yes!**
- QMC (Schlipf et al.): **no!**
- RDA: **no!** (much lower U_c)
- QMC (Rozenberg et al.): **yes!**



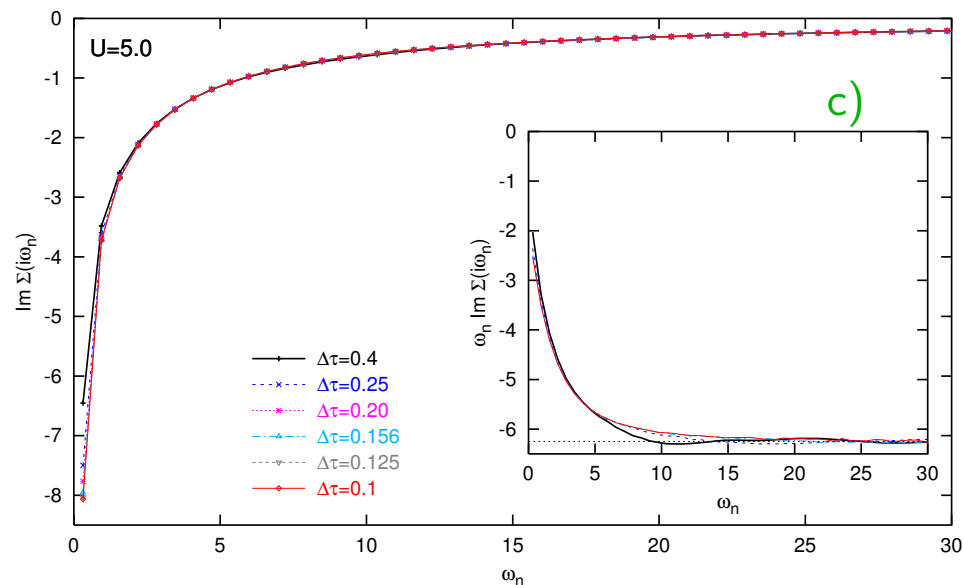
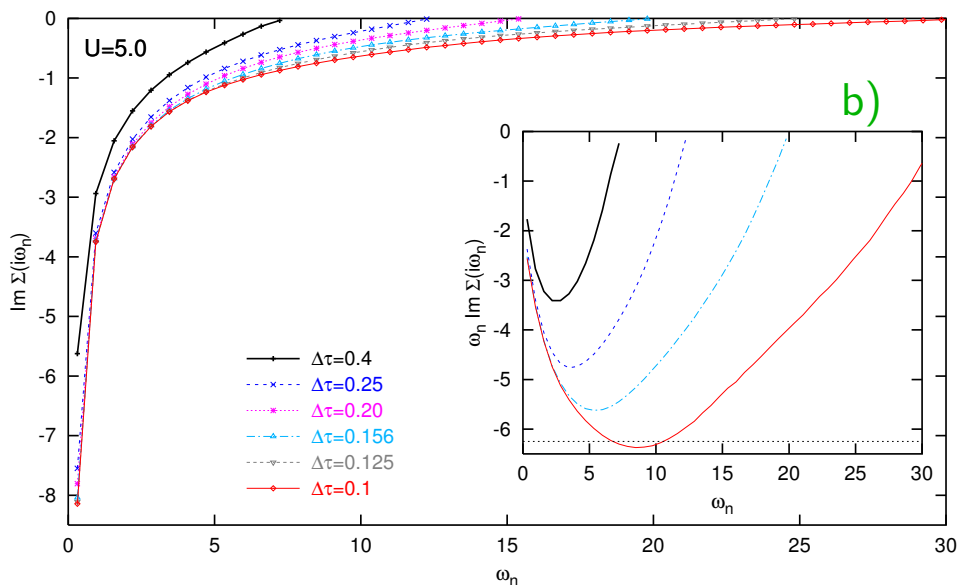
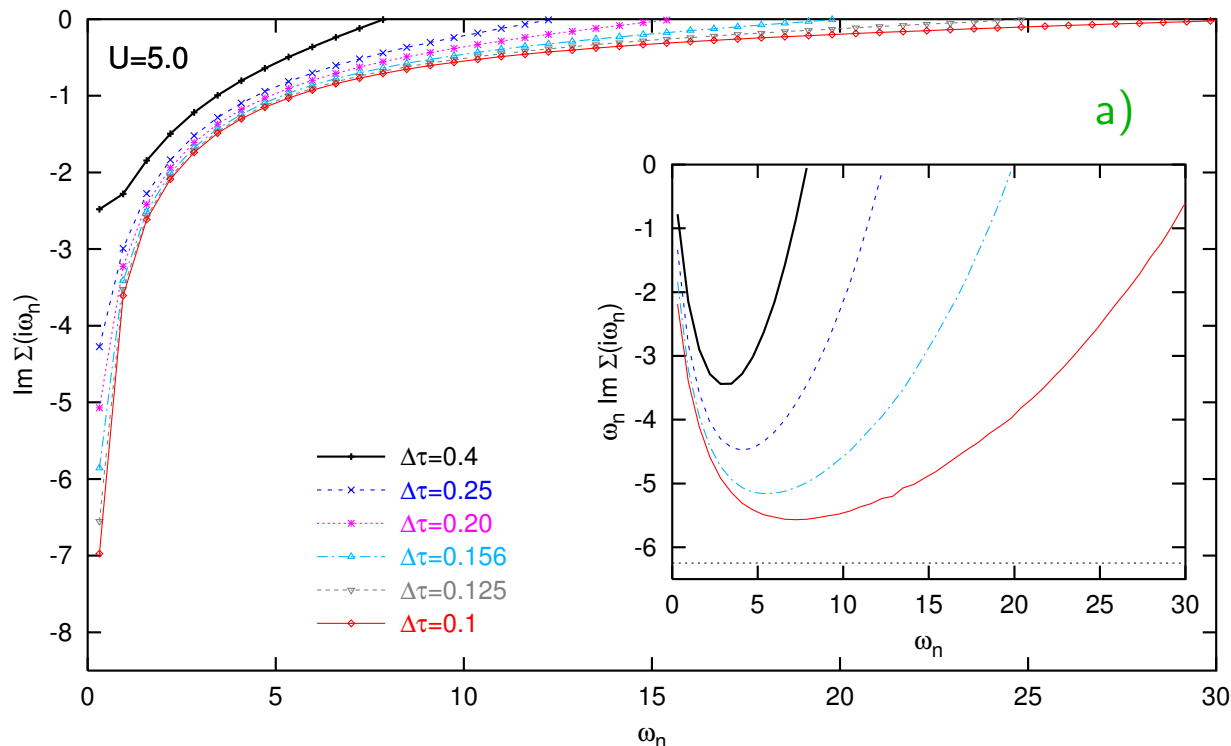
- Who is right / What went wrong?
- Precise coexistence phase diagram?
- Thermodynamic first order phase transition line?

Fourier transformation schemes:
 self-energy Σ ($T = 0.1, U = 5.0$)

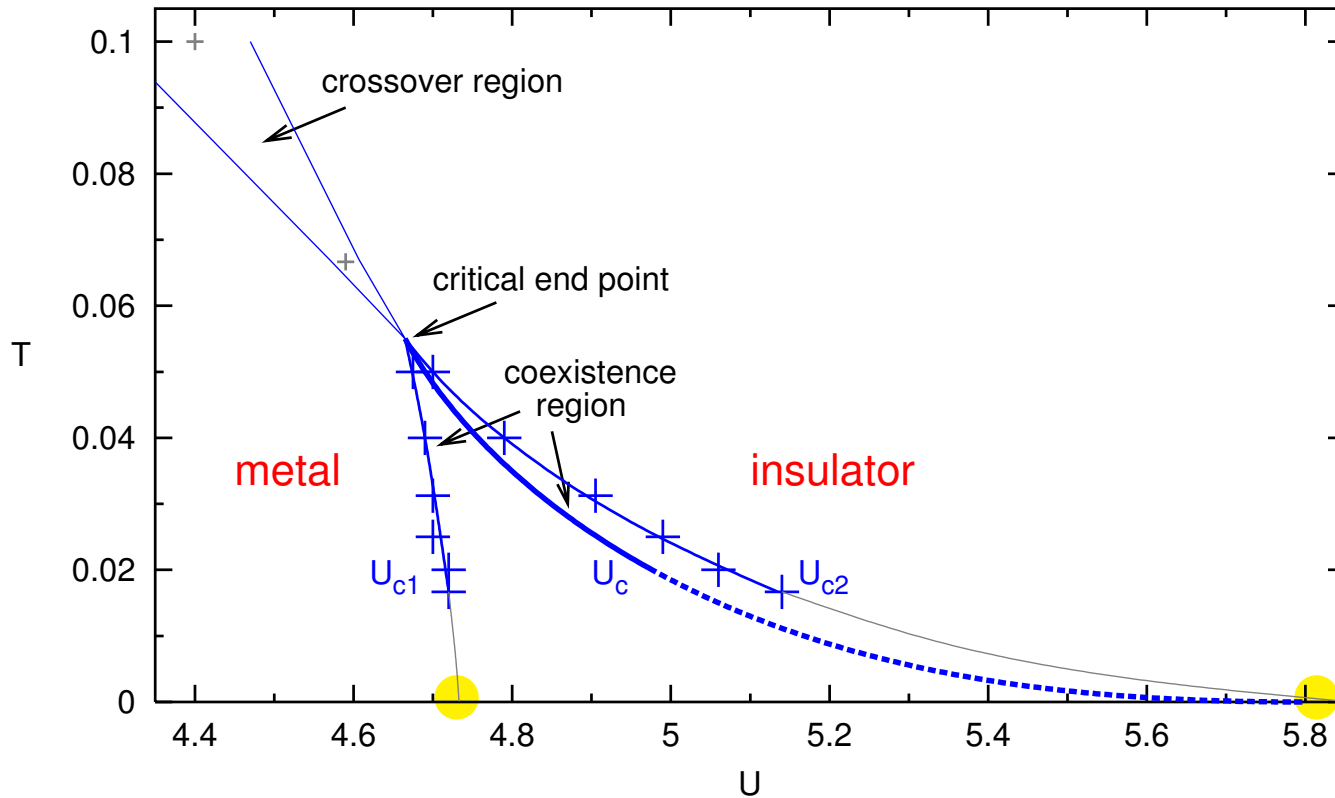
- a) "Ulmke smoothing",
- b) improved "Smoothing",
- c) cubic spline + analytic high-frequency corrections

large errors even for $\omega \rightarrow 0$ in a)

low-frequency errors of $\Sigma(\omega)$ small in b) and c)



Frustrated 1-band Hubbard model: 1st order MIT + coexistence



Georges and Krauth (1993)
 Rozenberg, Kotliar, Zhang (1994)
 Georges et al. (RMP, 1996)
 Schlipf et al. (1999)
 Rozenberg, Chitra, Kotliar (1999)
 Krauth (2000)
 Bulla (1999, 2001)
 Joo, Oudovenko (2001)
 Tong (2001)
 Blümer (2000, 2002)

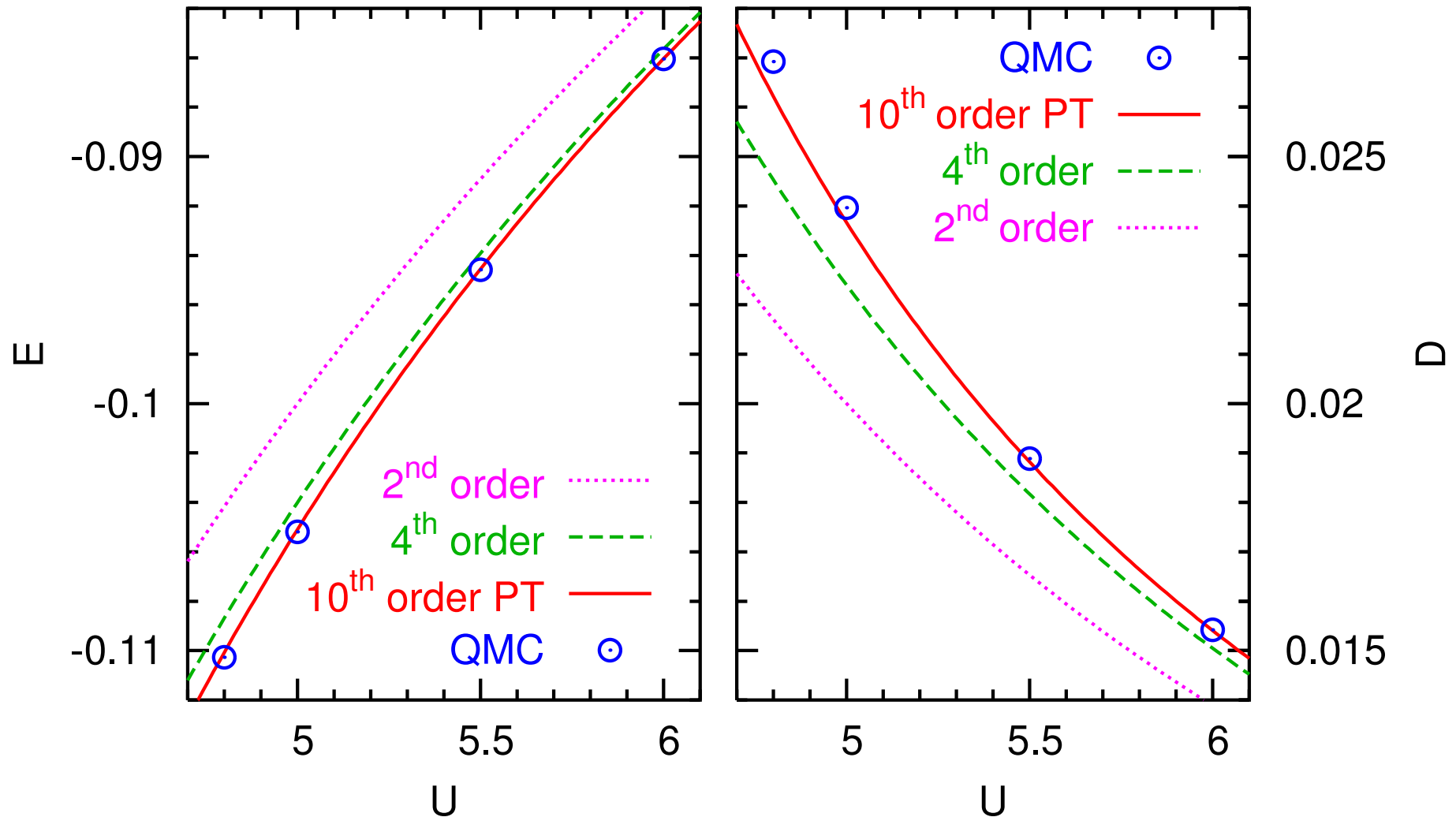
1st order line from
$$\frac{dU_c(T)}{dT} = f(T, U_c(T)); \quad f(T, U) := \frac{\Delta E(T, U)}{T \Delta D(T, U)}$$

low- T asymptotics from
$$U_c(T) = U_c^0 - \sqrt{\frac{2S_0 T}{a}} + \frac{\gamma_0}{4S_0} T + \mathcal{O}(T^{3/2})$$

High-precision energetics needed, even for $T \rightarrow 0$

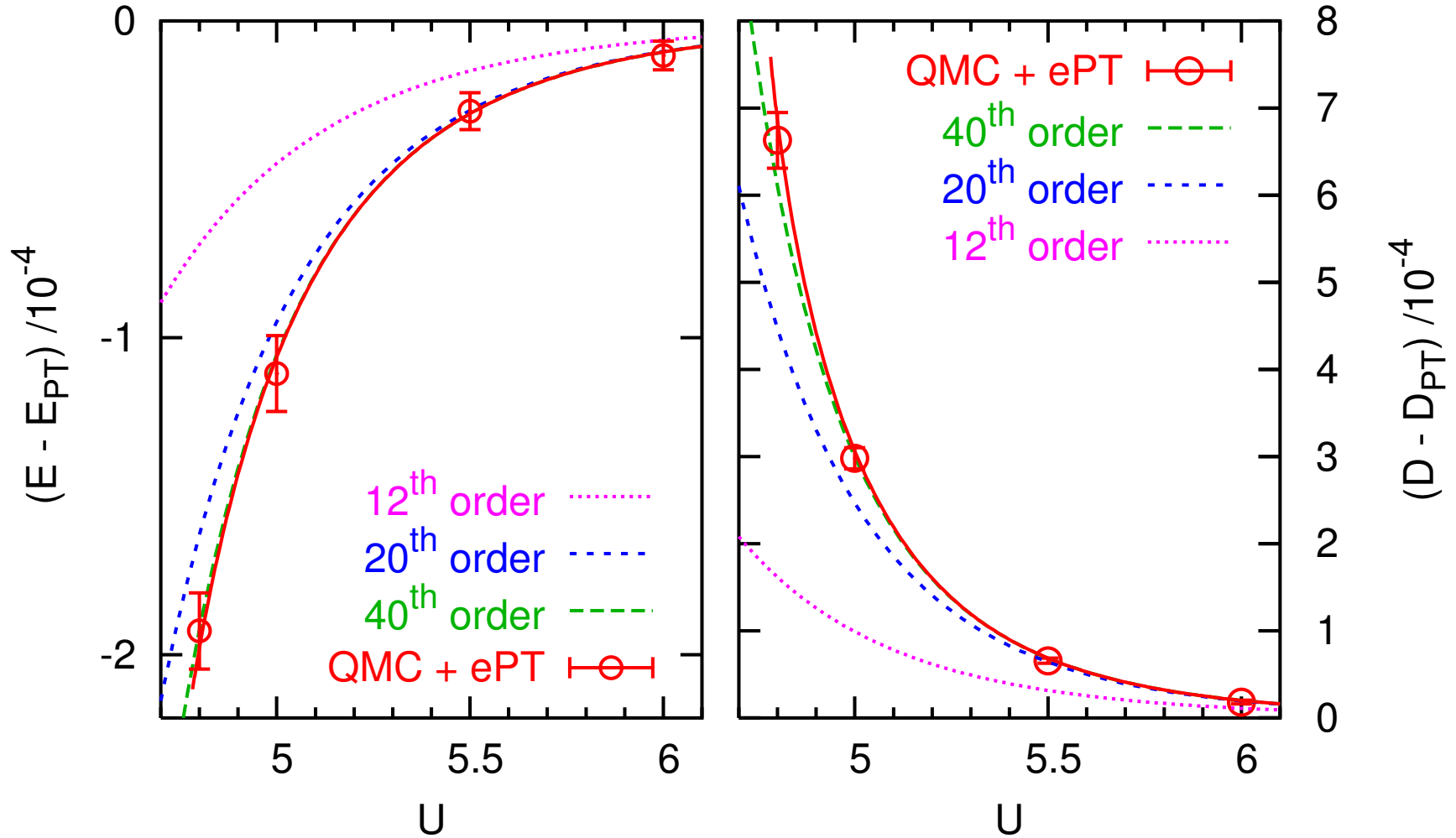
High-precision ground state estimates from QMC

QMC (+ high-frequency expansion) vs. strong-coupling PT for insulating phase



Excellent agreement at $U = 6.0$, deviations below.

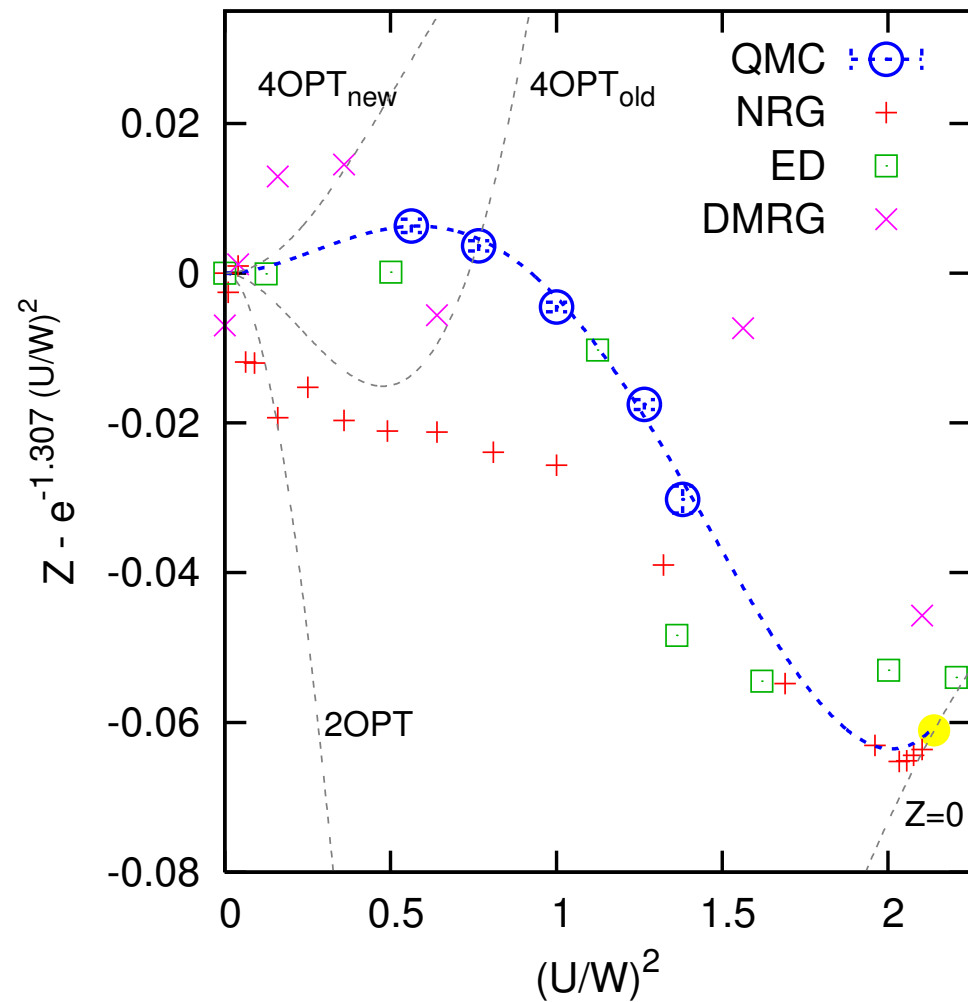
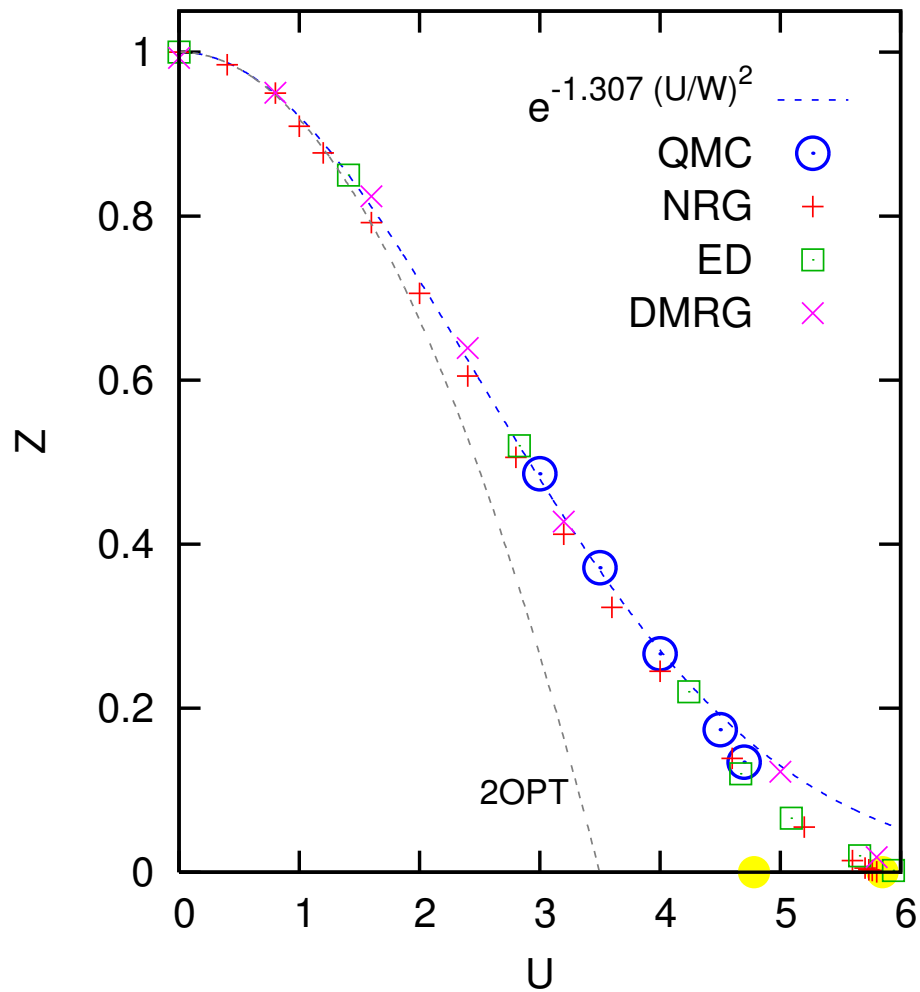
Higher resolution plots: differences w.r.t. 10th order PT



ePT: extrapolation of PT to infinite order [NB, Kalinowski, Phys. Rev B **71**, 195102 (2005)]
 \rightsquigarrow critical interaction U_{c1} , critical exponents, benchmark ($\Delta E \lesssim 10^{-6}$)

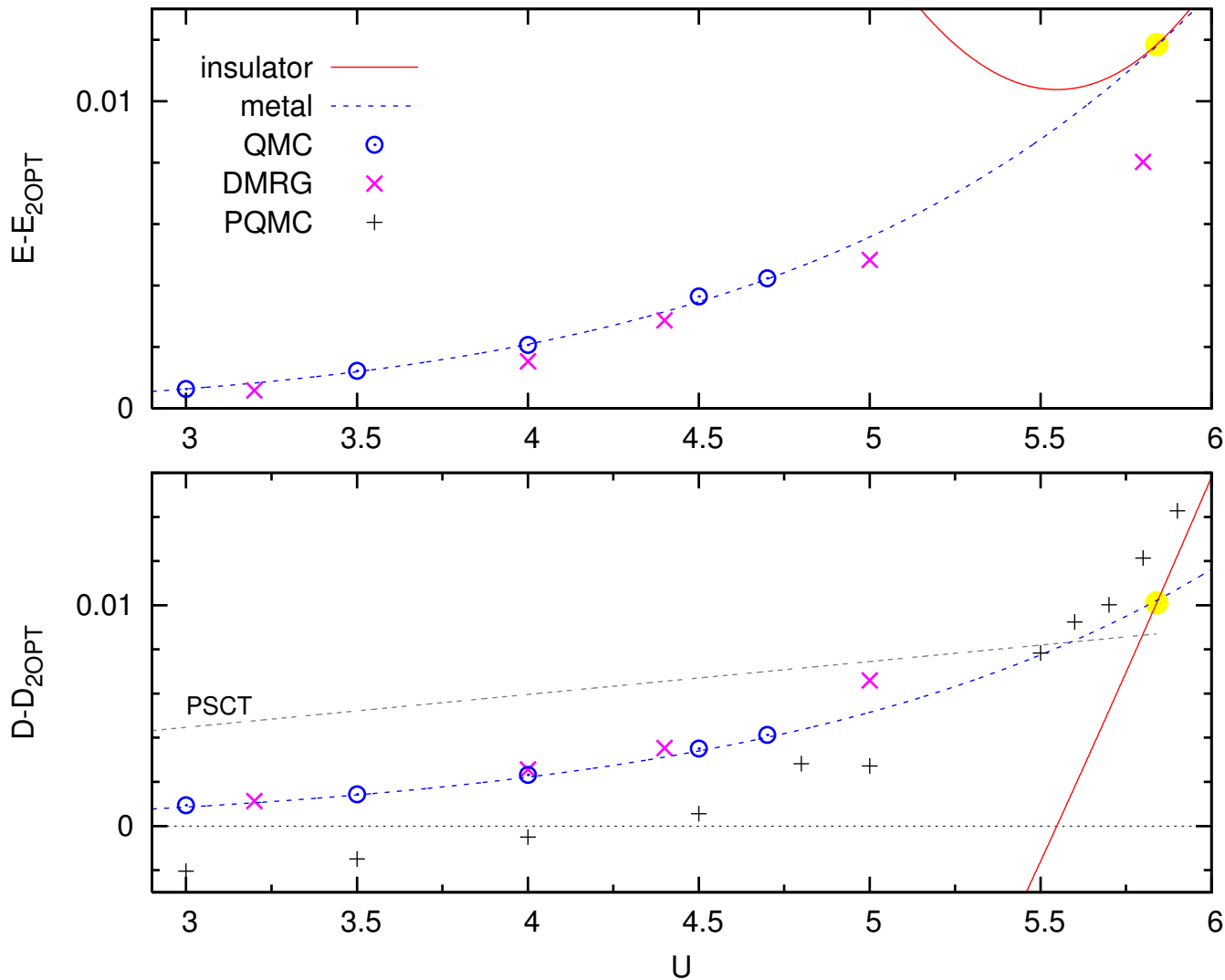
QMC algorithm has passed only available authoritative (1-band) test!

Quasiparticle weight/mass enhancement $Z = m/m^*$ in metal at $T = 0$



T -extrapolated QMC even beats all ground state methods!

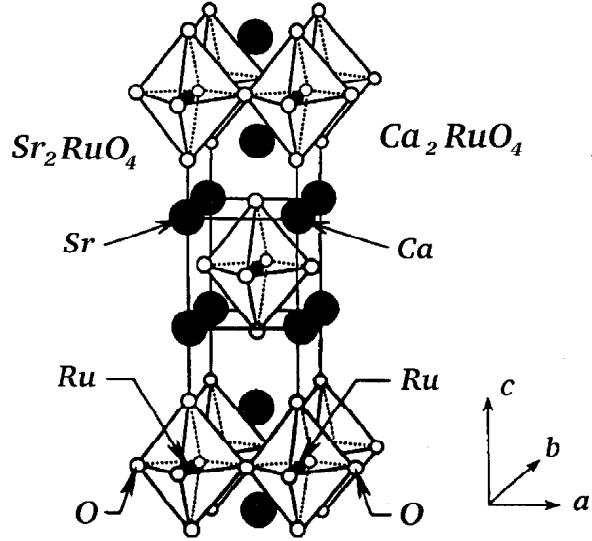
Energetics: differences w.r.t. 2nd order weak-coupling PT for E and D



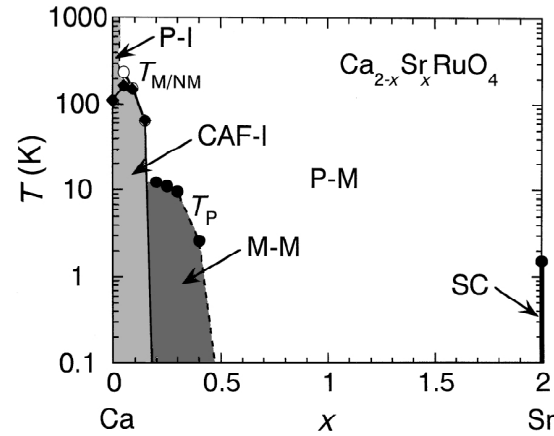
QMC-fit consistent with ePT for insulator; large deviations of PQMC, PSCT.

4th order PT coefficient corrected by QMC ($-62 \rightarrow +5$)

Orbital-selective Mott transition in 2-band Hubbard model

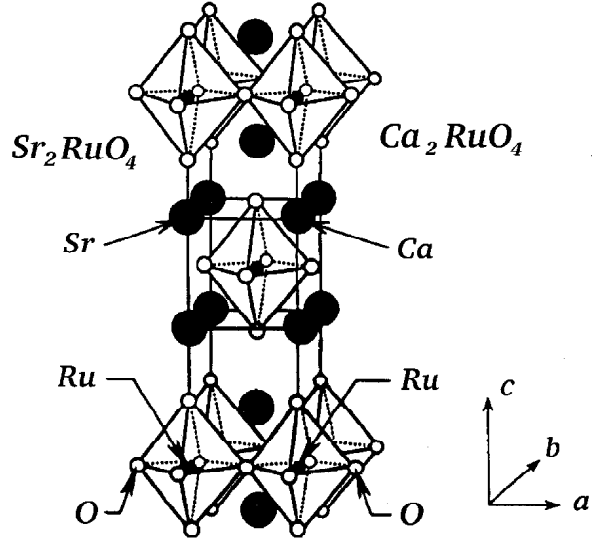


Motivation: $\text{Ca}_{2-x}\text{Sr}_x\text{RuO}_4$

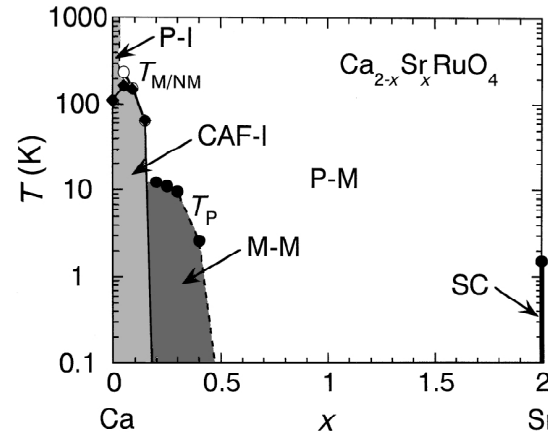


[Nakatsuj, Maeno (2002)]

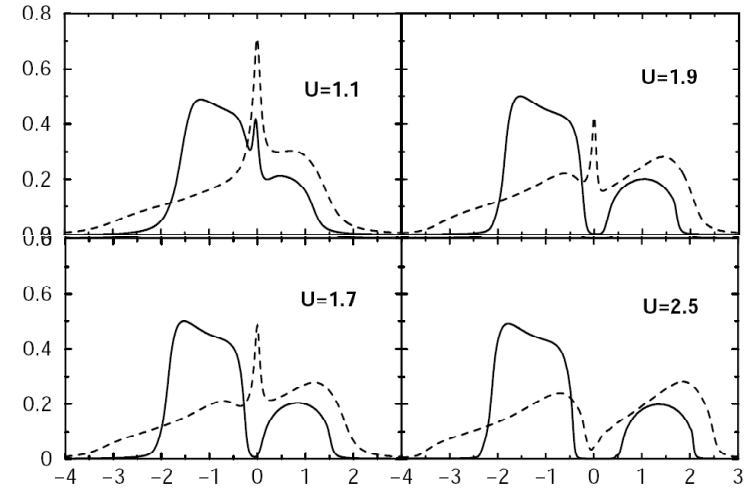
Orbital-selective Mott transition in 2-band Hubbard model



Motivation: $\text{Ca}_{2-x}\text{Sr}_x\text{RuO}_4$



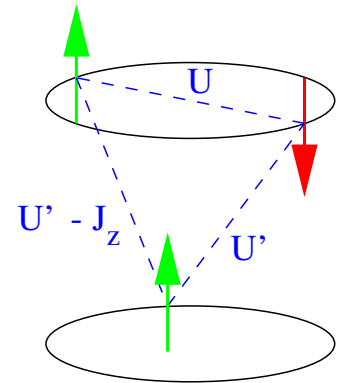
[Nakatsuj, Maeno (2002)]



[Anisimov et al. (2002)]

$$\begin{aligned}
 H = & \sum_{m=1}^2 \left[- \sum_{\langle ij \rangle \sigma} t_m c_{im\sigma}^\dagger c_{jm\sigma} + U \sum_i n_{im\uparrow} n_{im\downarrow} \right] \\
 & + \sum_{i\sigma\sigma'} (U' - \delta_{\sigma\sigma'} J_z) n_{i1\sigma} n_{i2\sigma'} \\
 & + \frac{1}{2} J_\perp \sum_{i\sigma} \left[c_{i1\sigma}^\dagger \left(c_{i2\bar{\sigma}}^\dagger c_{i1\bar{\sigma}} + c_{i1\bar{\sigma}}^\dagger c_{i2\bar{\sigma}} \right) c_{i2\sigma} + \text{h.c.} \right]
 \end{aligned}$$

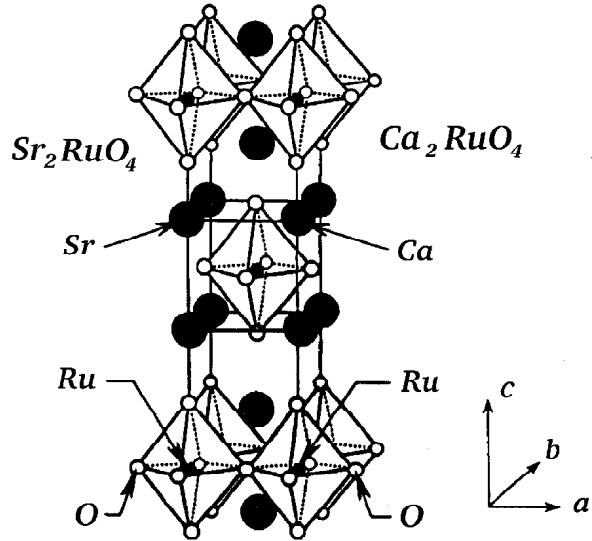
$m=2$



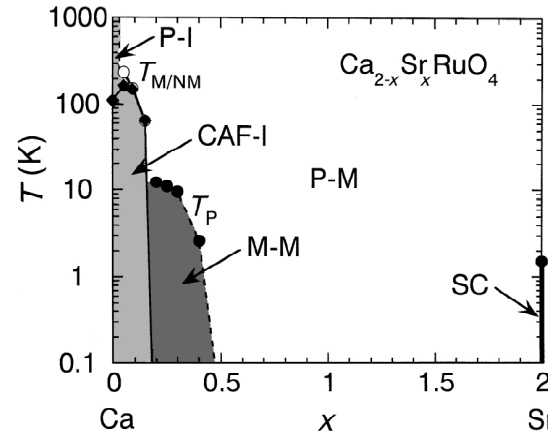
$m=1$

For Bethe DOS, $t_2 = 2t_1$: two 1st order MITs for $U' = J_z = J_\perp = 0$

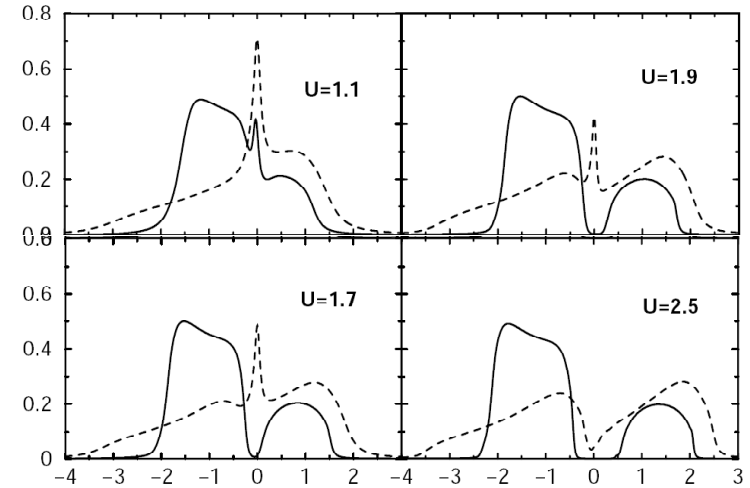
Orbital-selective Mott transition in 2-band Hubbard model



Motivation: $Ca_{2-x}Sr_xRuO_4$



[Nakatsuj, Maeno (2002)]



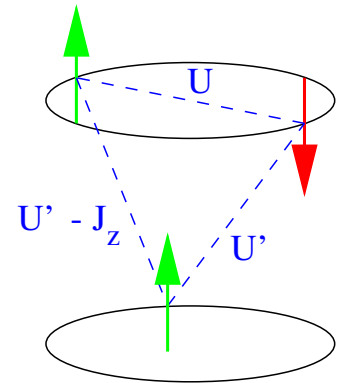
[Anisimov et al. (2002)]

$$H = \sum_{m=1}^2 \left[- \sum_{\langle ij \rangle \sigma} t_m c_{im\sigma}^\dagger c_{jm\sigma} + U \sum_i n_{im\uparrow} n_{im\downarrow} \right]$$

$$+ \sum_{i\sigma\sigma'} (U' - \delta_{\sigma\sigma'} J_z) n_{i1\sigma} n_{i2\sigma'}$$

$$+ \frac{1}{2} J_\perp \sum_{i\sigma} \left[c_{i1\sigma}^\dagger (c_{i2\bar{\sigma}}^\dagger c_{i1\bar{\sigma}} + c_{i1\bar{\sigma}}^\dagger c_{i2\bar{\sigma}}) c_{i2\sigma} + \text{h.c.} \right]$$

m = 2



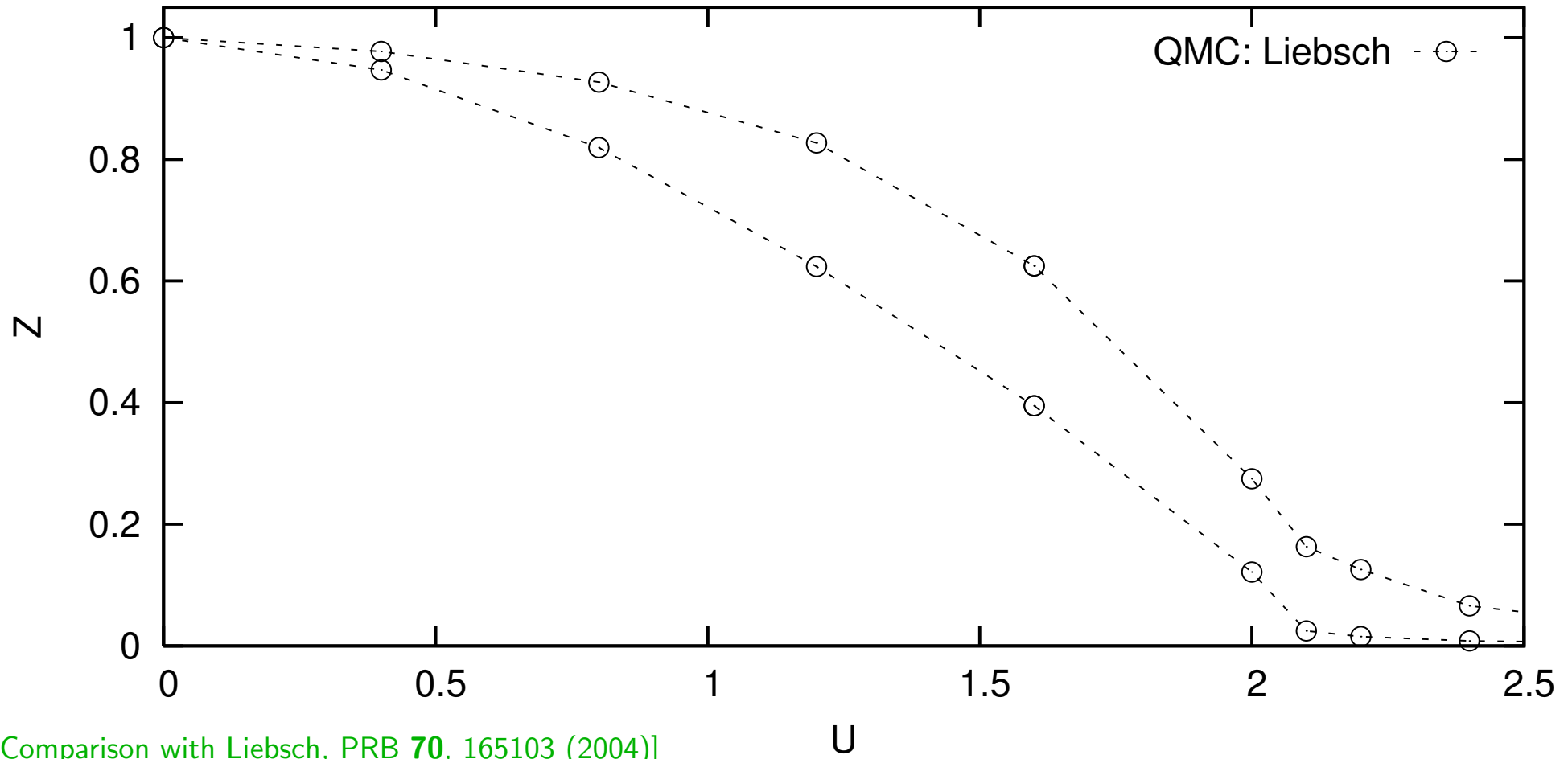
m = 1

For Bethe DOS, $t_2 = 2t_1$: two 1st order MITs for $U' = J_z = J_\perp = 0$

two distinct MITs for $J_z = J_\perp = U/4$

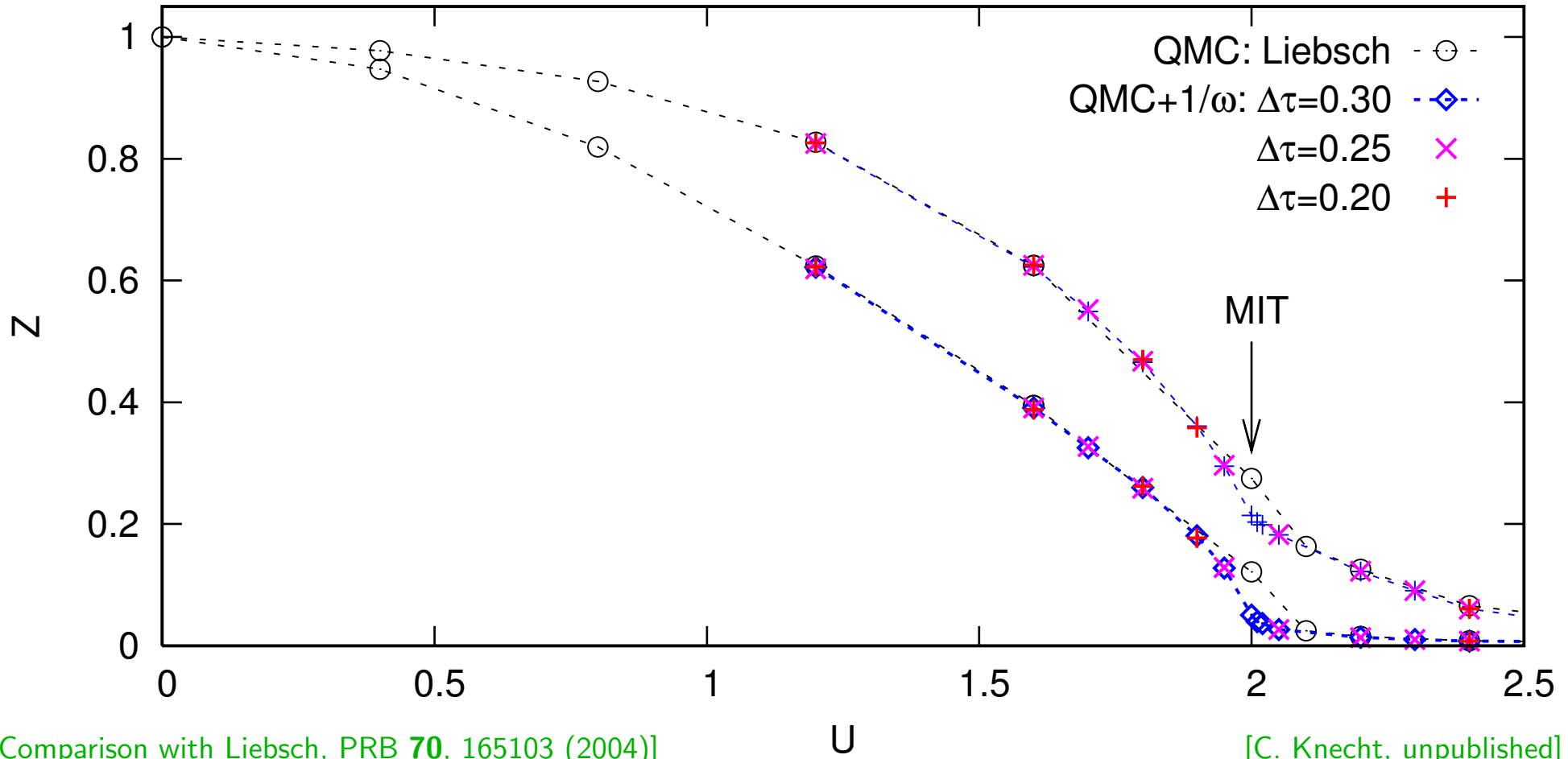
single MIT or **OSMTs** for $J_\perp = 0, J_z = U/4$?

Test for multiband-QMC: quasiparticle weights $Z = m/m^*$ in 2-band model



[Comparison with Liebsch, PRB **70**, 165103 (2004)]

Test for multiband-QMC: quasiparticle weights $Z = m/m^*$ in 2-band model

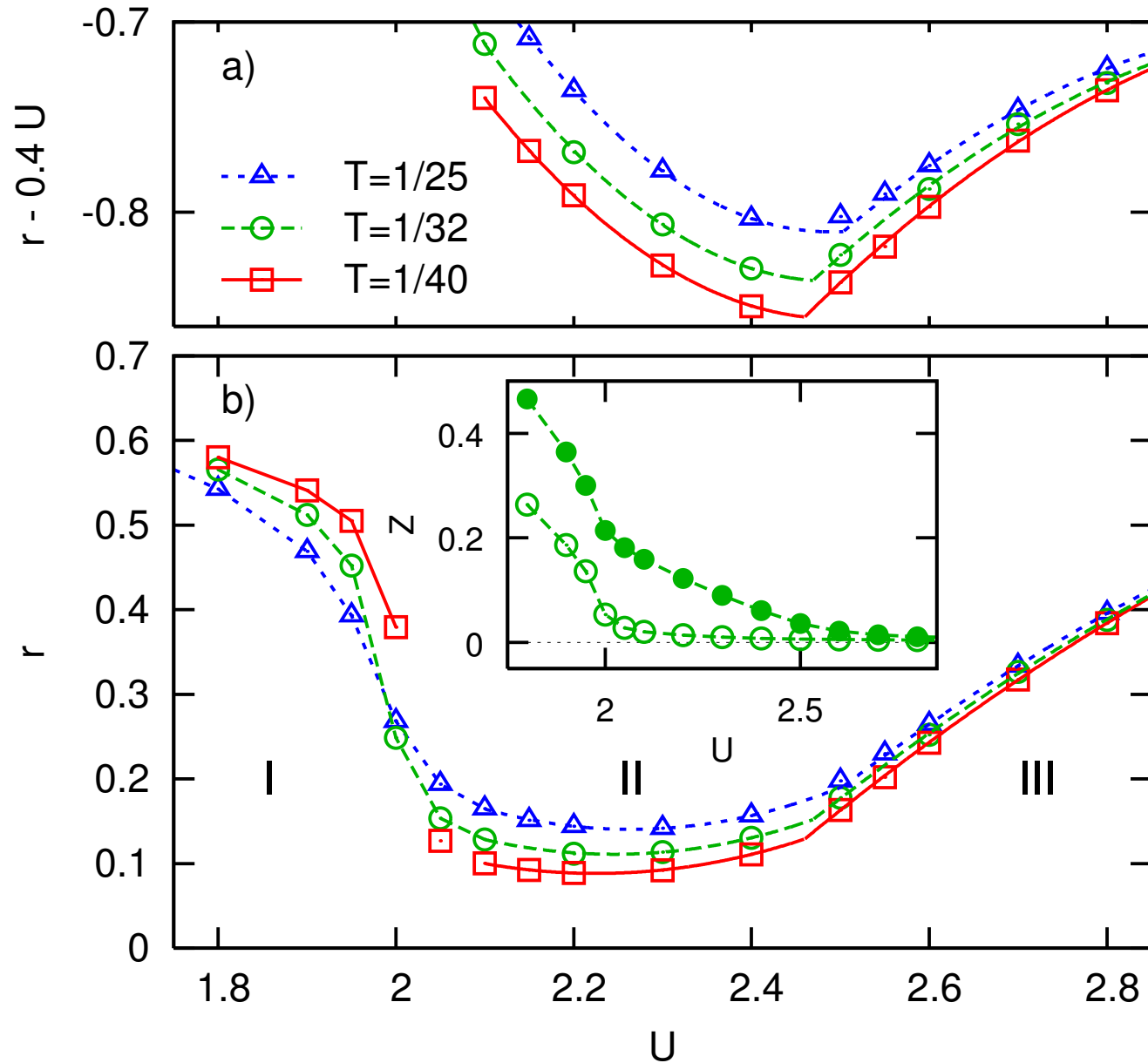


Very small dependence on discretization $\Delta\tau$.

Conclusion in 3/2005: New algorithm clearly exposes (single) metal-insulator transition (MIT)

But: wide band still "quite metallic" for $U > 2.0$ – 2nd transition?

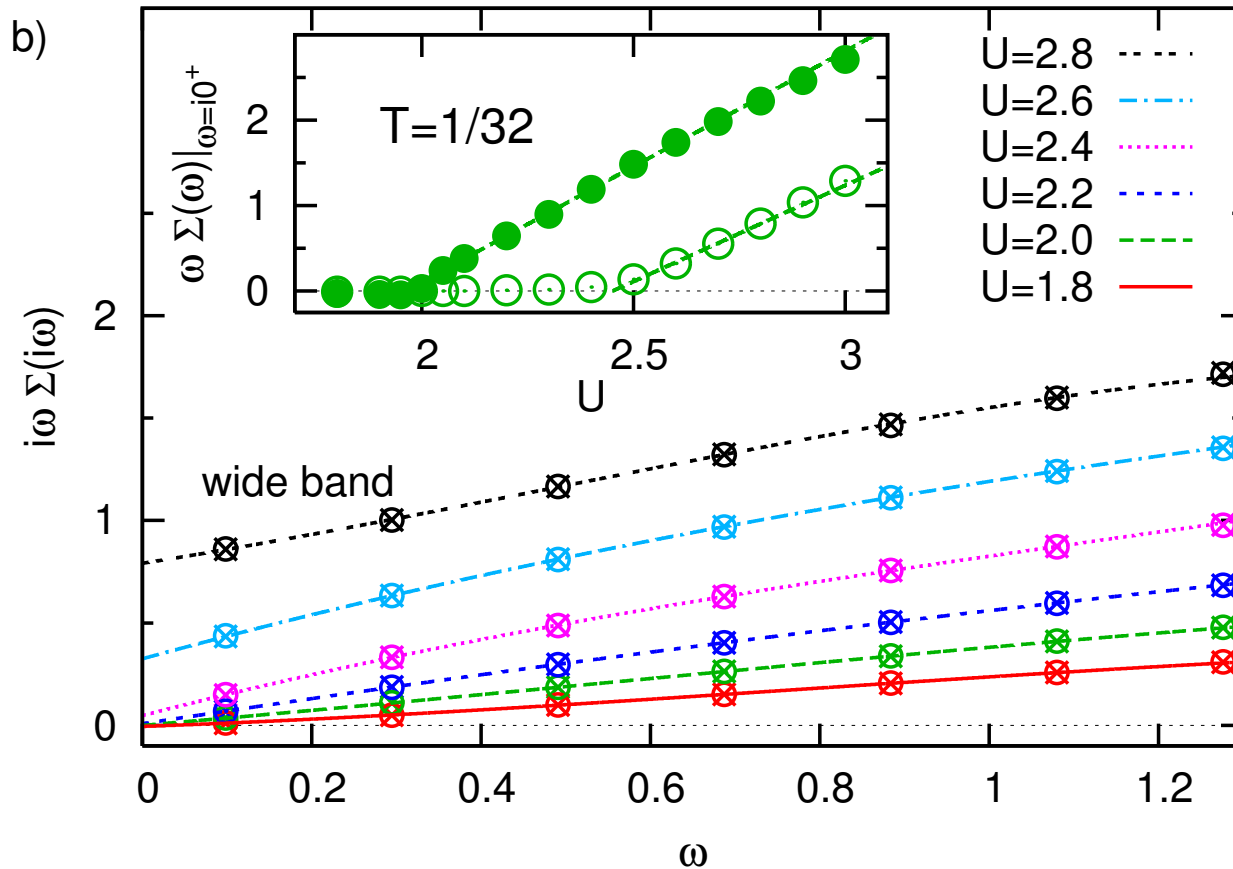
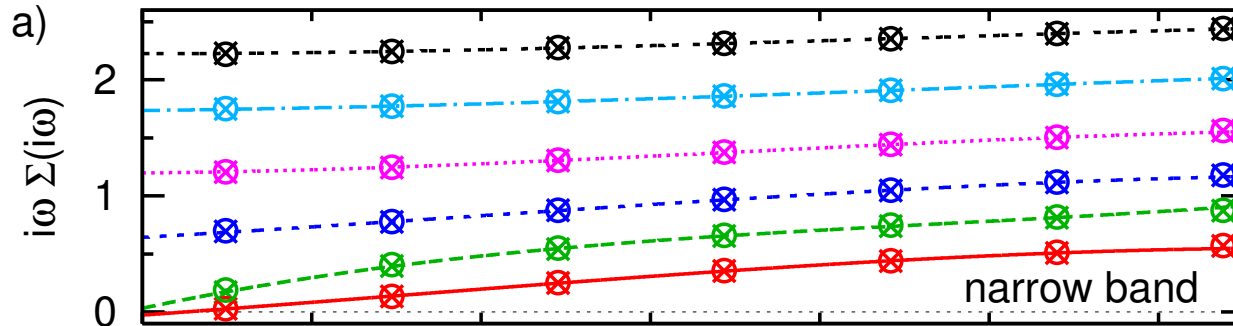
Ratio of quasiparticle weights $r = Z_{\text{narrow}}/Z_{\text{wide}}$



3 regions of different character

kinks indicate 2nd transition at $U \approx 2.5$

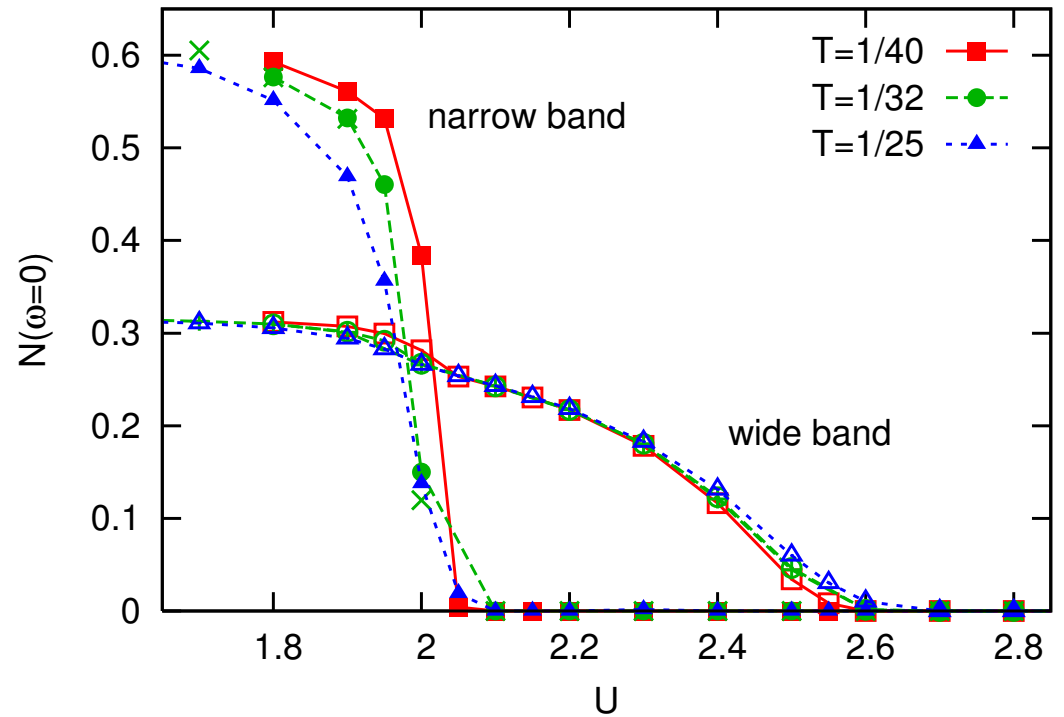
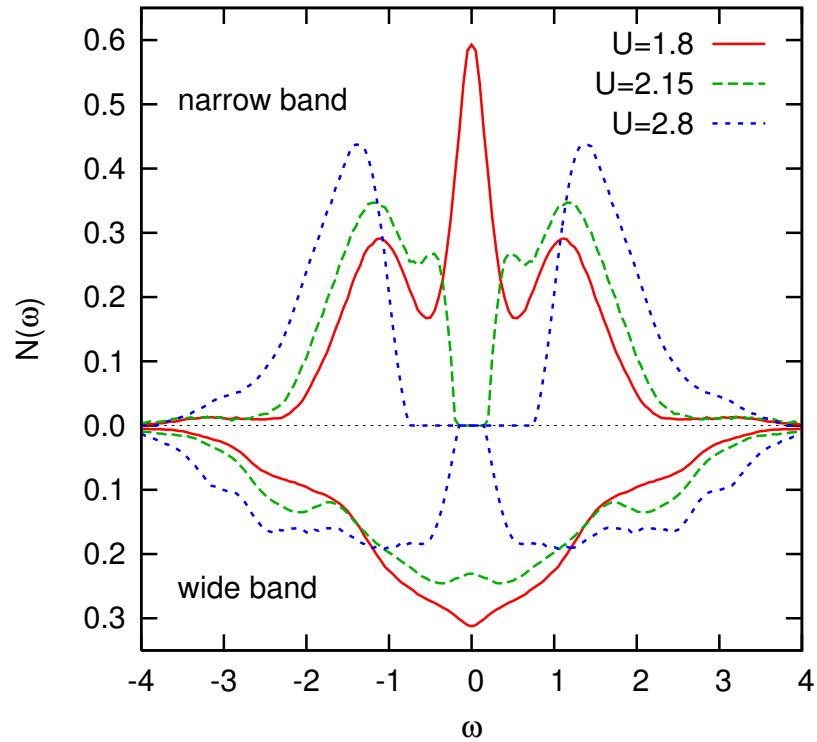
Low-frequency analysis of self-energy



for regular self-energy:
 $\omega \Sigma(\omega) \xrightarrow{\omega \rightarrow 0} 0$

singularities (\sim gap) for
 $U > 2, U > 2.5$

Spectral function (interacting DOS)



Clear indications for second singularity

Wide band remains metallic at $U \approx 2.0$

⇒ **two orbital-selective Mott transitions**

[Knecht, NB, van Dongen, cond-mat/0505106]

Summary and Outlook

DMFT+QMC: valuable numerical approach for correlated electron systems

- ab initio approach in combination with DFT(LDA)
- even broader applicability for cluster extensions
- Fourier transformation scheme crucial for reliability and efficiency

Mott transition in frustrated 1-band Hubbard model

High-precision ground state estimates from QMC

Orbital-selective Mott transition in 2-band Hubbard model

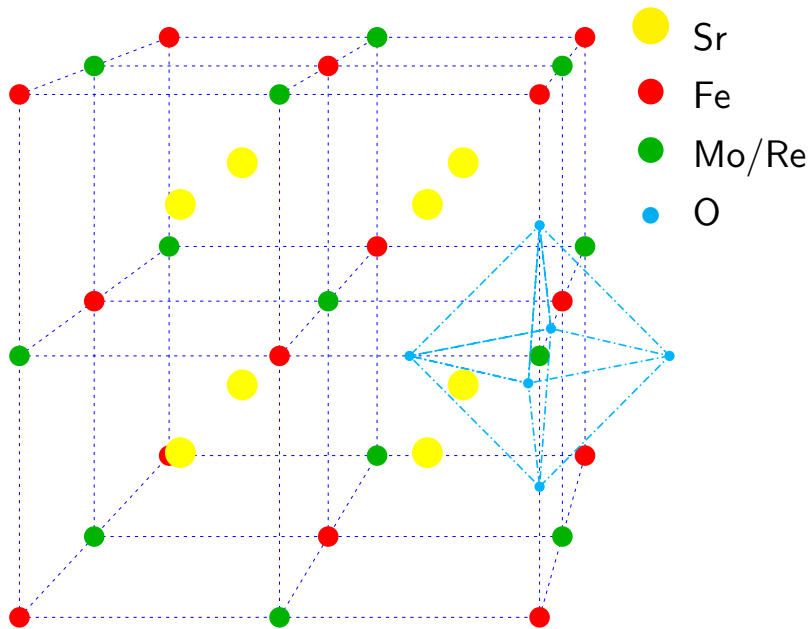
Outlook: realistic material-specific calculations with LDA+DMFT

Thanks to: Carsten Knecht, Krunoslav Pozgajcic, Peter van Dongen

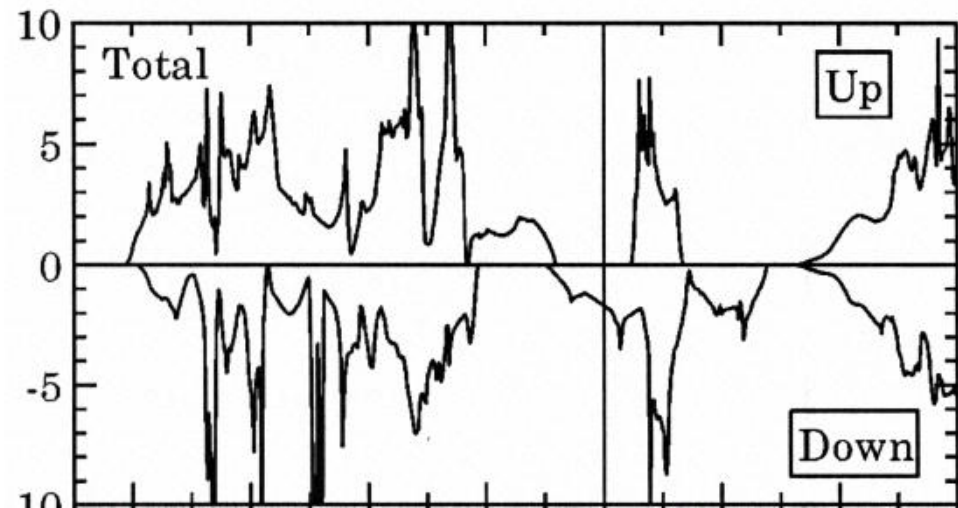
NIC Jülich, DFG (BI775/1)

Outlook: theory of half-metallic double perovskites

$\text{Sr}_2\text{FeMoO}_6$ and $\text{Sr}_2\text{FeReO}_6$



Valences: Sr^{2+} [Kr]
 Fe^{3+} [Ar] $3d^5$
 Mo^{5+} [Kr] $4d^1$
 O^{2-} [Ne]



[LSDA+U for $\text{Sr}_2\text{FeMoO}_6$, Saitoh et al. (2002)]

$$\begin{aligned}
 H = & \epsilon^f \sum_{i\alpha} n_{i\alpha}^f + \epsilon^m \sum_{i\alpha} n_{i\alpha}^m + \sum_{i, \alpha \neq \alpha'} U_{\alpha\alpha'}^f n_{i\alpha}^f n_{i\alpha'}^f + \sum_{j, \alpha \neq \alpha'} U_{\alpha\alpha'}^m n_{j\alpha}^m n_{j\alpha'}^m \\
 & + \sum_{\langle ij \rangle \alpha} t^{fm} (f_{i\alpha}^\dagger m_{j\alpha} + \text{hc}) + \sum_{\langle jj' \rangle \alpha} t^{mm} m_{j\alpha}^\dagger m_{j'\alpha} + \sum_{\langle ii' \rangle \alpha} t^{ff} f_{i\alpha}^\dagger f_{i'\alpha}
 \end{aligned}$$

DPG research group 559 on “New materials with high spin polarization”

Speakers: Felser (Mainz), Hillebrands (Kaiserslautern)

List of projects

- | | | |
|----|-------------------------|---|
| 1 | Felser | Synthesis: Heusler compounds |
| 2 | Jacob, Adrian | Thin films |
| 3 | Jourdan, Jacob, Adrian | Tunnel spectroscopy |
| 4 | Tremel | Synthesis: double perovskites |
| 5 | Elmers | Surface magnetization |
| 6 | Blümer, van Dongen | Theory of double perovskites: LDA+DMFT(QMC) |
| 7 | Schönhense, Felser | Spin resolved photoemission and DFT(LDA) |
| 8 | Ksenofontov, Felser | Mößbauer spectroscopy |
| 9 | Demokritov, Hillebrands | Brillouin light scattering spectroscopy |
| 10 | Aeschlimann, Bauer | Spectroscopy of unoccupied states (2PPE) |

Additional material for discussions

Comparison at $T = 1/32$ with Liebsch, PRB **70**, 165103 (2004)

