

Critical exponents in strongly correlated electron systems

Nils Blümer, KOMET Theory

Critical exponents in strongly correlated electron systems

Nils Blümer, KOMET Theory

The trick

(i) mathematical puzzles

1	2	4	8	16	?
---	---	---	---	----	---

Critical exponents in strongly correlated electron systems

Nils Blümer, KOMET Theory

The trick

(i) mathematical puzzles

1	2	4	8	16	32
---	---	---	---	----	----

4	2	5	2	6	?
---	---	---	---	---	---

Critical exponents in strongly correlated electron systems

Nils Blümer, KOMET Theory

The trick

(i) mathematical puzzles

1	2	4	8	16	32
---	---	---	---	----	----

4	2	5	2	6	2
---	---	---	---	---	---

3	1	4	1	5	?
---	---	---	---	---	---

Critical exponents in strongly correlated electron systems

Nils Blümer, KOMET Theory

The trick

(i) mathematical puzzles

1	2	4	8	16	32
---	---	---	---	----	----

4	2	5	2	6	2
---	---	---	---	---	---

3	1	4	1	5	9
---	---	---	---	---	---

- explicit task: calculate next element(s)

Critical exponents in strongly correlated electron systems

Nils Blümer, KOMET Theory

The trick

(i) mathematical puzzles

1 2 4 8 16 32 ? ? ? ...

4 2 5 2 6 2

3 1 4 1 5 9

- explicit task: calculate next element(s)
- implicit task: derive law (determining all elements)
- hypothesis testable by prediction of selected elements

(ii) Taylor expansions and perturbation series

$$f(x) = \sum_{k=0}^{\infty} a_k x^k$$

Function $f(x)$ fully determined by coefficients $\{a_k\}_{k=0}^{\infty}$ within radius of convergence

(ii) Taylor expansions and perturbation series

$$f(x) = \sum_{k=0}^{\infty} a_k x^k$$

Function $f(x)$ fully determined by coefficients $\{a_k\}_{k=0}^{\infty}$ within radius of convergence

a) $\{a_k\} = \left[\frac{1}{1} \right] \left[\frac{1}{2} \right] \left[\frac{1}{6} \right] \left[\frac{1}{24} \right] \left[\frac{1}{120} \right] \dots \left[\frac{1}{k!} \right] \rightsquigarrow f(x) = \sum_{k=0}^{\infty} \frac{x^k}{k!} = e^x$ exponential function

(ii) Taylor expansions and perturbation series

$$f(x) = \sum_{k=0}^{\infty} a_k x^k$$

Function $f(x)$ fully determined by coefficients $\{a_k\}_{k=0}^{\infty}$ within radius of convergence

$$\text{a) } \{a_k\} = \left[\frac{1}{1} \right] \left[\frac{1}{2} \right] \left[\frac{1}{6} \right] \left[\frac{1}{24} \right] \left[\frac{1}{120} \right] \dots \left[\frac{1}{k!} \right]$$

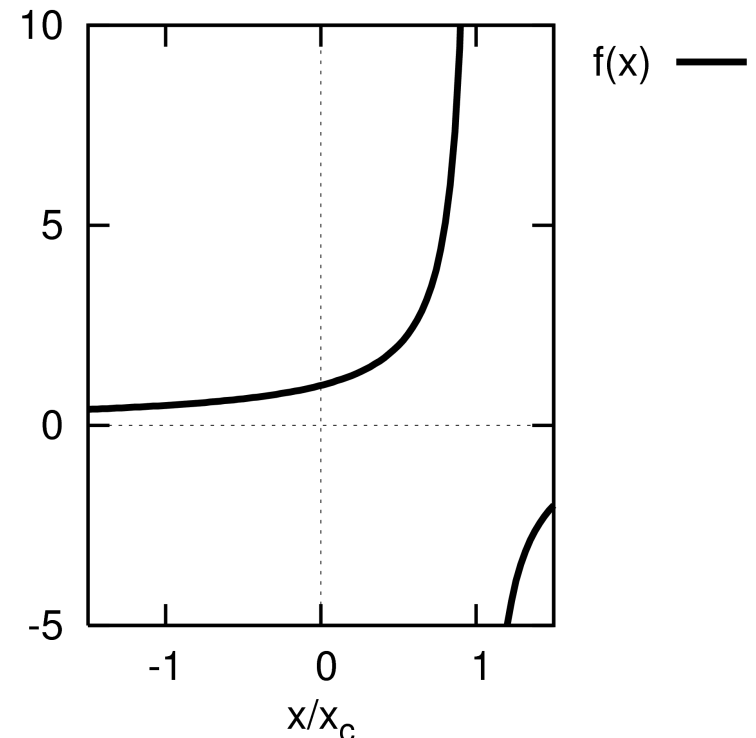
$$\rightsquigarrow f(x) = \sum_{k=0}^{\infty} \frac{x^k}{k!} = e^x \quad \text{exponential function}$$

$$\text{b) } \{a_k\} = \left[q^0 \right] \left[q^1 \right] \left[q^2 \right] \left[q^3 \right] \left[q^4 \right] \dots \left[q^k \right]$$

$$\rightsquigarrow f(x) = \sum_{k=0}^{\infty} (qx)^k = \frac{1}{1 - qx} = \frac{1}{q} (x_c - x)^{-1}$$

for $|x| < x_c = q^{-1}$ geometric series

Note: singularity in $f(x)$ at x_c , exponent -1



(ii) Taylor expansions and perturbation series

$$f(x) = \sum_{k=0}^{\infty} a_k x^k$$

Function $f(x)$ fully determined by coefficients $\{a_k\}_{k=0}^{\infty}$ within radius of convergence

$$\text{a) } \{a_k\} = \left[\frac{1}{1} \right] \left[\frac{1}{2} \right] \left[\frac{1}{6} \right] \left[\frac{1}{24} \right] \left[\frac{1}{120} \right] \dots \left[\frac{1}{k!} \right]$$

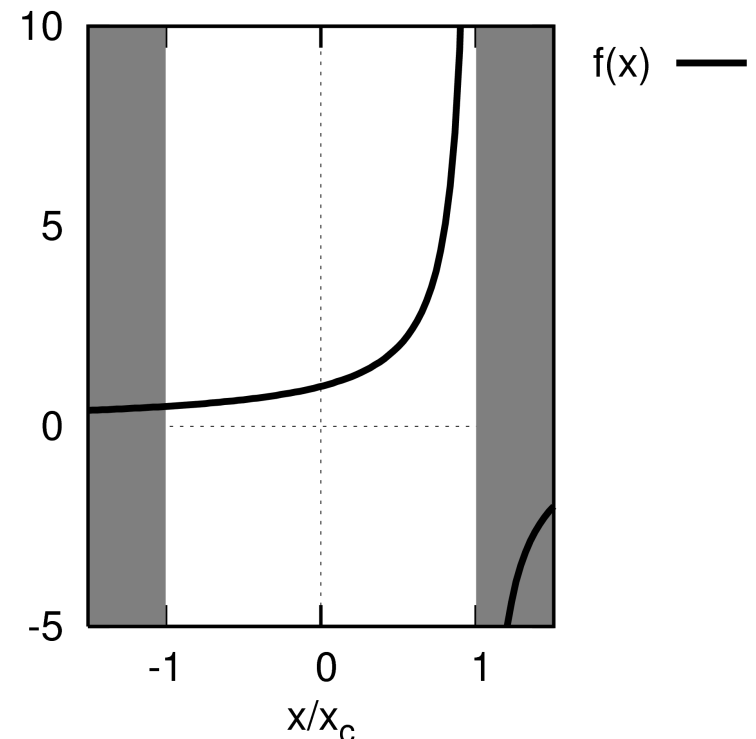
$$\rightsquigarrow f(x) = \sum_{k=0}^{\infty} \frac{x^k}{k!} = e^x \quad \text{exponential function}$$

$$\text{b) } \{a_k\} = \left[q^0 \right] \left[q^1 \right] \left[q^2 \right] \left[q^3 \right] \left[q^4 \right] \dots \left[q^k \right]$$

$$\rightsquigarrow f(x) = \sum_{k=0}^{\infty} (qx)^k = \frac{1}{1 - qx} = \frac{1}{q} (x_c - x)^{-1}$$

for $|x| < x_c = q^{-1}$ geometric series

Note: singularity in $f(x)$ at x_c , exponent -1



(ii) Taylor expansions and perturbation series

$$f(x) = \sum_{k=0}^{\infty} a_k x^k$$

Function $f(x)$ fully determined by coefficients $\{a_k\}_{k=0}^{\infty}$ within radius of convergence

a) $\{a_k\} = \left[\frac{1}{1} \right] \left[\frac{1}{2} \right] \left[\frac{1}{6} \right] \left[\frac{1}{24} \right] \left[\frac{1}{120} \right] \dots \left[\frac{1}{k!} \right] \rightsquigarrow f(x) = \sum_{k=0}^{\infty} \frac{x^k}{k!} = e^x$ exponential function

b) $\{a_k\} = \left[q^0 \right] \left[q^1 \right] \left[q^2 \right] \left[q^3 \right] \left[q^4 \right] \dots \left[q^k \right]$

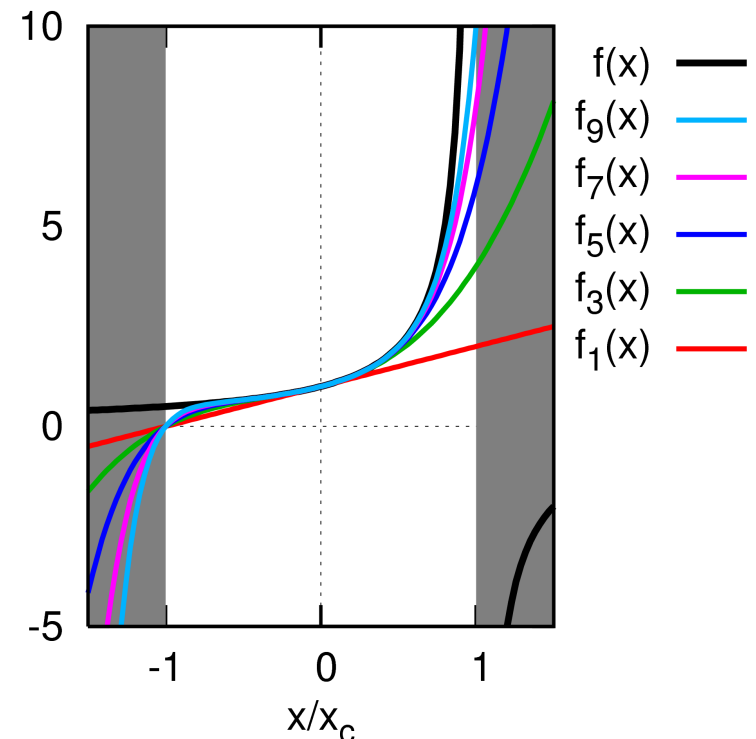
$$\rightsquigarrow f(x) = \sum_{k=0}^{\infty} (qx)^k = \frac{1}{1 - qx} = \frac{1}{q} (x_c - x)^{-1}$$

for $|x| < x_c = q^{-1}$ geometric series

Note: singularity in $f(x)$ at x_c , exponent -1

However: finite order approximations

$$f_n(x) = \sum_{k=0}^n a_k x^k \text{ regular for all } x$$



(iii) Essence of extrapolated perturbation theory (ePT)

1) obtain as many coefficients as possible

2) derive law

3) test law

4) characterize singularity and extract critical exponents

(iii) Essence of extrapolated perturbation theory (ePT)

1) obtain as many coefficients as possible – from strong-coupling perturbation theory

$$\boxed{a_2} \quad \boxed{a_4} \quad \boxed{a_6} \quad \boxed{a_8} \quad \boxed{a_{10}}$$

2) derive law

3) test law

4) characterize singularity and extract critical exponents

(iii) Essence of extrapolated perturbation theory (ePT)

1) obtain as many coefficients as possible – from strong-coupling perturbation theory

$$\boxed{a_2} \quad \boxed{a_4} \quad \boxed{a_6} \quad \boxed{a_8} \quad \boxed{a_{10}}$$

2) derive law – in good (?) approximation

$$\boxed{a_2} \quad \boxed{a_4} \quad \boxed{a_6} \quad \boxed{a_8} \quad \boxed{a_{10}} \quad \boxed{\tilde{a}_k} \quad \tilde{a}_k \sim k^\tau q^k$$

3) test law

4) characterize singularity and extract critical exponents

(iii) Essence of extrapolated perturbation theory (ePT)

1) obtain as many coefficients as possible – from strong-coupling perturbation theory

$$\boxed{a_2} \quad \boxed{a_4} \quad \boxed{a_6} \quad \boxed{a_8} \quad \boxed{a_{10}}$$

2) derive law – in good (?) approximation

$$\boxed{a_2} \quad \boxed{a_4} \quad \boxed{a_6} \quad \boxed{a_8} \quad \boxed{a_{10}} \quad \boxed{\tilde{a}_k} \quad \tilde{a}_k \sim k^\tau q^k$$

3) test law – comparison with quantum Monte Carlo

$$\boxed{a_2} \quad \boxed{a_4} \quad \boxed{a_6} \quad \boxed{a_8} \quad \boxed{a_{10}} \quad \boxed{\tilde{a}_{12}} \quad \boxed{\tilde{a}_{14}} \quad \boxed{\tilde{a}_{16}} \quad \boxed{\tilde{a}_{18}} \quad \boxed{\tilde{a}_{20}} \quad \boxed{\tilde{a}_{22}} \quad \boxed{\tilde{a}_{24}} \quad \boxed{\tilde{a}_{26}} \quad \boxed{\tilde{a}_{28}} \quad \boxed{\tilde{a}_{30}} \rightsquigarrow \text{observables}$$

4) characterize singularity and extract critical exponents

(iii) Essence of extrapolated perturbation theory (ePT)

1) obtain as many coefficients as possible – from strong-coupling perturbation theory

$$\boxed{a_2} \quad \boxed{a_4} \quad \boxed{a_6} \quad \boxed{a_8} \quad \boxed{a_{10}}$$

2) derive law – in good (?) approximation

$$\boxed{a_2} \quad \boxed{a_4} \quad \boxed{a_6} \quad \boxed{a_8} \quad \boxed{a_{10}} \quad \boxed{\tilde{a}_k} \quad \tilde{a}_k \sim k^\tau q^k$$

3) test law – comparison with quantum Monte Carlo

$$\boxed{a_2} \quad \boxed{a_4} \quad \boxed{a_6} \quad \boxed{a_8} \quad \boxed{a_{10}} \quad \boxed{\tilde{a}_{12}} \quad \boxed{\tilde{a}_{14}} \quad \boxed{\tilde{a}_{16}} \quad \boxed{\tilde{a}_{18}} \quad \boxed{\tilde{a}_{20}} \quad \boxed{\tilde{a}_{22}} \quad \boxed{\tilde{a}_{24}} \quad \boxed{\tilde{a}_{26}} \quad \boxed{\tilde{a}_{28}} \quad \boxed{\tilde{a}_{30}} \rightsquigarrow \text{observables}$$

4) characterize singularity and extract critical exponents

$$f^{\text{irr}}(x) \propto |x - q^{-1}|^{\tau-1} \dots$$

Critical exponents in strongly correlated electron systems

Nils Blümer, KOMET Theory

Outline

Essence of extrapolated perturbation theory (ePT)

Phase transitions and critical exponents

Strongly correlated electron systems

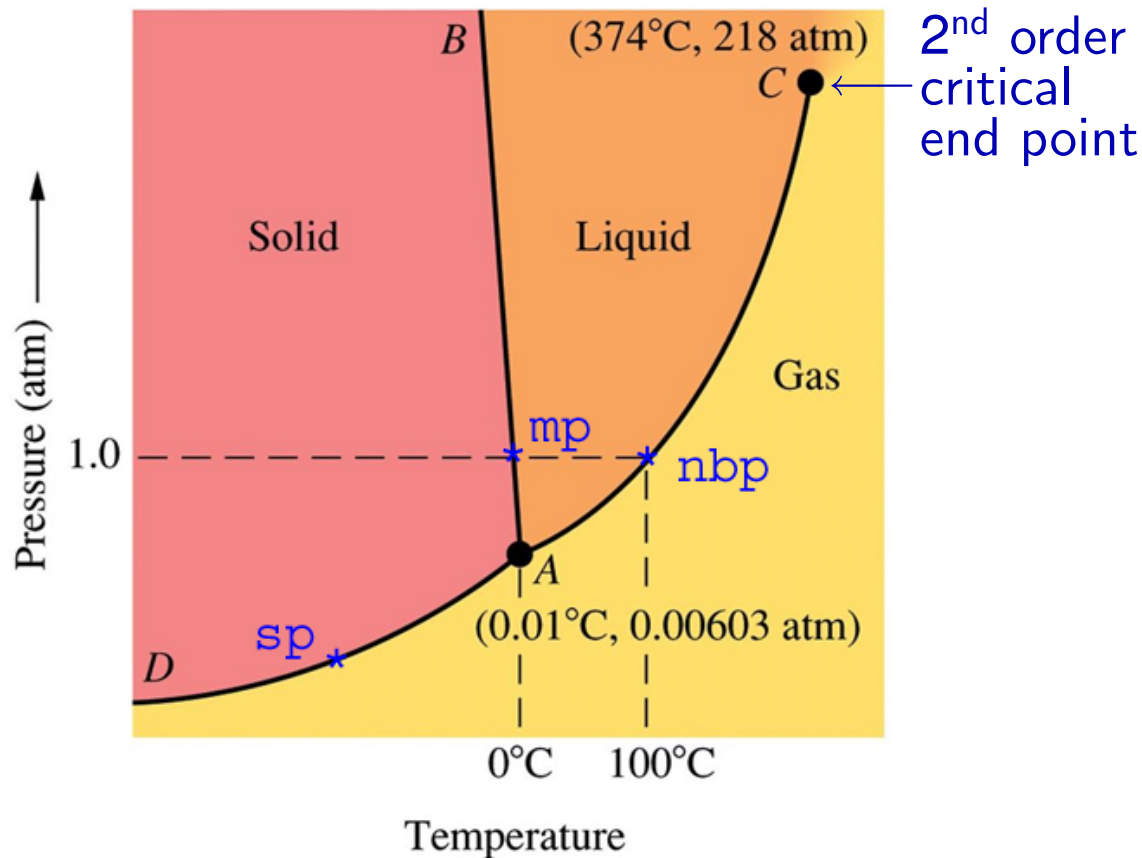
Mott metal insulator transition

Critical exponents

Summary

Phase transitions and critical exponents

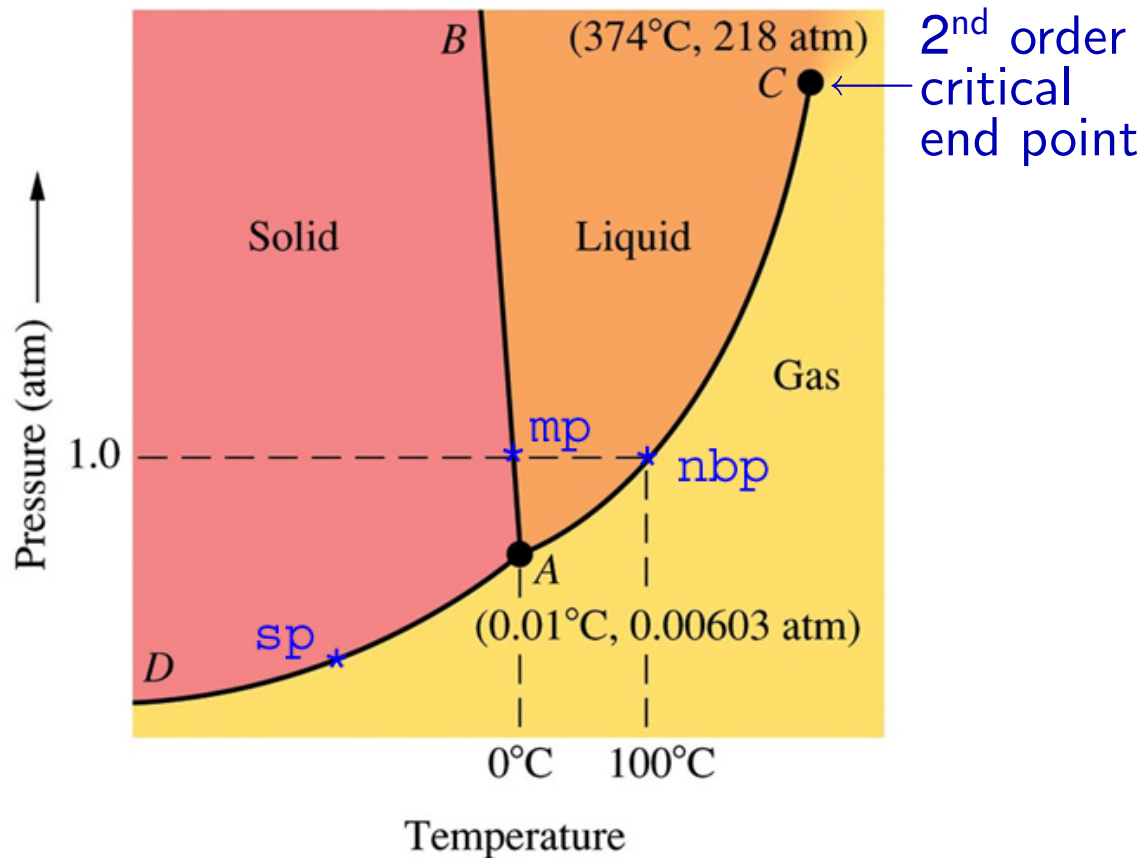
Liquid-gas transition of water



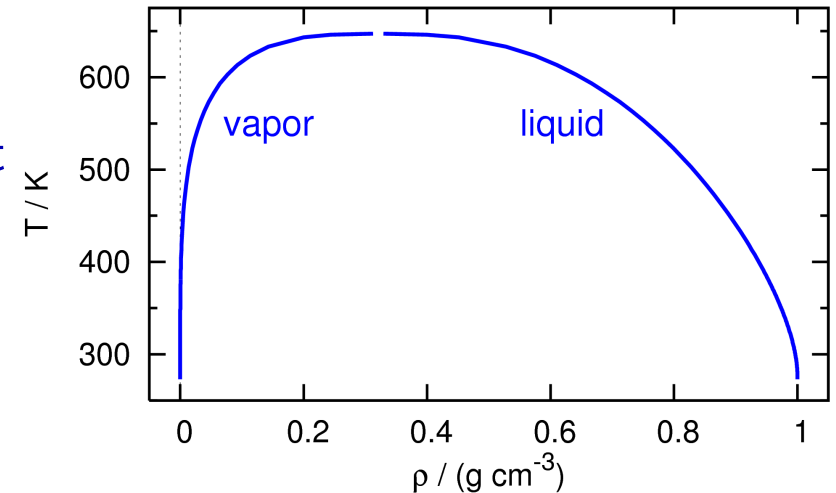
liquid-gas transition: 1st order at $T_t < T < T^*$
2nd order at $T = T^*$

Phase transitions and critical exponents

Liquid-gas transition of water



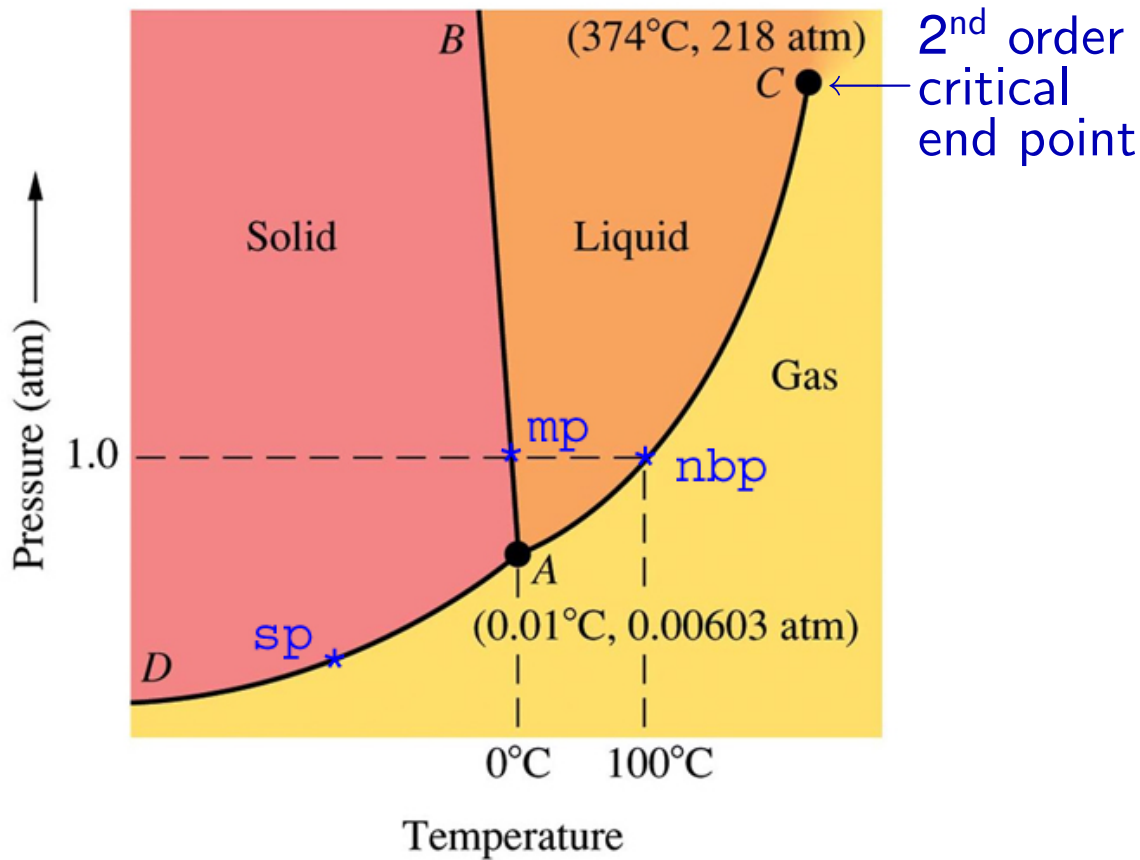
Coexistence curves of vapor + water



liquid-gas transition: 1st order at $T_t < T < T^*$
2nd order at $T = T^*$

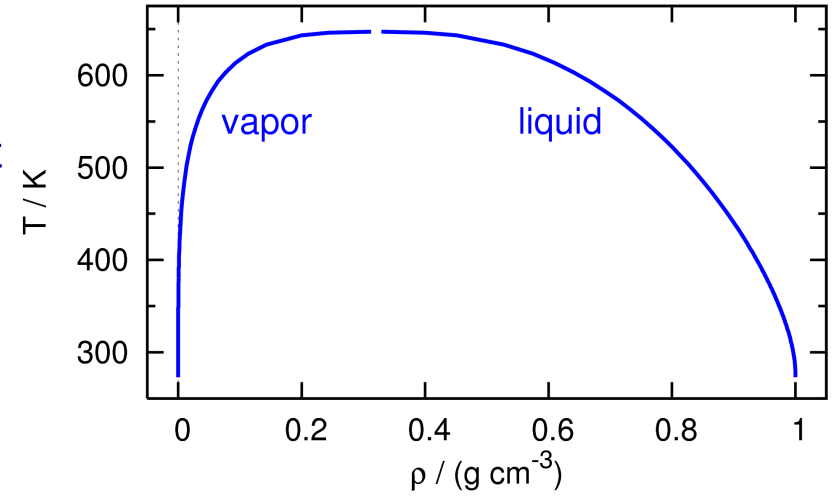
Phase transitions and critical exponents

Liquid-gas transition of water

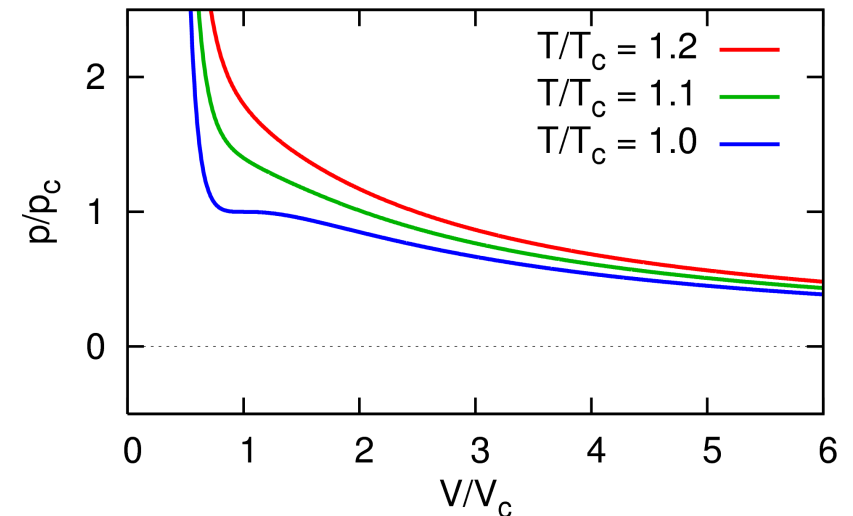


liquid-gas transition: 1st order at $T_t < T < T^*$
 2nd order at $T = T^*$

Coexistence curves of vapor + water

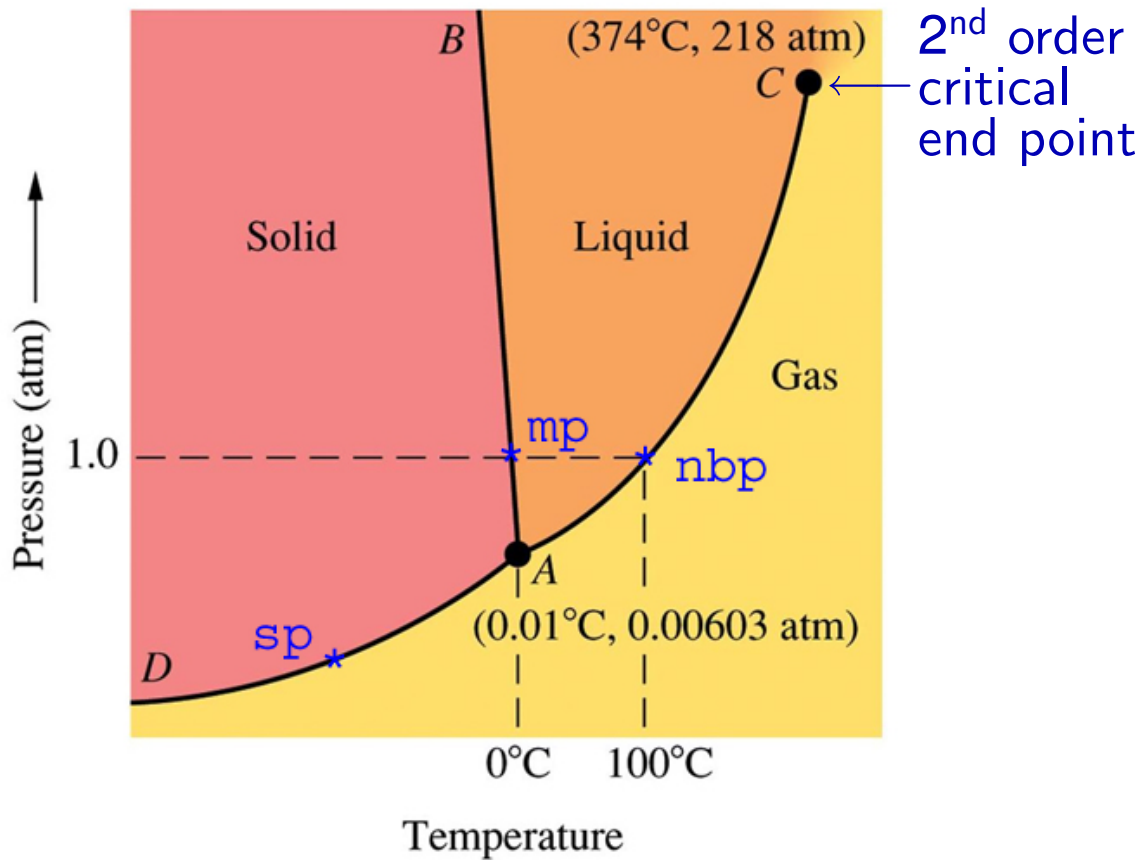


$p(V)$ curves for van der Waals system

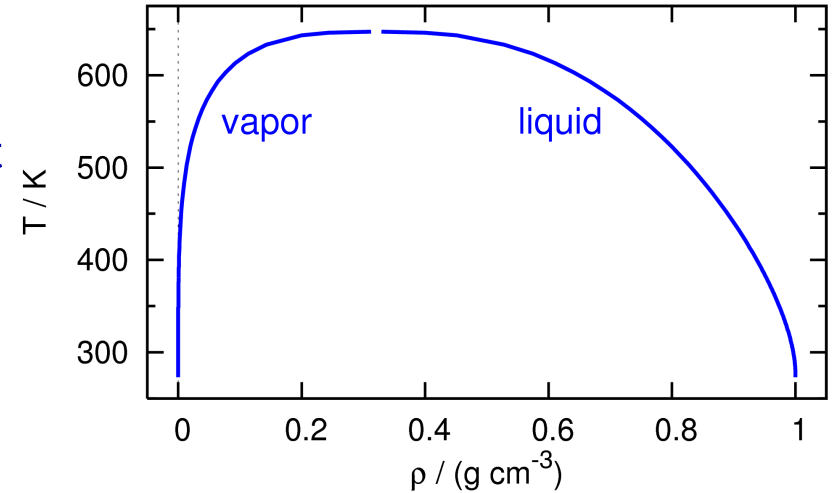


Phase transitions and critical exponents

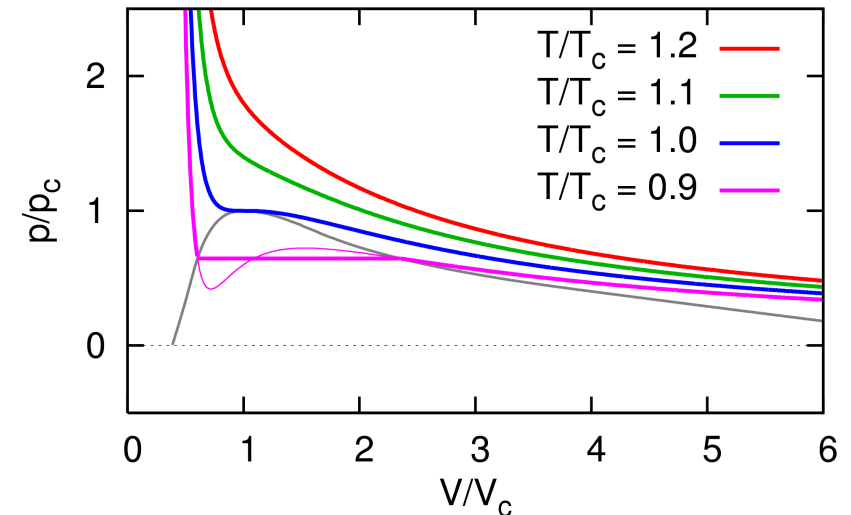
Liquid-gas transition of water



Coexistence curves of vapor + water



$p(V)$ curves for van der Waals system

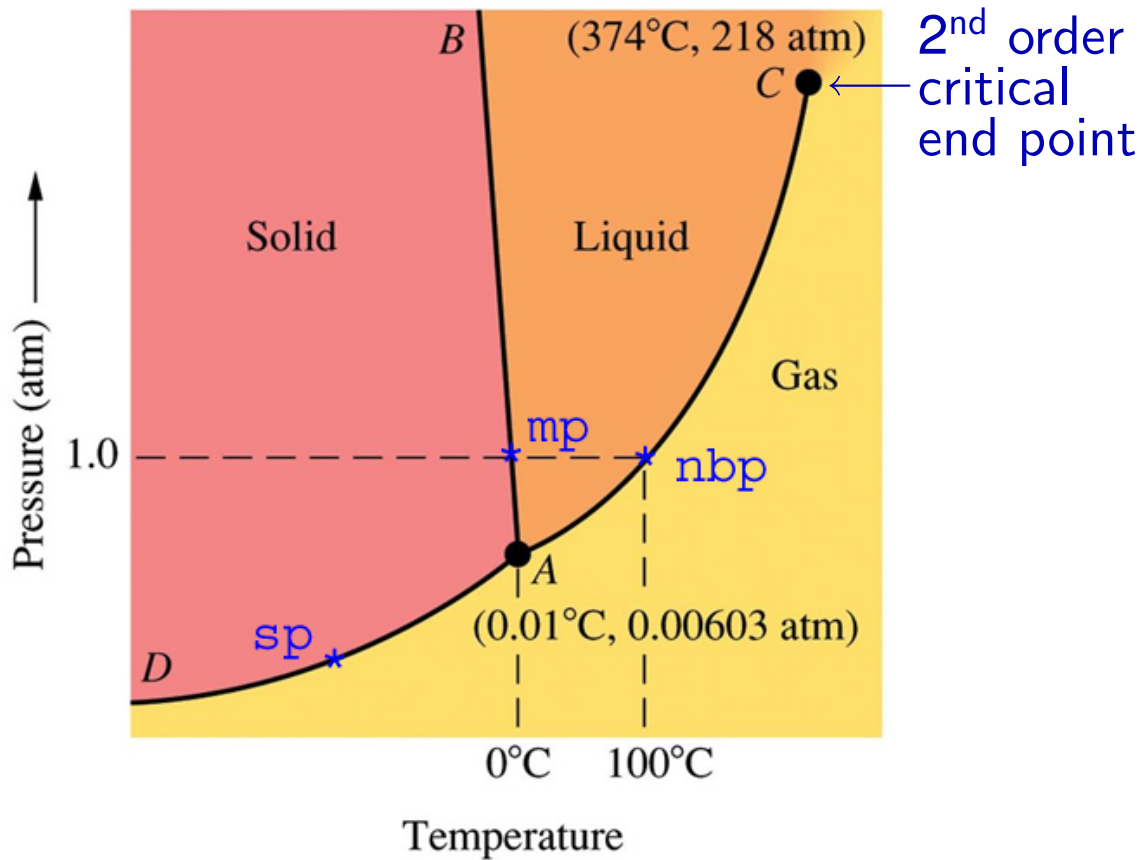


liquid-gas transition: 1st order at $T_t < T < T^*$
 2nd order at $T = T^*$

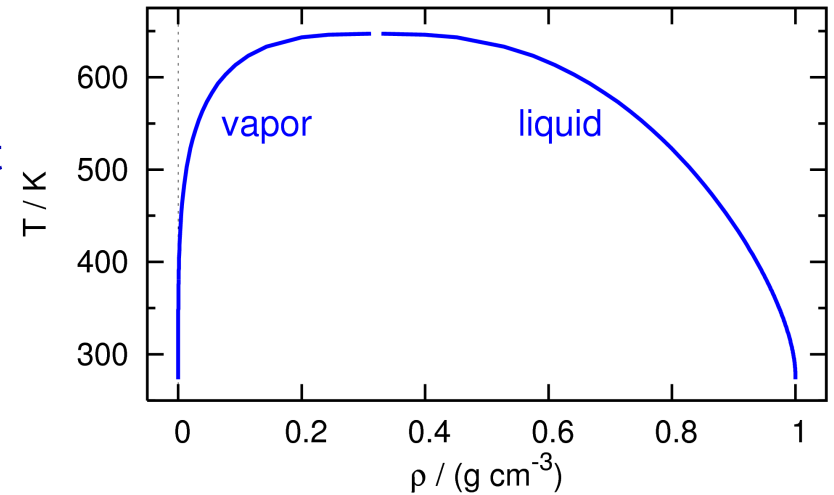
metastable: superheated/cooled water, . . .

Phase transitions and critical exponents

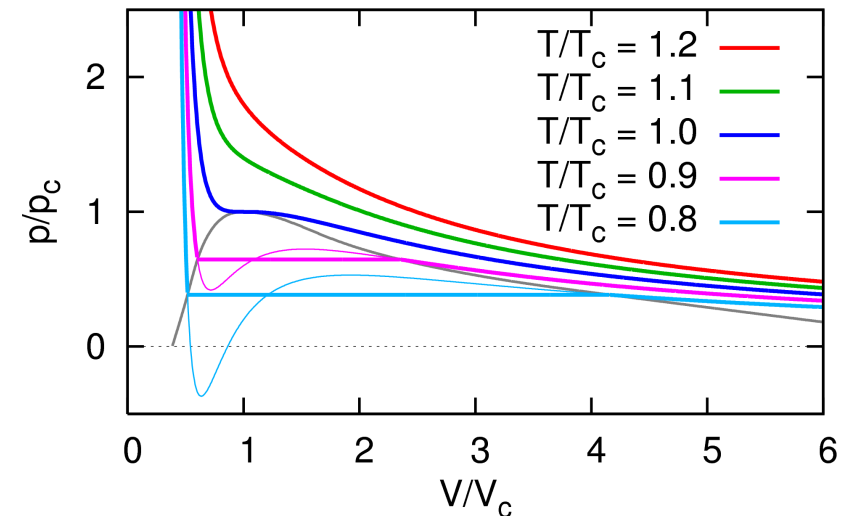
Liquid-gas transition of water



Coexistence curves of vapor + water



$p(V)$ curves for van der Waals system

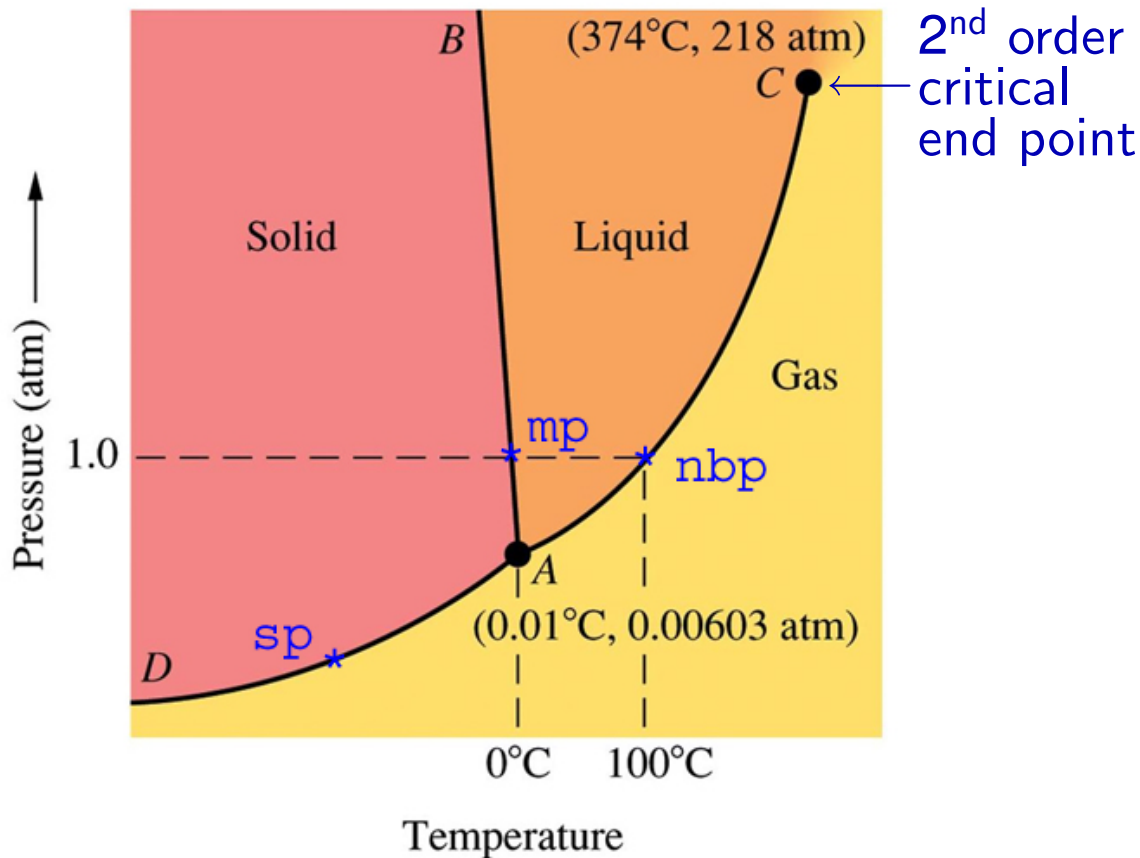


liquid-gas transition: 1st order at $T_t < T < T^*$
 2nd order at $T = T^*$

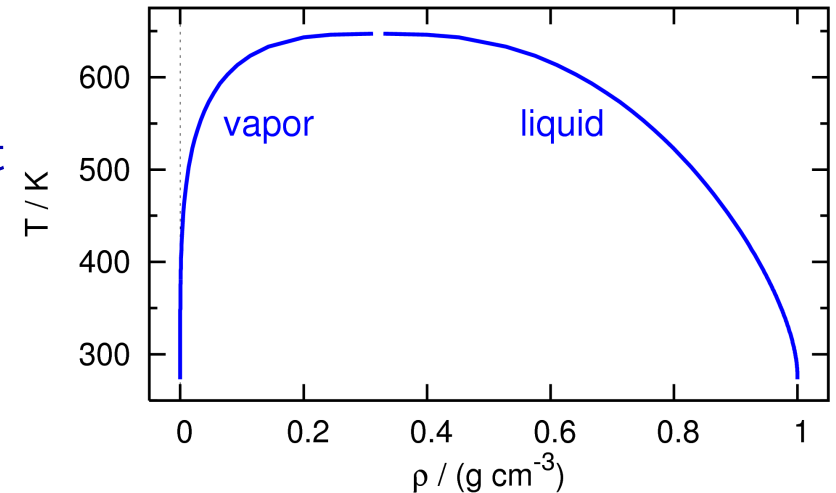
metastable: superheated/cooled water, . . .

Phase transitions and critical exponents

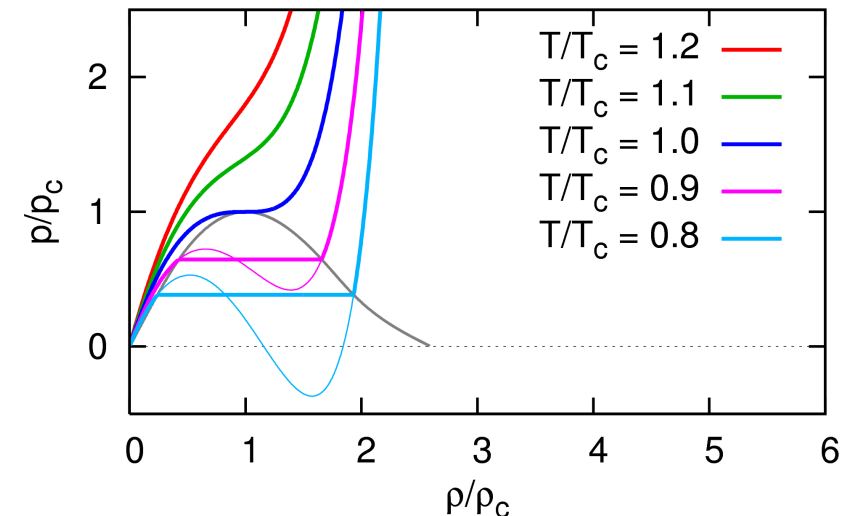
Liquid-gas transition of water



Coexistence curves of vapor + water



$p(V)$ curves for van der Waals system

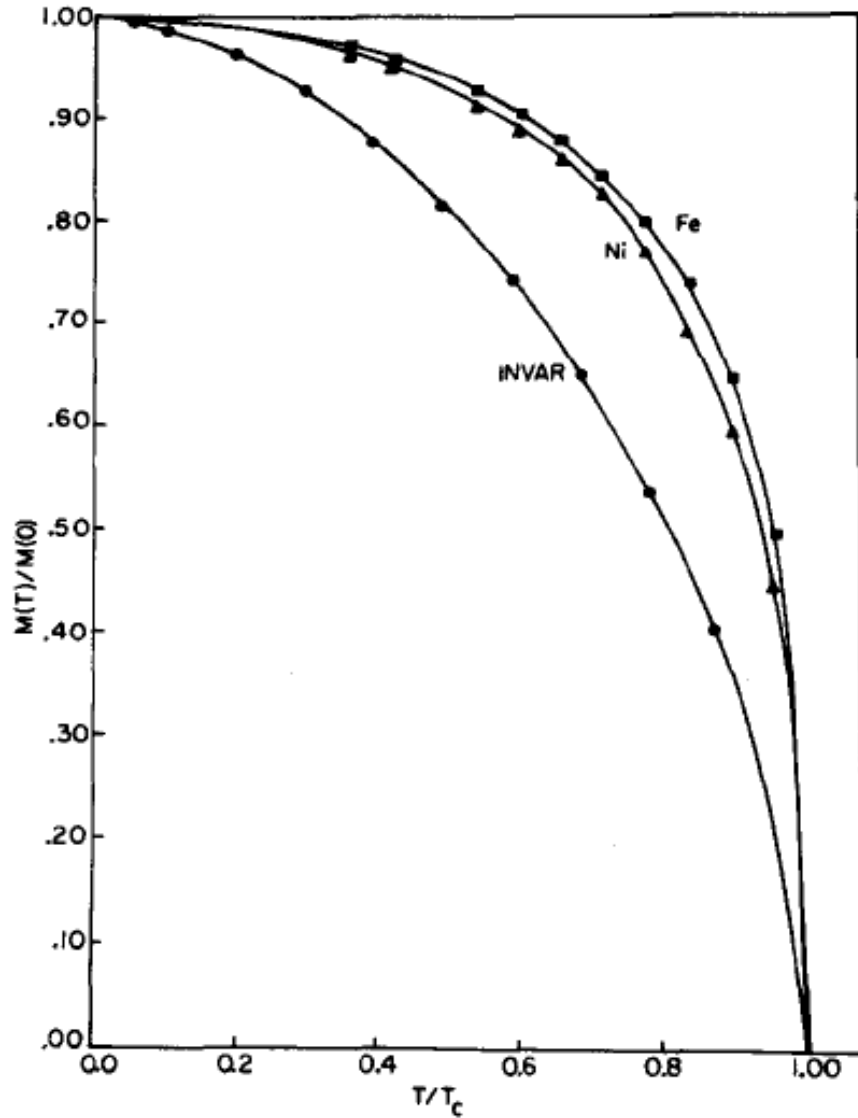


liquid-gas transition: 1st order at $T_t < T < T^*$
 2nd order at $T = T^*$

metastable: superheated/cooled water, . . .

Ferromagnetism

Magnetization of Fe, Ni, INVAR



Critical behavior at 2nd order transition:

$$M(T) \sim |T - T_c|^\beta \quad \text{magnetization}$$

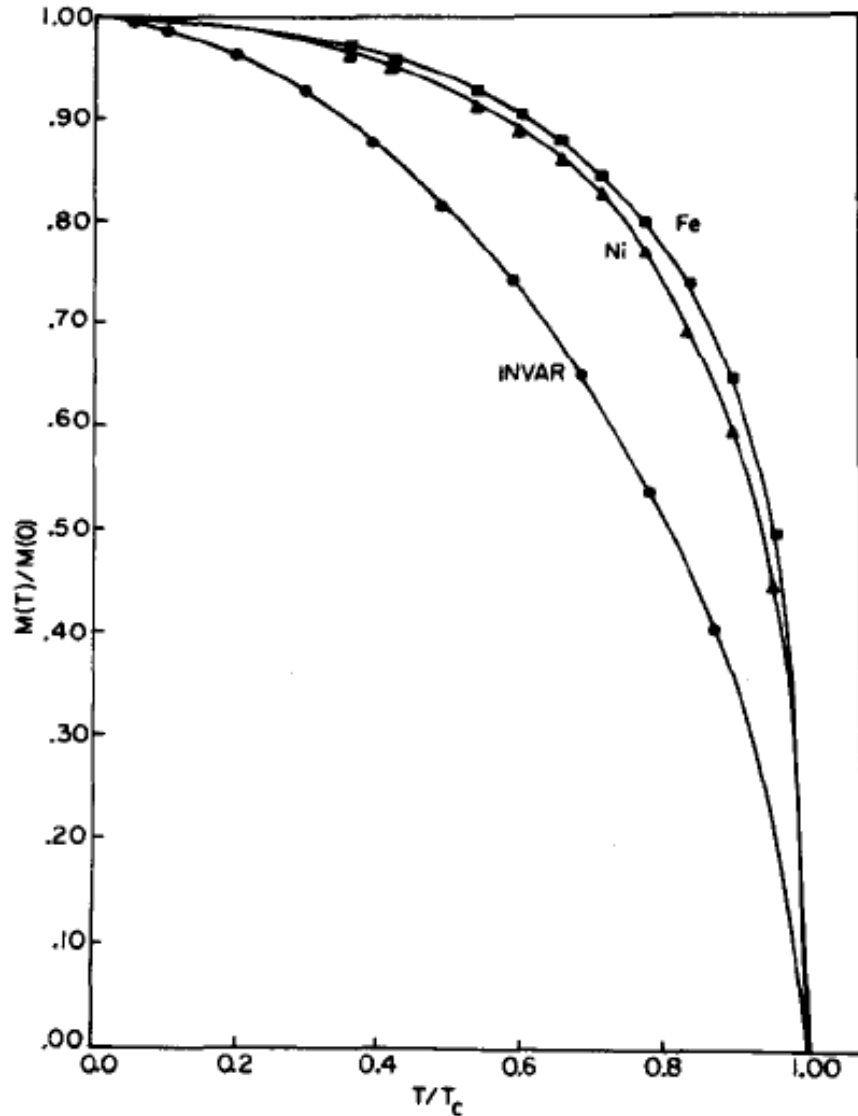
$$\chi(T) \sim |T - T_c|^{-\gamma} \quad \text{susceptibility}$$

$$C(T) \sim |T - T_c|^{-\alpha} \quad \text{specific heat}$$

[Schlosser, Phys. Lett. 40, 195 (1972)]

Ferromagnetism

Magnetization of Fe, Ni, INVAR



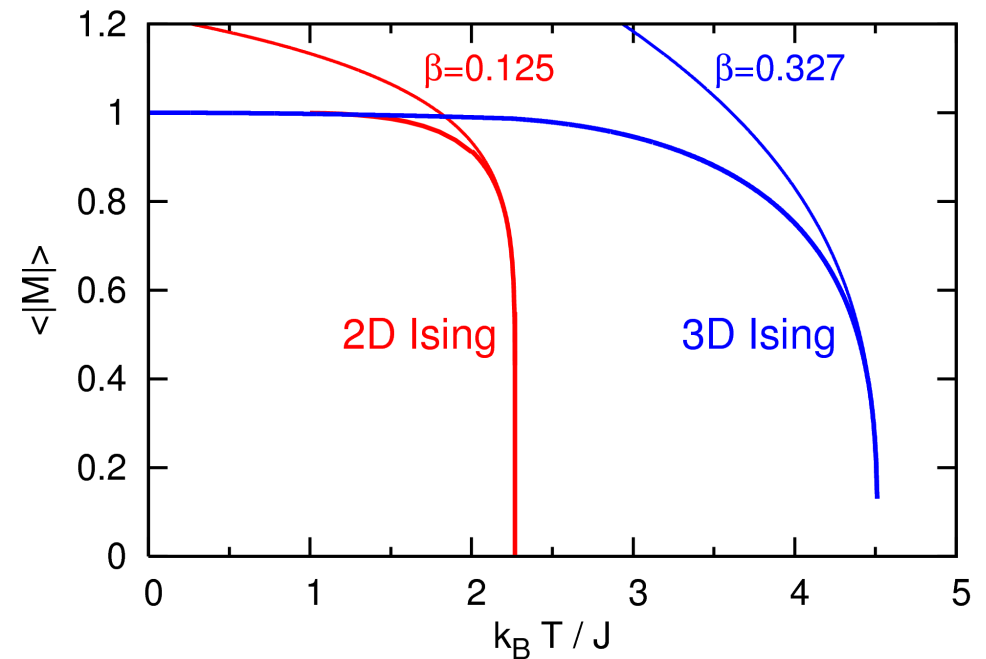
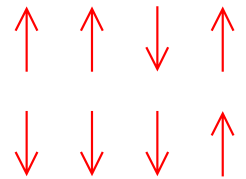
Critical behavior at 2nd order transition:

$$M(T) \sim |T - T_c|^\beta \quad \text{magnetization}$$

$$\chi(T) \sim |T - T_c|^{-\gamma} \quad \text{susceptibility}$$

$$C(T) \sim |T - T_c|^{-\alpha} \quad \text{specific heat}$$

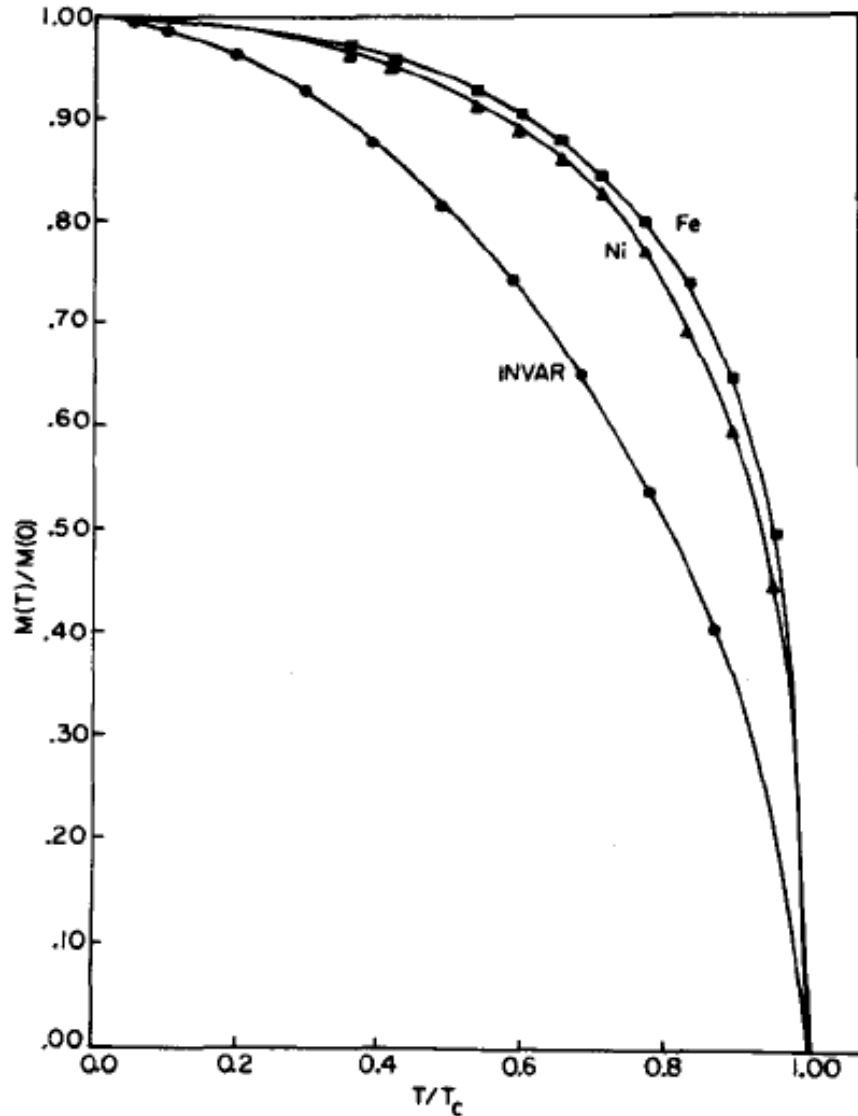
Ising model $H = J \sum_{\langle ij \rangle} s_i^z s_j^z$



[Schlosser, Phys. Lett. 40, 195 (1972)]

Ferromagnetism

Magnetization of Fe, Ni, INVAR



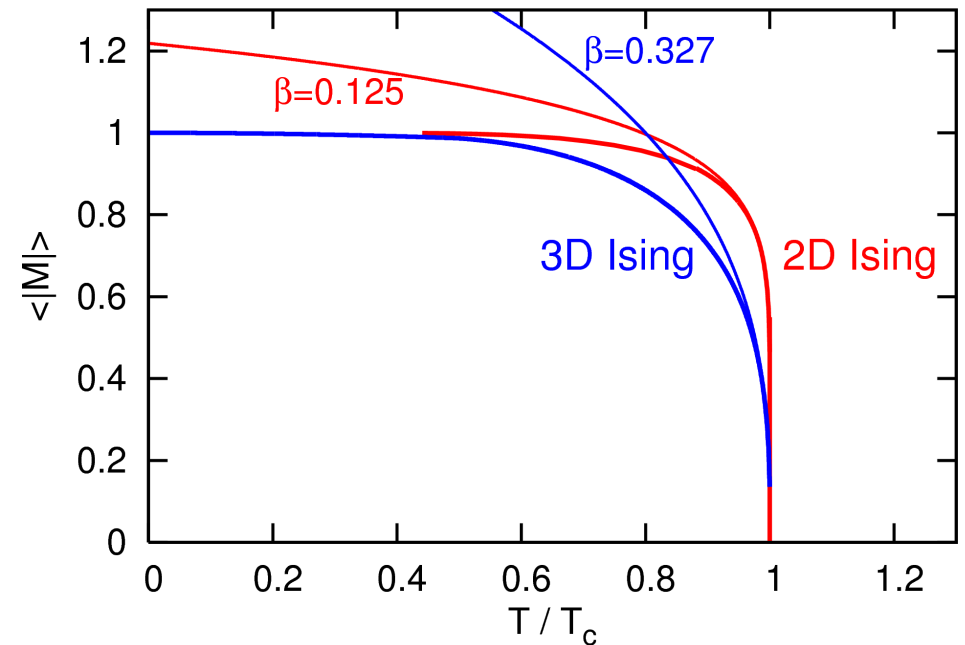
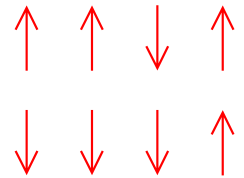
Critical behavior at 2nd order transition:

$$M(T) \sim |T - T_c|^\beta \quad \text{magnetization}$$

$$\chi(T) \sim |T - T_c|^{-\gamma} \quad \text{susceptibility}$$

$$C(T) \sim |T - T_c|^{-\alpha} \quad \text{specific heat}$$

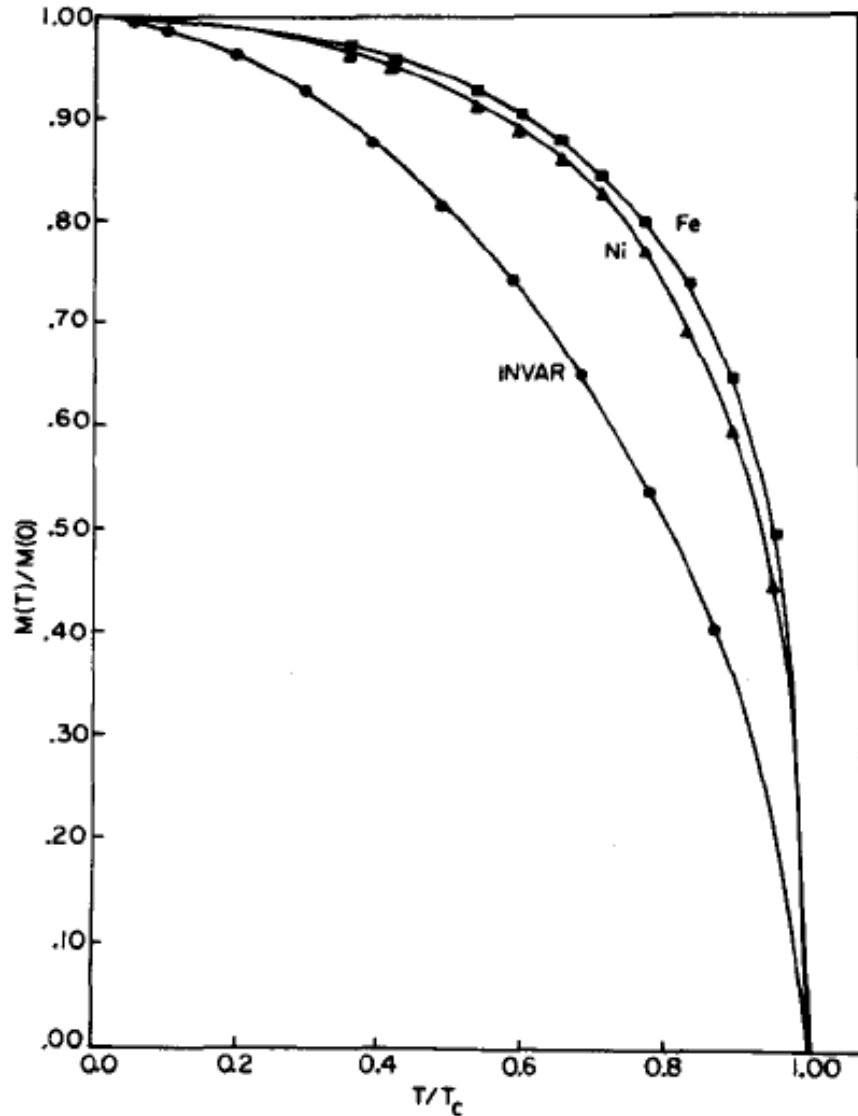
Ising model $H = J \sum_{\langle ij \rangle} s_i^z s_j^z$



[Schlosser, Phys. Lett. 40, 195 (1972)]

Ferromagnetism

Magnetization of Fe, Ni, INVAR



[Schlosser, Phys. Lett. 40, 195 (1972)]

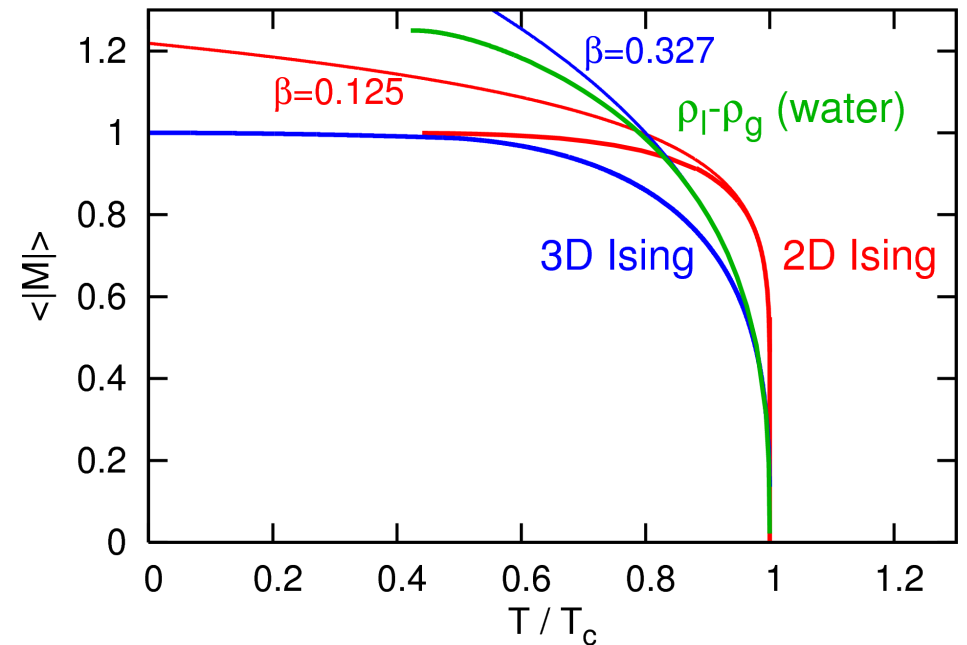
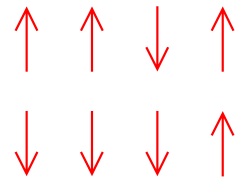
Critical behavior at 2nd order transition:

$$M(T) \sim |T - T_c|^\beta \quad \text{magnetization}$$

$$\chi(T) \sim |T - T_c|^{-\gamma} \quad \text{susceptibility}$$

$$C(T) \sim |T - T_c|^{-\alpha} \quad \text{specific heat}$$

Ising model $H = J \sum_{\langle ij \rangle} s_i^z s_j^z$



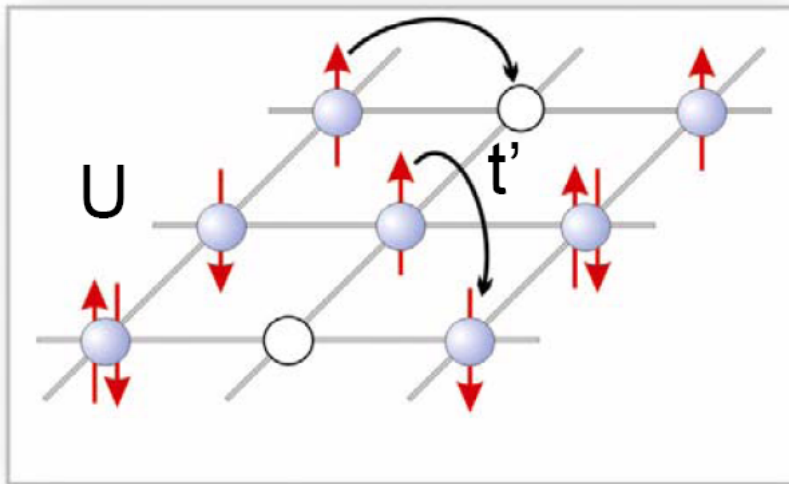
Liquid-gas: 3D Ising universality class

Strongly correlated electron systems

Valence electrons in solids:

correlated by Coulomb repulsion and
Pauli principle \rightsquigarrow Hubbard type models

t



Charge and spin degrees of freedom

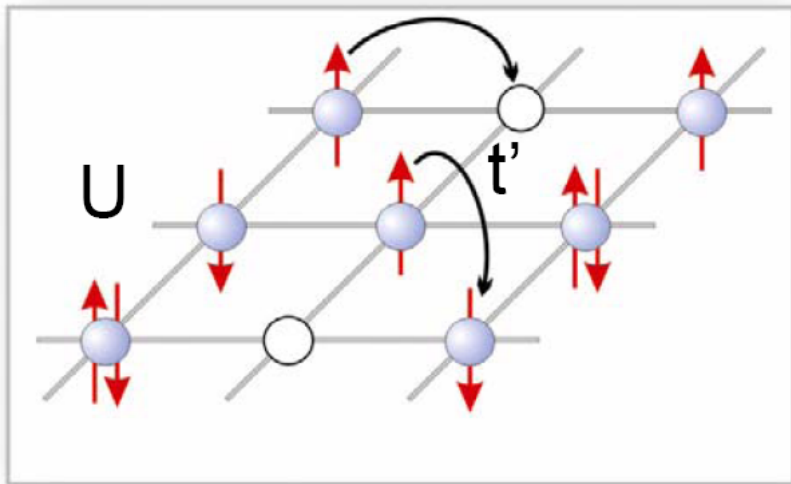
Competition between kinetic energy
($W(t, t')$) and potential energy (U)

Strongly correlated electron systems

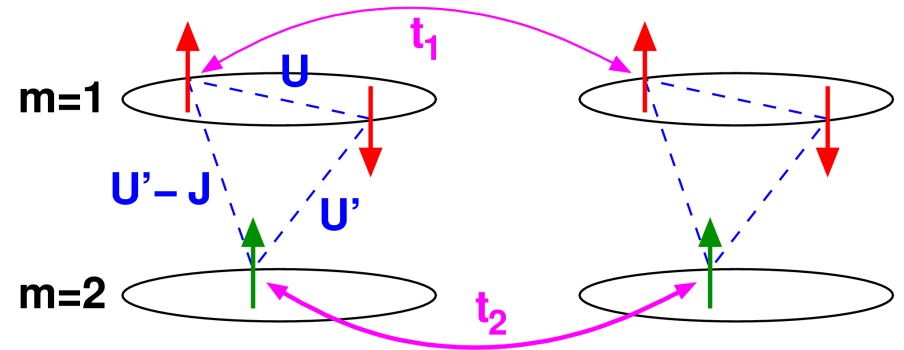
Valence electrons in solids:

correlated by Coulomb repulsion and Pauli principle \rightsquigarrow Hubbard type models

t



Multi-band case: Hund rule couplings

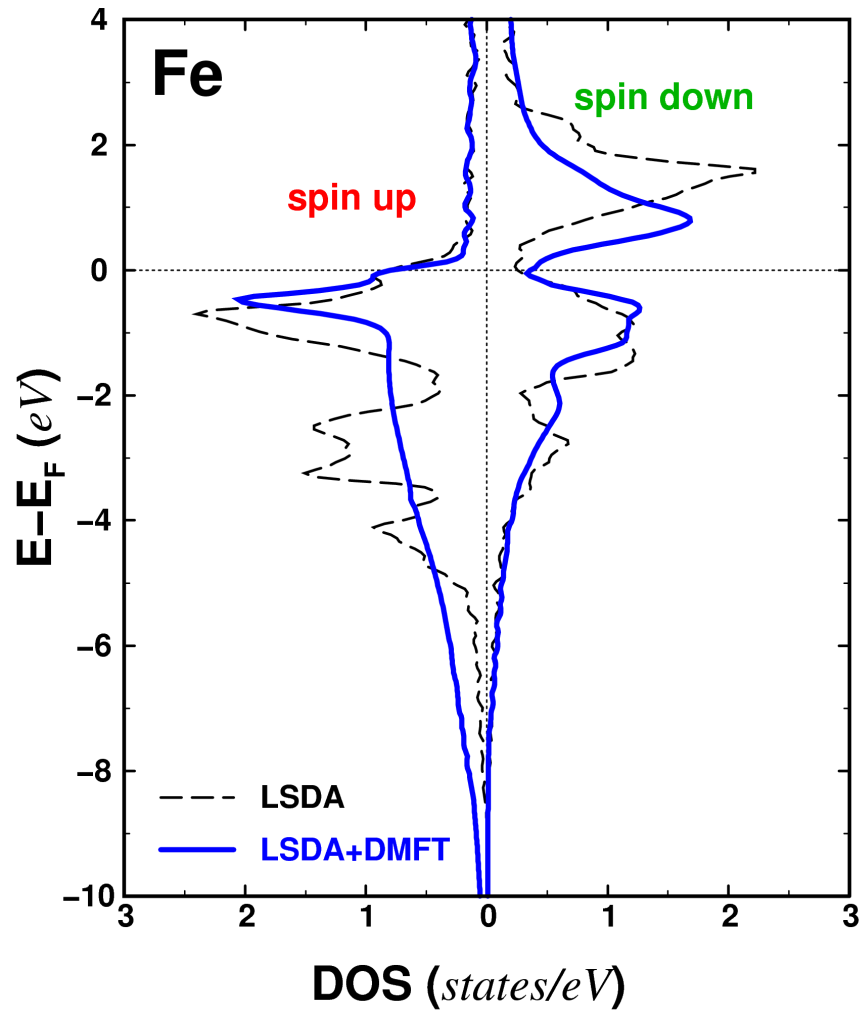


Charge and spin degrees of freedom

Competition between kinetic energy ($W(t, t')$) and potential energy (U)

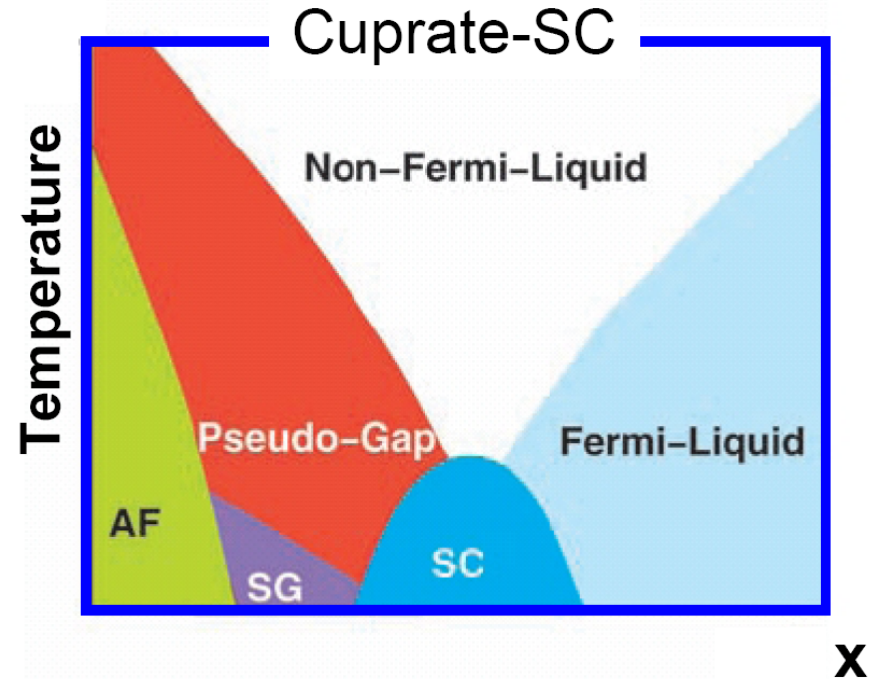
Systems and phenomena

- Itinerant (ferro)magnetism



[Chioncel et. al, PRB (2003)]

- High- T_c superconductivity



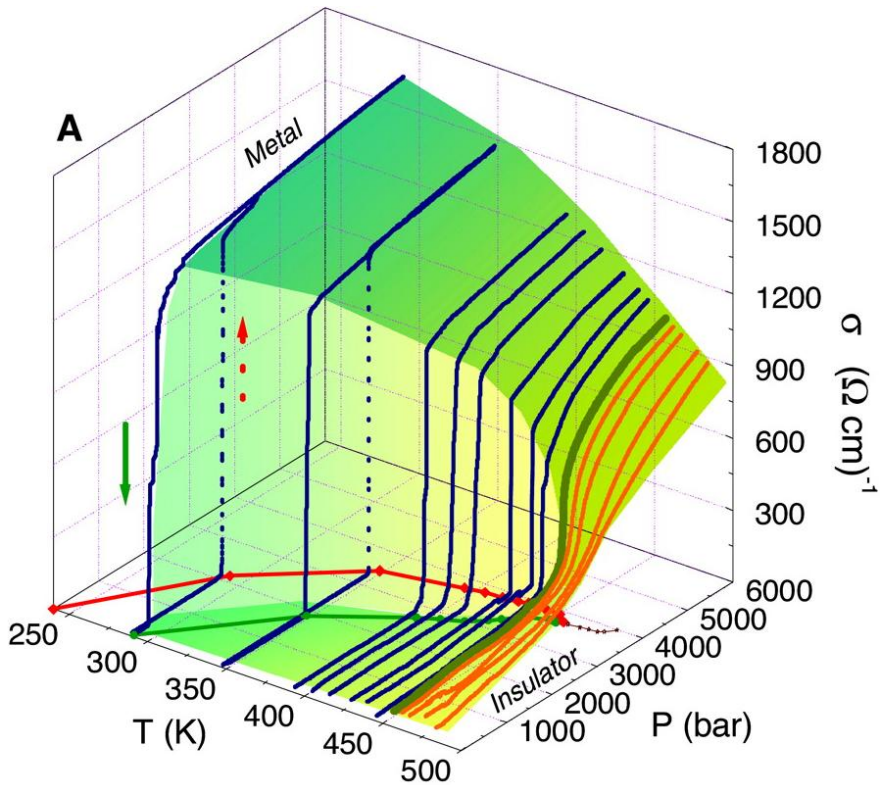
- Metal-insulator transitions
- Ultracold atoms on optical lattices

...

The Mott metal-insulator transition

Mott transition in $(V_{1-x}Cr_x)_2O_3$:

Hysteresis in conductivity

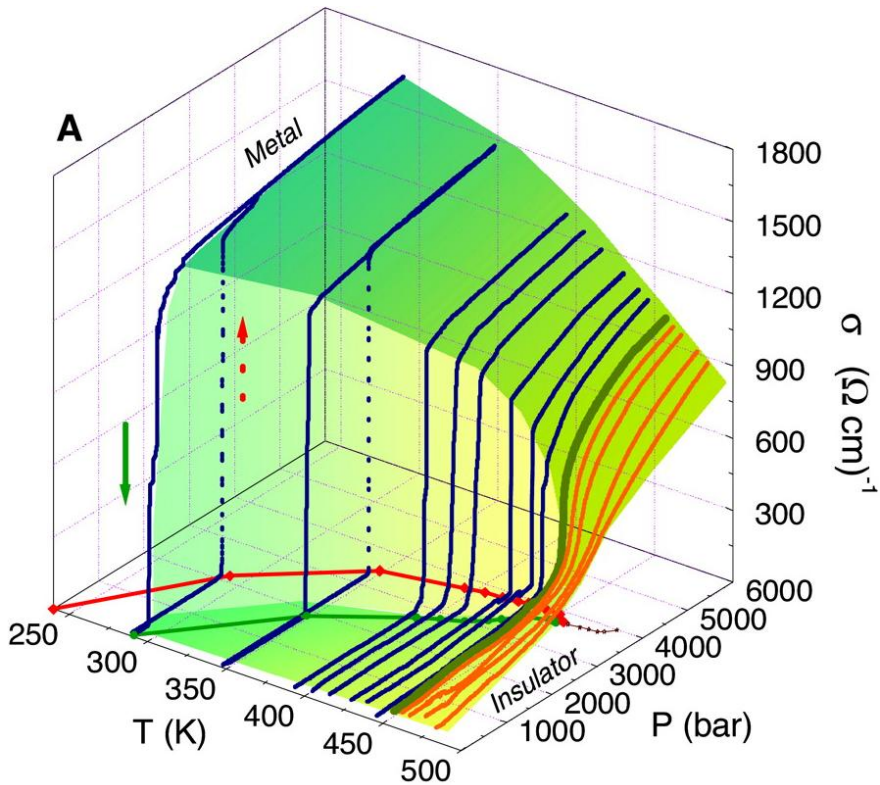


[Limelette et al., Science 302, 89 (2003)]

The Mott metal-insulator transition

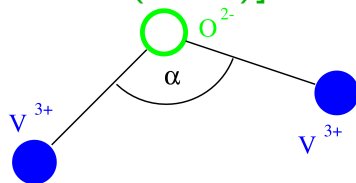
Mott transition in $(V_{1-x}Cr_x)_2O_3$:

Hysteresis in conductivity



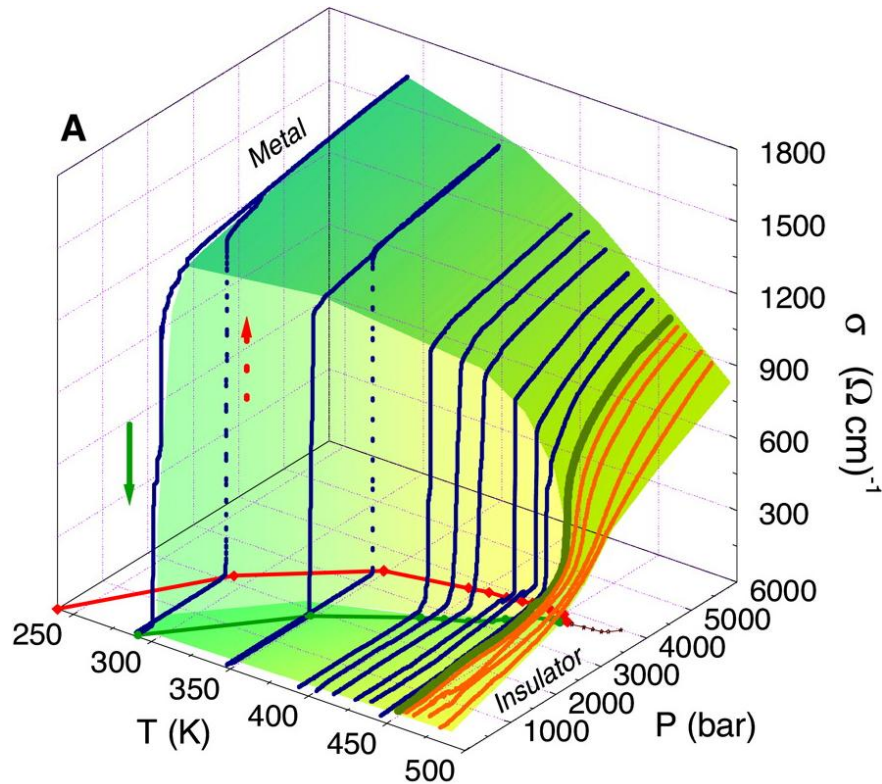
[Limelette et al., Science 302, 89 (2003)]

Pressure
changes W/U



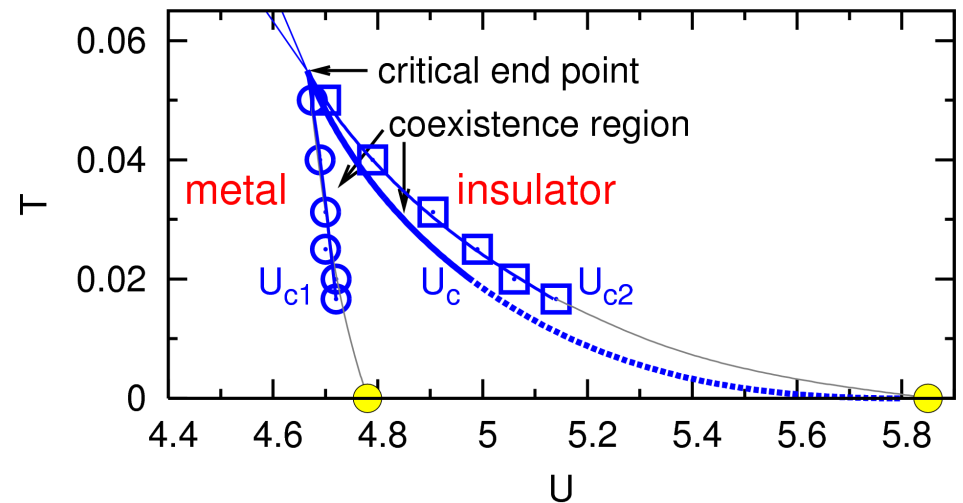
The Mott metal-insulator transition

Mott transition in $(V_{1-x}Cr_x)_2O_3$:
Hysteresis in conductivity



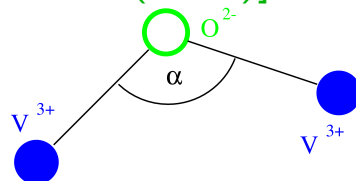
Theoretical phase diagram:

Dynamical mean-field theory (DMFT)
Quantum Monte Carlo (QMC)



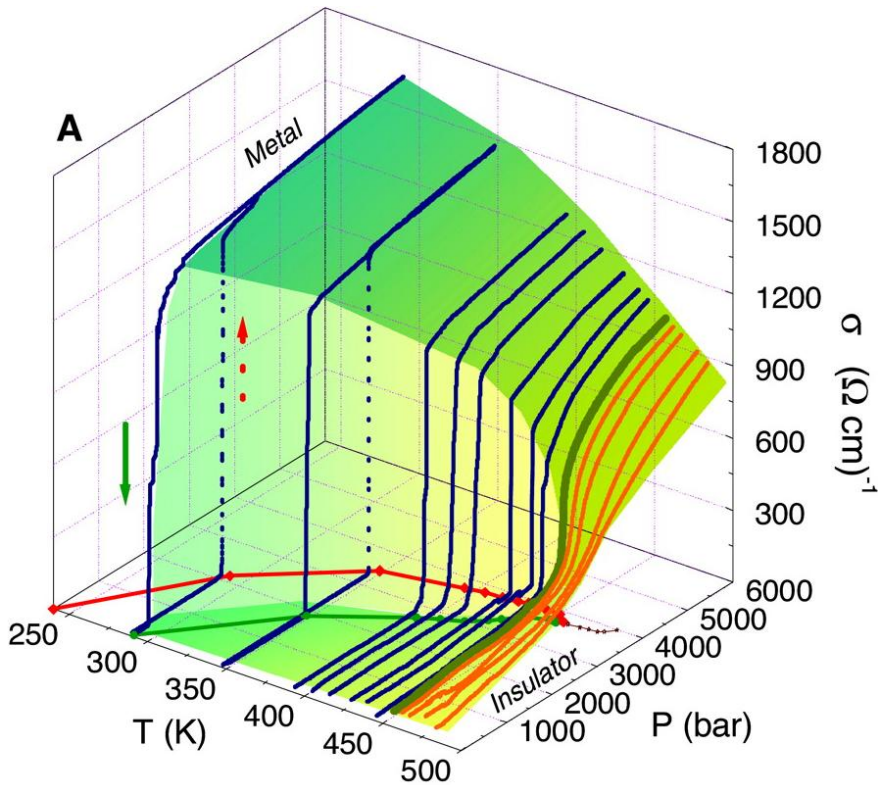
[Limelette et al., Science 302, 89 (2003)]

Pressure
changes W/U



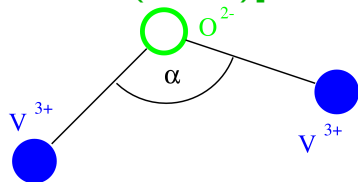
The Mott metal-insulator transition

Mott transition in $(V_{1-x}Cr_x)_2O_3$:
Hysteresis in conductivity



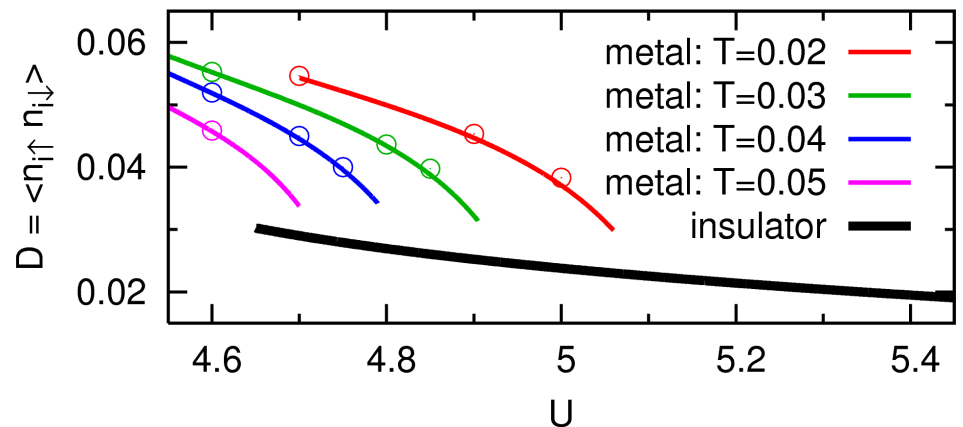
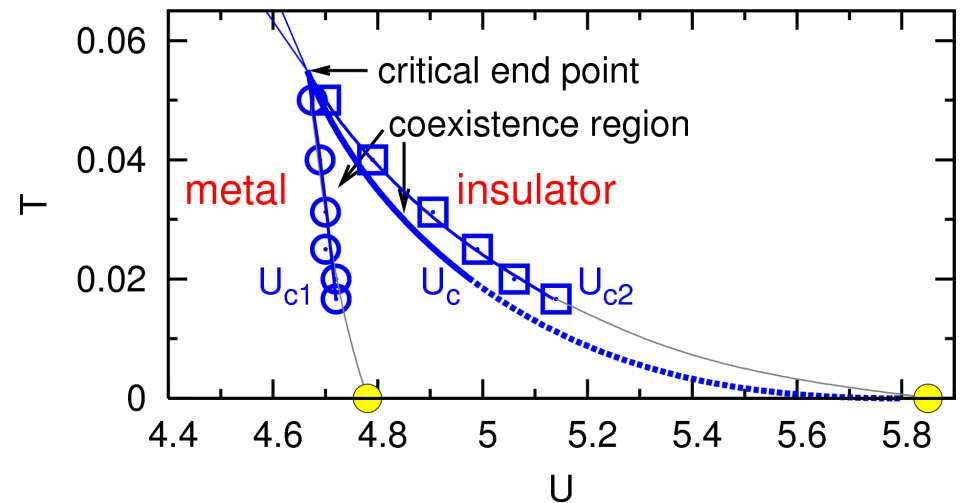
[Limelette et al., Science 302, 89 (2003)]

Pressure
changes W/U



Theoretical phase diagram:

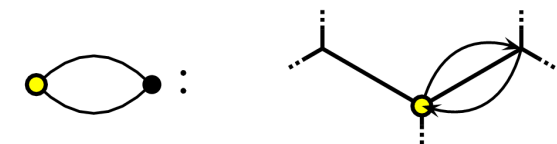
Dynamical mean-field theory (DMFT)
Quantum Monte Carlo (QMC)

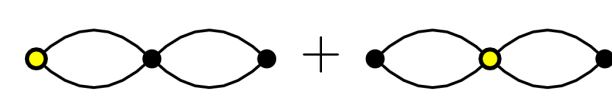


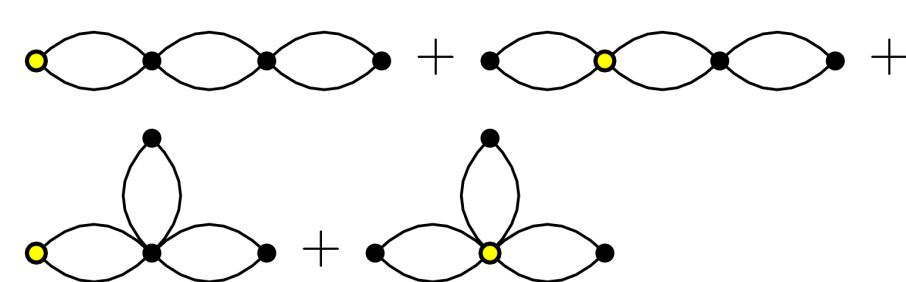
Critical exponents from ePT (using PT + QMC)

Kato-Takahashi strong-coupling perturbation theory ($T=0$)

$$E_{\text{PT}}(U) = -\frac{1}{2} \frac{t^2}{U} - \frac{1}{2} \frac{t^4}{U^3} - \frac{19}{8} \frac{t^6}{U^5} - \frac{593}{32} \frac{t^8}{U^7} - \frac{23877}{128} \frac{t^{10}}{U^9} + \mathcal{O}\left(\frac{t^{12}}{U^{11}}\right)$$

$$\mathcal{O}\left(\frac{t^2}{U}\right):$$


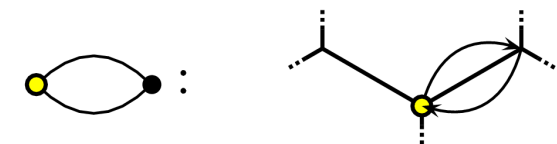
$$\mathcal{O}\left(\frac{t^4}{U^3}\right):$$


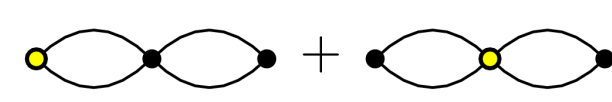
$$\mathcal{O}\left(\frac{t^6}{U^5}\right):$$


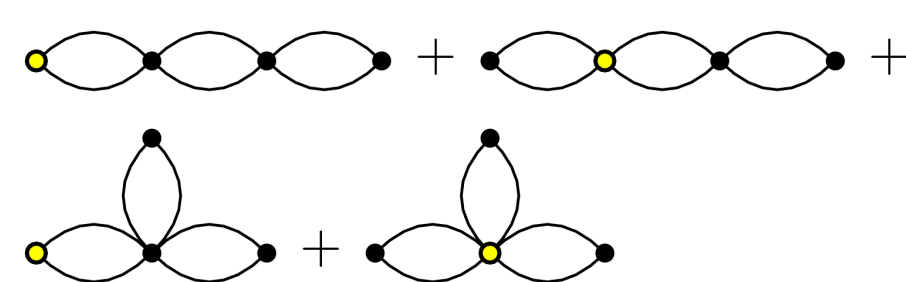
Critical exponents from ePT (using PT + QMC)

Kato-Takahashi strong-coupling perturbation theory ($T=0$)

$$E_{\text{PT}}(U) = -\frac{1}{2} \frac{t^2}{U} - \frac{1}{2} \frac{t^4}{U^3} - \frac{19}{8} \frac{t^6}{U^5} - \frac{593}{32} \frac{t^8}{U^7} - \frac{23877}{128} \frac{t^{10}}{U^9} + \mathcal{O}\left(\frac{t^{12}}{U^{11}}\right)$$

$$\mathcal{O}\left(\frac{t^2}{U}\right):$$


$$\mathcal{O}\left(\frac{t^4}{U^3}\right):$$


$$\mathcal{O}\left(\frac{t^6}{U^5}\right):$$


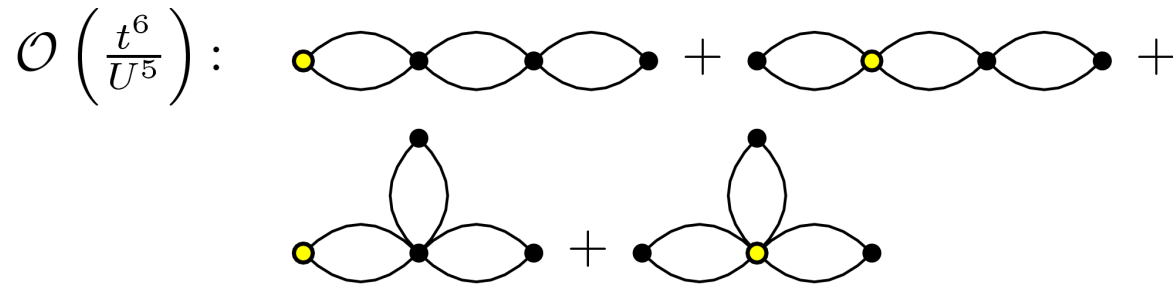
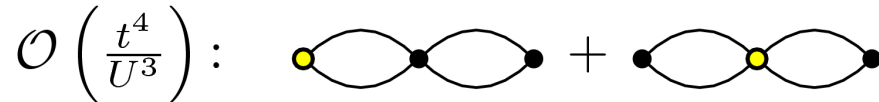
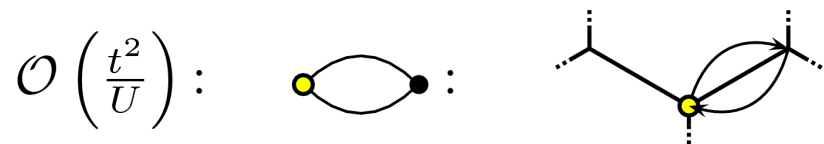
10th order PT accurate (only) at $U \gtrsim 6$: $\Delta E_{\text{PT}} \leq 10^{-5}$

Critical exponents from ePT (using PT + QMC)

Kato-Takahashi strong-coupling perturbation theory ($T=0$)

$$E_{\text{PT}}(U) = -\frac{1}{2} \frac{t^2}{U} - \frac{1}{2} \frac{t^4}{U^3} - \frac{19}{8} \frac{t^6}{U^5} - \frac{593}{32} \frac{t^8}{U^7} - \frac{23877}{128} \frac{t^{10}}{U^9} + \mathcal{O}\left(\frac{t^{12}}{U^{11}}\right)$$

coefficient ratios: 1 4.8 7.8 10.1



10th order PT accurate (only) at $U \gtrsim 6$: $\Delta E_{\text{PT}} \leq 10^{-5}$

extrapolation?

Extrapolated perturbation theory: ePT

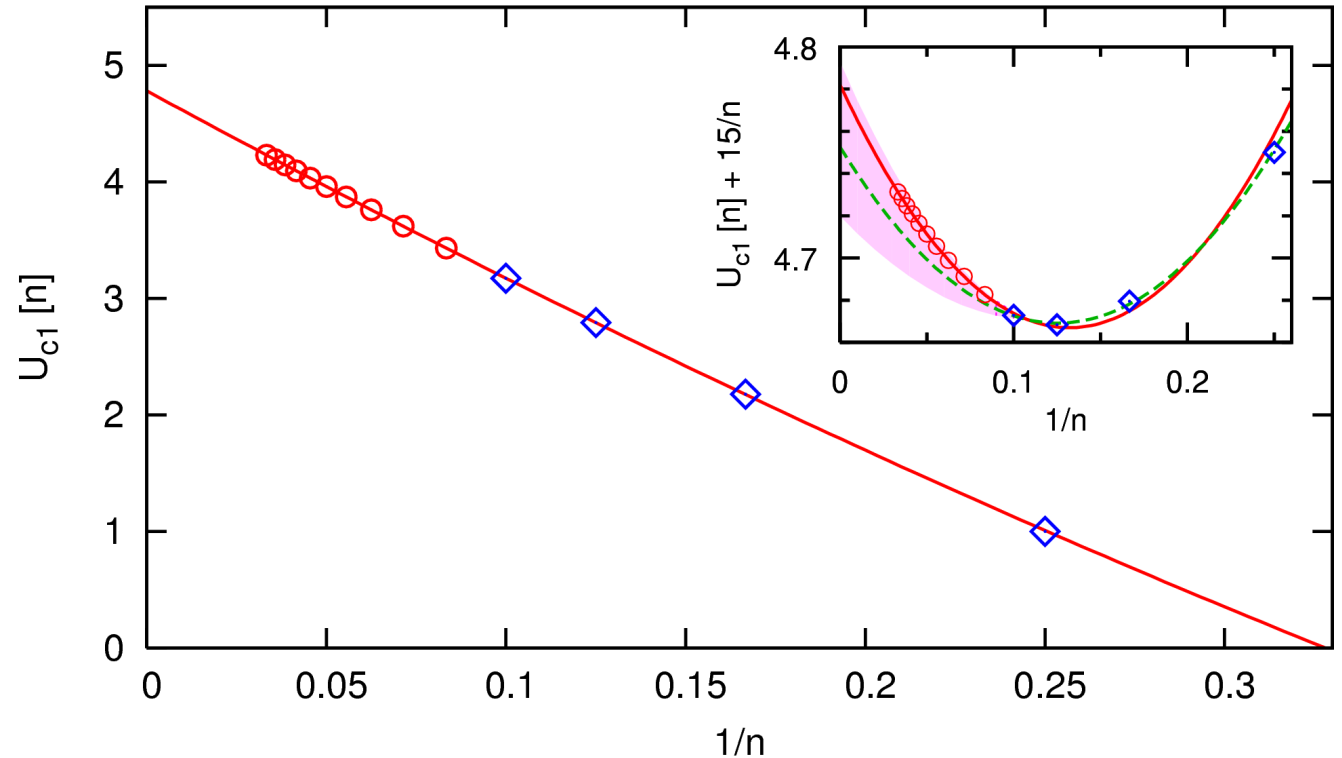
Extrapolate coefficients

$$E_{\text{PT}} = \sum_{i=1}^{\infty} a_{2i} U^{1-2i}$$

by fitting ratios

$$U_{c1}[2i] \equiv \sqrt{a_{2i+2}/a_{2i}} \text{ to}$$

$$U_{c1}[n] \approx U_{c1} + u_1 n^{-1} + u_2 n^{-2}$$



Extrapolated perturbation theory: ePT

Extrapolate coefficients

$$E_{\text{PT}} = \sum_{i=1}^{\infty} a_{2i} U^{1-2i}$$

by fitting ratios

$$U_{c1}[2i] \equiv \sqrt{a_{2i+2}/a_{2i}} \quad \text{to}$$

$$U_{c1}[n] \approx U_{c1} + u_1 n^{-1} + u_2 n^{-2}$$

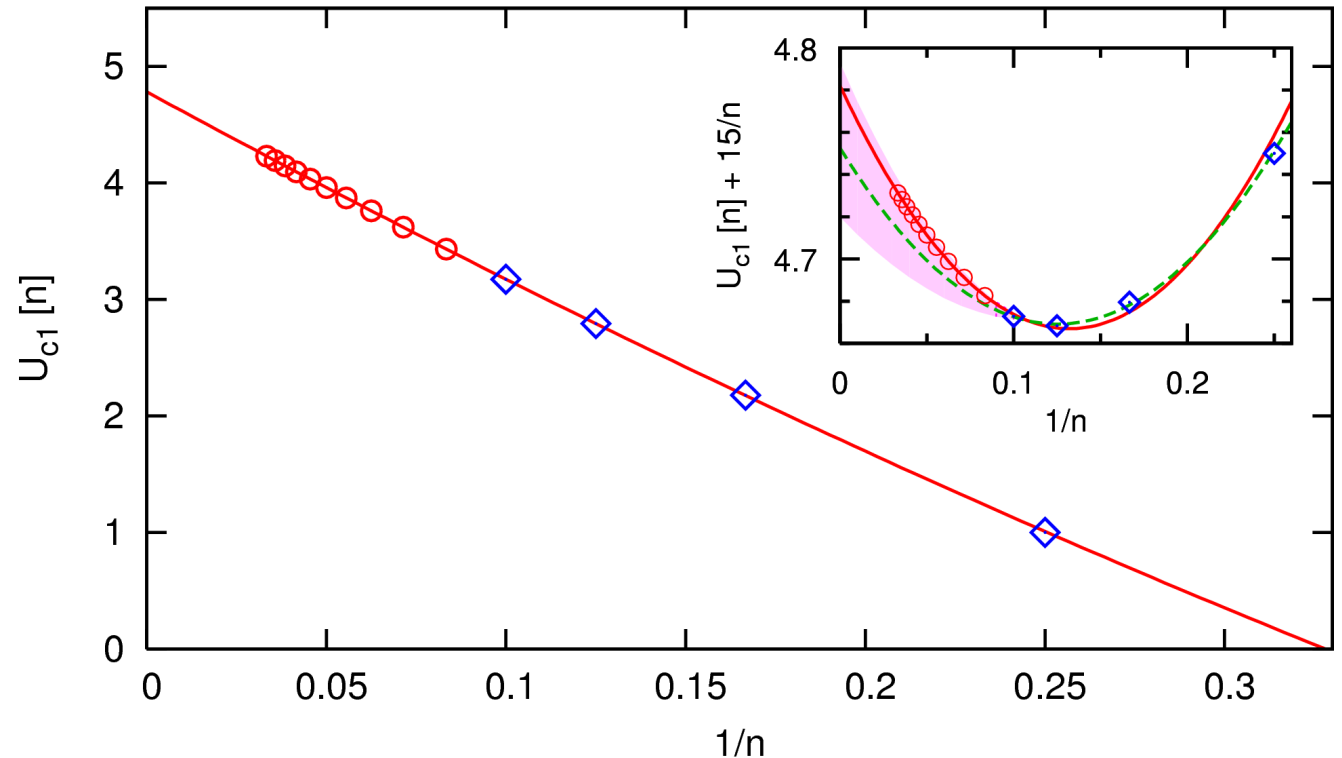
General consequences:

$$U_{c1} = \lim_{i \rightarrow \infty} U_{c1}[2i]$$

$$a_n \propto n^{\tau} U_{c1}^n; \quad \tau = -\frac{u_1}{U_{c1}}$$

$$E^{\text{irr}}(U) \propto (U - U_{c1})^{\tau-1}$$

$$D^{\text{irr}}(U) \propto (U - U_{c1})^{\tau-2}$$



Specifics / numerical results of extrapolation:

Unrestricted quadratic fit $\rightsquigarrow \tau \approx 3.44$, $U_{c1} \approx 4.75$

Extrapolated perturbation theory: ePT

Extrapolate coefficients

$$E_{\text{PT}} = \sum_{i=1}^{\infty} a_{2i} U^{1-2i}$$

by fitting ratios

$$U_{c1}[2i] \equiv \sqrt{a_{2i+2}/a_{2i}} \quad \text{to}$$

$$U_{c1}[n] \approx U_{c1} + u_1 n^{-1} + u_2 n^{-2}$$

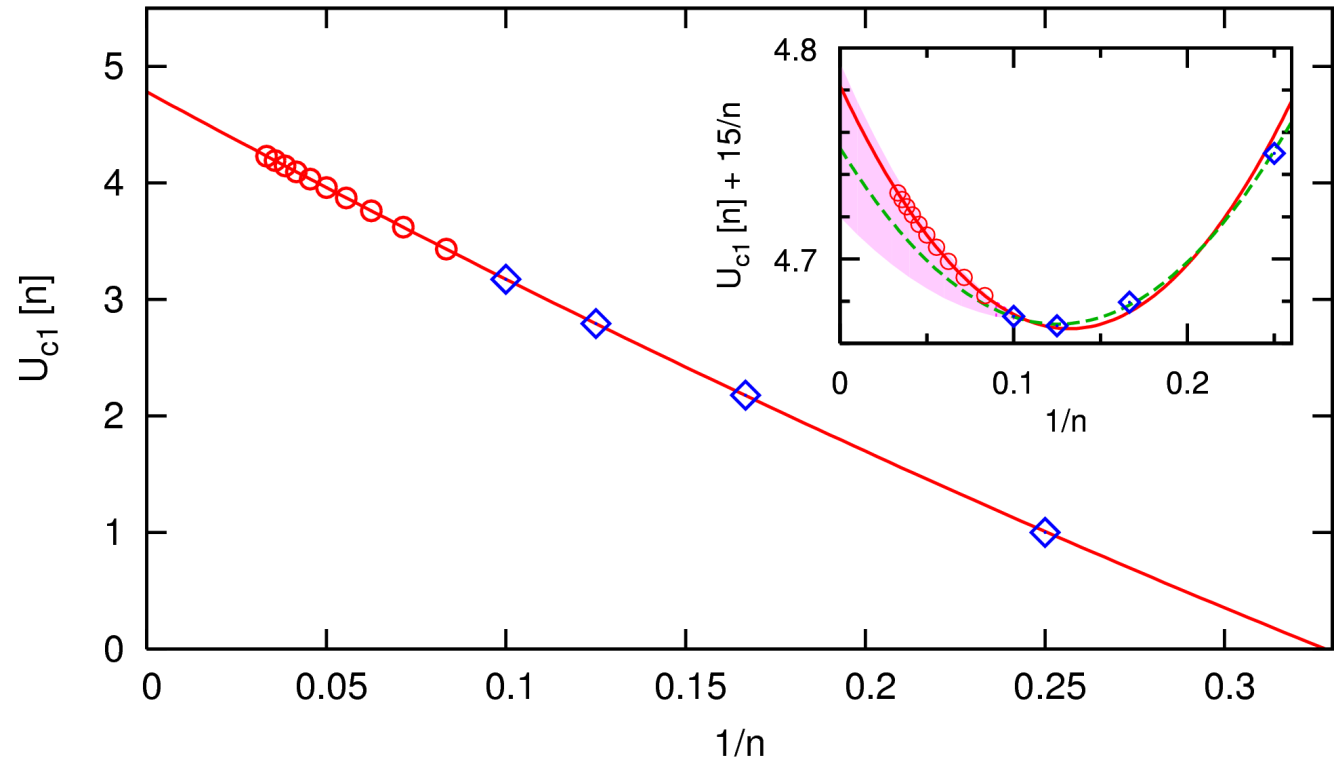
General consequences:

$$U_{c1} = \lim_{i \rightarrow \infty} U_{c1}[2i]$$

$$a_n \propto n^{\tau} U_{c1}^n; \quad \tau = -\frac{u_1}{U_{c1}}$$

$$E^{\text{irr}}(U) \propto (U - U_{c1})^{\tau-1}$$

$$D^{\text{irr}}(U) \propto (U - U_{c1})^{\tau-2}$$



Specifics / numerical results of extrapolation:

Unrestricted quadratic fit $\rightsquigarrow \tau \approx 3.44$, $U_{c1} \approx 4.75$

Comparisons with QMC $\rightsquigarrow 3.36 \leq \tau \leq 3.53$

Extrapolated perturbation theory: ePT

Extrapolate coefficients

$$E_{\text{PT}} = \sum_{i=1}^{\infty} a_{2i} U^{1-2i}$$

by fitting ratios

$$U_{c1}[2i] \equiv \sqrt{a_{2i+2}/a_{2i}} \quad \text{to}$$

$$U_{c1}[n] \approx U_{c1} + u_1 n^{-1} + u_2 n^{-2}$$

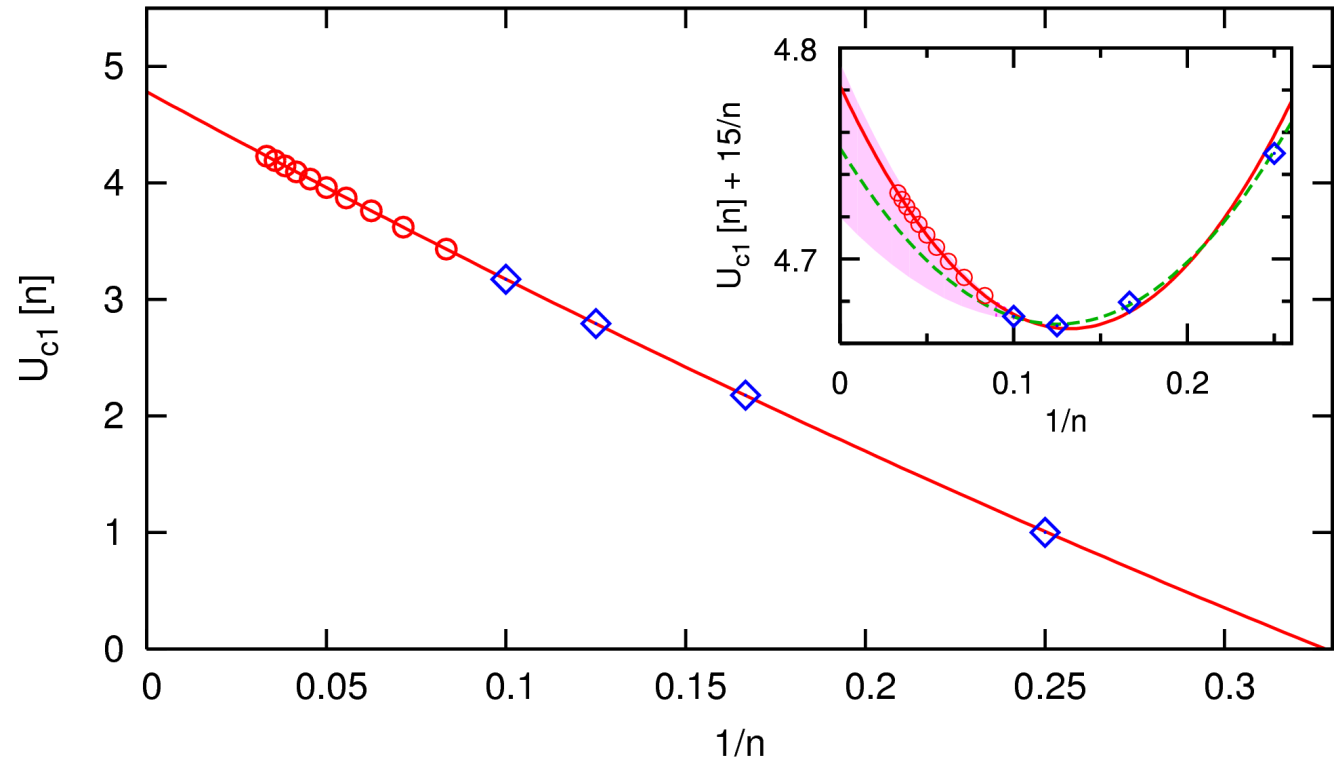
General consequences:

$$U_{c1} = \lim_{i \rightarrow \infty} U_{c1}[2i]$$

$$a_n \propto n^{\tau} U_{c1}^n; \quad \tau = -\frac{u_1}{U_{c1}}$$

$$E^{\text{irr}}(U) \propto (U - U_{c1})^{\tau-1}$$

$$D^{\text{irr}}(U) \propto (U - U_{c1})^{\tau-2}$$



Specifics / numerical results of extrapolation:

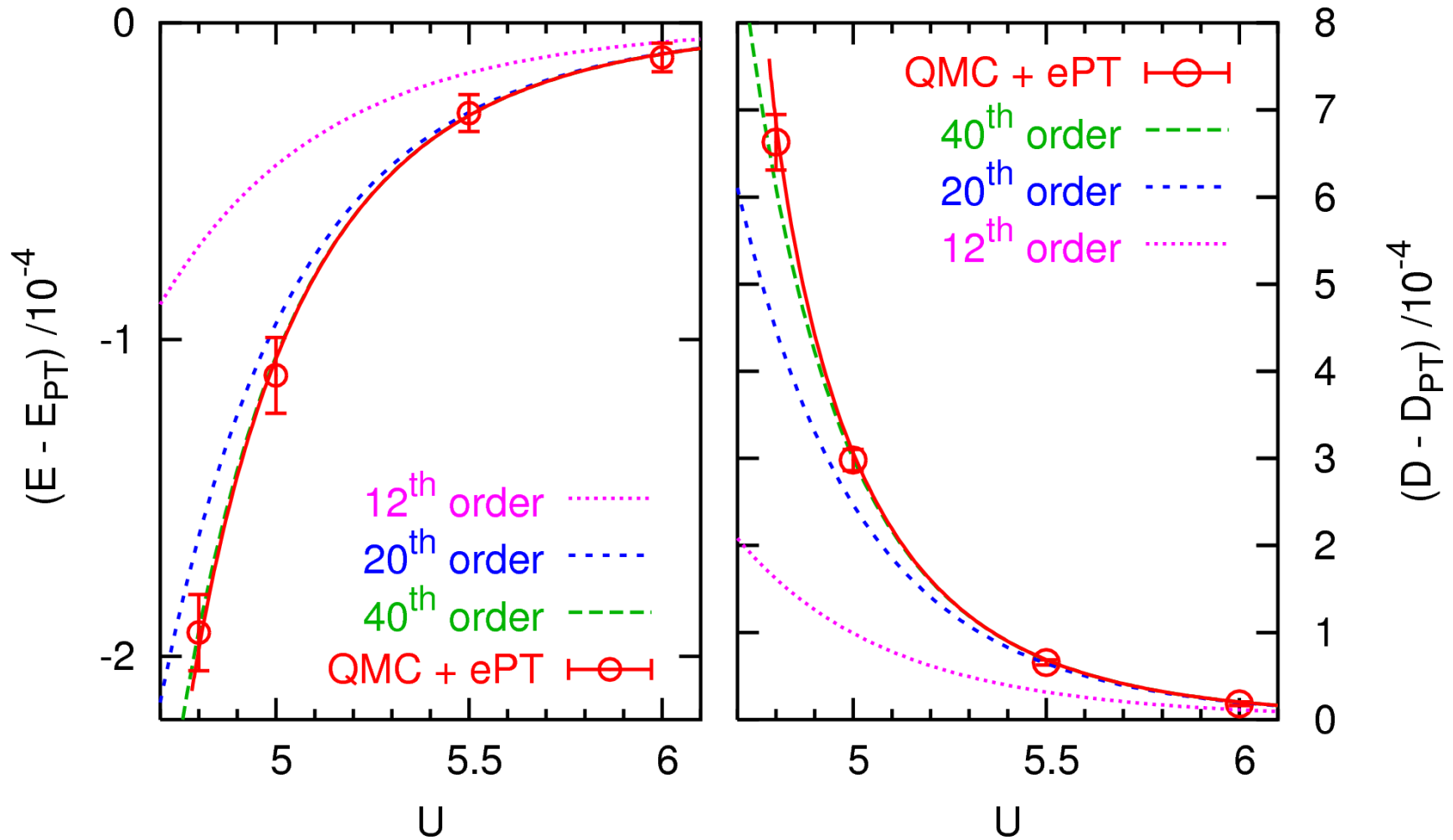
Unrestricted quadratic fit $\rightsquigarrow \tau \approx 3.44$, $U_{c1} \approx 4.75$

Comparisons with QMC $\rightsquigarrow 3.36 \leq \tau \leq 3.53$

Half-integer exponents likely for mean-field theories

Assume $\tau = 3.5 \rightsquigarrow U_{c1} = 4.782$, $E_{\text{ePT}}(U)$, $D_{\text{ePT}}(U)$

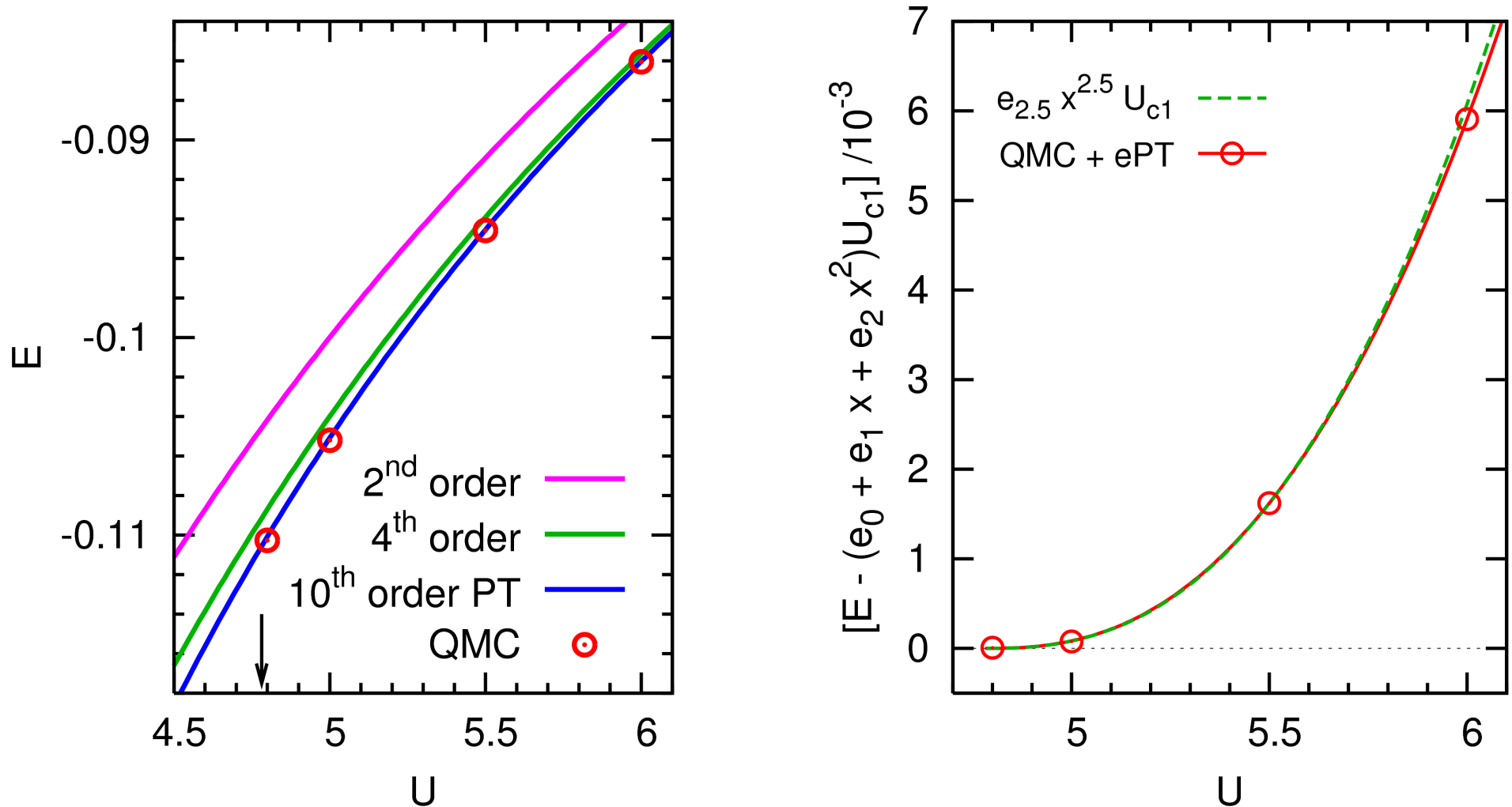
Comparisons: energy E and double occupancy $D = dE/dU$



Excellent agreement \rightsquigarrow reliable exponents, fully parametrized benchmark results
 [Blümer, Kalinowski, Phys. Rev. B **71**, 195102 (2005)]

Prediction of spinodal point U_{c1} confirmed by (most) DMRG, SFA results

A posteriori: coefficients of irregular parts confirmed by QMC



$$E(U) = e_0 + e_1 x + e_2 x^2 + e_{2.5} x^{2.5} + \mathcal{O}(x^3); \quad x \equiv \frac{U - U_{c1}}{U_{c1}}$$

Summary

Critical exponents in strongly correlated electron systems:
within Dynamical mean-field theory (DMFT)
for insulating ground state at spinodal U_{c1}
by extrapolation of strong-coupling PT (+ QMC)

Summary

Critical exponents in strongly correlated electron systems:

within Dynamical mean-field theory (DMFT)

for insulating ground state at spinodal U_{c1}

by extrapolation of strong-coupling PT (+ QMC)

Similar approach: extrapolation of 10th order PT in J_{\perp}/J_{\parallel}
for 2D Heisenberg model [Singh, PRB (1989)]

Summary

Critical exponents in strongly correlated electron systems:

within Dynamical mean-field theory (DMFT)

for insulating ground state at spinodal U_{c1}

by extrapolation of strong-coupling PT (+ QMC)

Similar approach: extrapolation of 10th order PT in J_{\perp}/J_{\parallel}

for 2D Heisenberg model [Singh, PRB (1989)]

Not covered: High-frequency corrections in HF-QMC DMFT solver

Multi-band: orbital-selective Mott transitions

Theory of half-metallic double perovskites (Eberhard Jakobi)

Realistic material-specific calculations with LDA+DMFT

Ultracold quantum gases on optical lattices (SFB-TRR 49)

Summary

Critical exponents in strongly correlated electron systems:

within Dynamical mean-field theory (DMFT)

for insulating ground state at spinodal U_{c1}

by extrapolation of strong-coupling PT (+ QMC)

Similar approach: extrapolation of 10th order PT in J_{\perp}/J_{\parallel}

for 2D Heisenberg model [Singh, PRB (1989)]

Not covered: High-frequency corrections in HF-QMC DMFT solver

Multi-band: orbital-selective Mott transitions

Theory of half-metallic double perovskites (Eberhard Jakobi)

Realistic material-specific calculations with LDA+DMFT

Ultracold quantum gases on optical lattices (SFB-TRR 49)

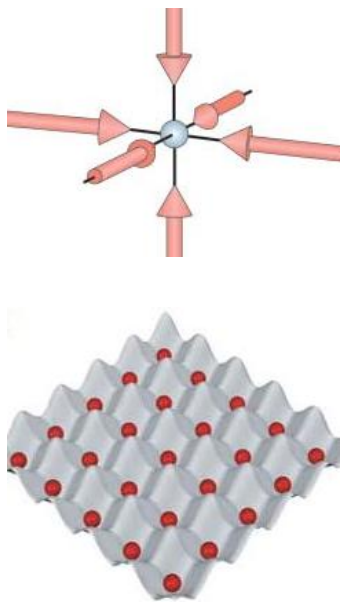
Thanks to: Eva Kalinowski, Carsten Knecht, Peter van Dongen

Florian Gebhard, Wolfgang Paul, Kurt Binder

Starting in 7/2007: SFB/TRR 49 (Frankfurt - Kaiserslautern - Mainz)

Condensed matter systems with variable many-body interactions

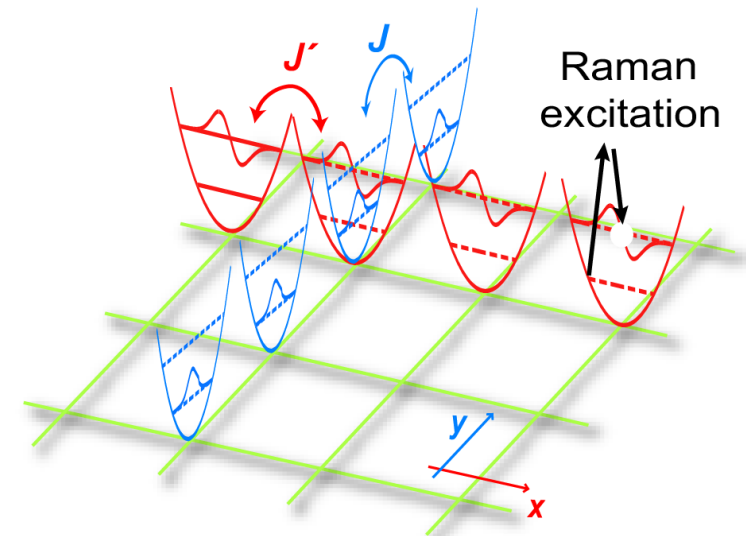
- A1 [Bloch] – Ultracold Fermi mixtures in optical lattices
 - A2 [Kuhr/Bloch] – Spatially addressable quantum gases in optical lattices
 - A3 [Hofstetter] – Inhomogeneous quantum phases in ultracold gases with strong correlations
 - A5 [Fleischhauer/Eggert] – Advanced numerical methods for correlated quantum gases
 - A6 [Blümer] – Flavour-selective Mott transitions of ultracold quantum gases on optical lattices
 - A7 [Hillebrands/Serha] – Collective effects and instabilities of a magnon gas
 - A8 [Kopietz] – Interacting magnons and critical behaviour of bosons
- project area B: real materials



A1 + A6: flavor selectivity in Fermi mixtures of different

- atomic species: ${}^6\text{Li}$ and ${}^{40}\text{K}$
 - hyperfine states
 - vibrational levels
- on optical lattices

Hopping amplitudes tunable and flavor-dependent!



[Müller, Fölling, Widera, Bloch (2007)]