

# Discriminating AF signatures in ultracold fermions using tunable dimensionality and geometric frustration

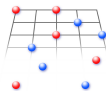
Nils Blümer

Institut für Physik, Johannes Gutenberg-Universität Mainz



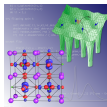
[Gorelik et al., Phys. Rev. A **85**, 061602(R) (2012)]

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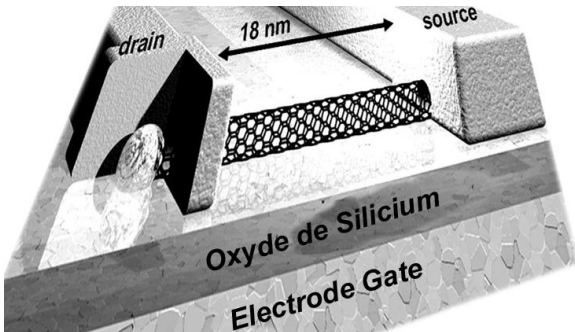
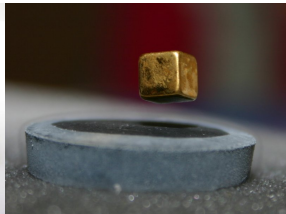
TR 49: *Condensed matter systems with variable many-body interactions*  
Frankfurt / Kaiserslautern / Mainz

FOR 1346  
LDA+DMFT  
Augsburg et al.



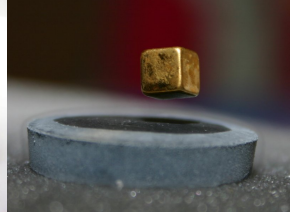
# Motivation: materials science, correlation phenomena

When and why are materials metallic, colored, magnetic, superconducting?

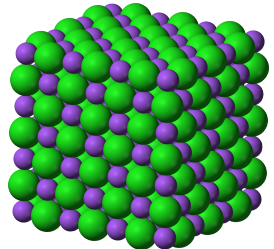
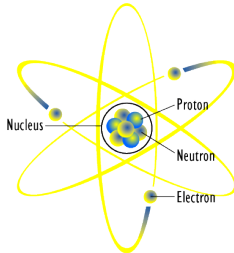


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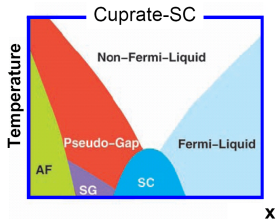


Structure: atoms (nuclei + electrons) in periodic arrangement

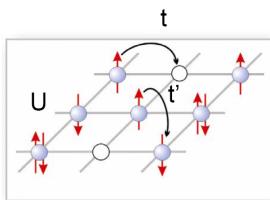


# Motivation: Ultracold lattice fermions as quantum simulators?

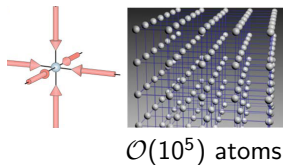
## Correlated materials



## Fermionic Hubbard model

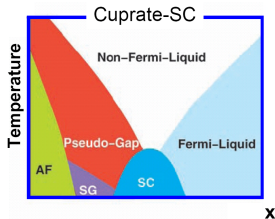


## Ultracold fermions on optical lattices

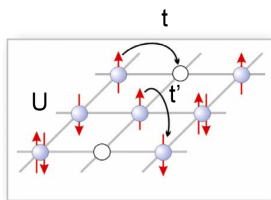


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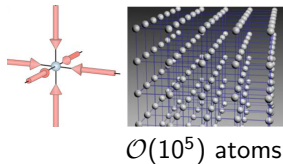
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## Ultracold fermions on optical lattices



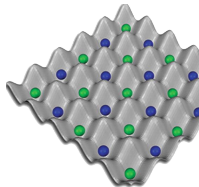
Ultracold fermions on optical lattices can realize Hubbard model (plus X)

- ↪ reproduce known physics ↪ **verify models** for correlated materials
- ↪ **access new regimes** and properties (e.g., time resolution)

## tunable parameters

- ↪ “play” with models, switch features continuously
- ↪ **understanding of phenomena**

Our focus: **antiferromagnetism (AF)**

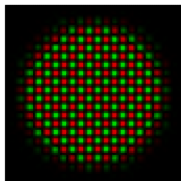


# Specifics of ultracold fermions on optical lattices

- 1) Finite systems ( $\lesssim 10^5$  fermions), trapping potential
  - $\rightsquigarrow$  Numerics expensive, e.g. using RDMFT
  - $\rightsquigarrow$  AF core radius (at best) only 10-20 lattice spacings

[Gorelik et al., PRL **105**, 065301 (2010)]
- 2) Tunable hopping and interactions (+ degeneracy)
- 3) Tunable anisotropy  $\rightsquigarrow$  switch dimensionality

magnetization

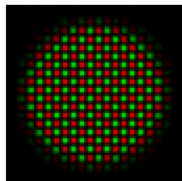


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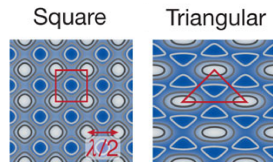
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- 4) Tunable frustration (and dimerization)

- 5) Temperature in trap unknown, but loading  $\sim$  adiabatic

$\rightsquigarrow$  use *entropy* as thermal parameter



Characteristics of AF in finite systems? Impact of dimensionality and frustration?

Dimensionality of (finite) lattice models

Probes of AF correlations and entropy in unfrustrated lattices

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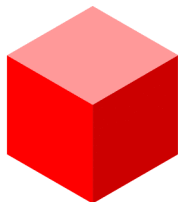
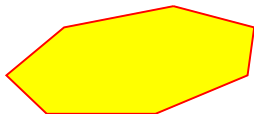
Poster by E. Gorelik:

AF in honeycomb  $\leftrightarrow$  square  $\leftrightarrow$  dimerized lattice

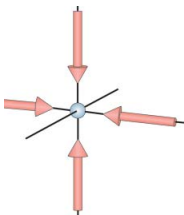
spin correlations near localized impurities (larger  $U$ , reduced  $t$ )

## Dimensionality of finite-size systems: continuous case

- 1) Lines, surfaces, bodies (mathematical objects in  $\mathbb{R}^3$ ):  
dimension defined, e.g., by mapping from  $\mathbb{R}^d (d \leq 3)$



- 2) Continuum systems with finite widths - compare scales



## Dimensionality of finite-size systems: lattice case

General Hubbard model ( $N$  sites,  $t_{ij} = t_{ji}$ )

$$H = \sum_{i,j=1}^N \sum_{\sigma} t_{ij} c_{i\sigma}^{\dagger} c_{j\sigma} + \sum_{i=1}^N U_i n_{i\uparrow} n_{i\downarrow}$$

4 ●

3 ●

5 ●

2 ●

1 ●

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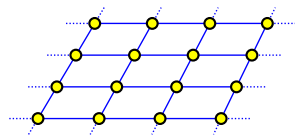
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Familiar case: ● sites form (infinite) **mathematical lattice**,  
with basis vectors  $\tau_{\alpha}$  ( $1 \leq \alpha \leq d$ )

- regular hopping:  $t_{ij} \equiv t_{\tau}$  (where  $\tau = \mathbf{r}_i - \mathbf{r}_j$ )

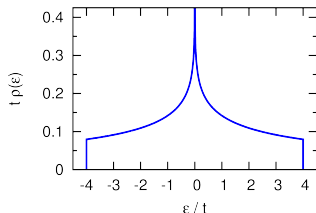
Example: **square lattice**



Clear:  $d = 2$

Important property:  
noninteracting  
**density of states**  
(DOS)  $\rho(\epsilon)$

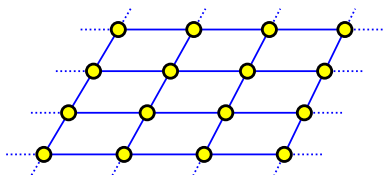
impact of cutoff?



## Dimensionality of finite-size systems: lattice case

Consider  $4 \times 4$  lattice with PBC's

– clearly two-dimensional character!?

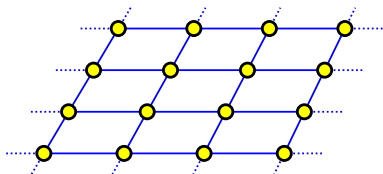
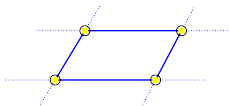


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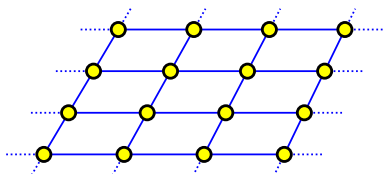
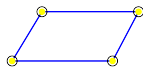
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$\rightsquigarrow d = 1$  ( $\rightsquigarrow$  BA)



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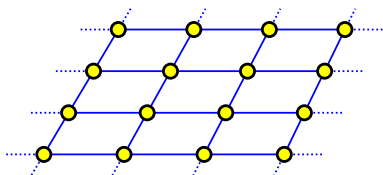
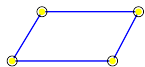
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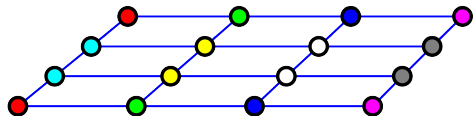
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Now back to  $4 \times 4$  (PBC's not drawn)



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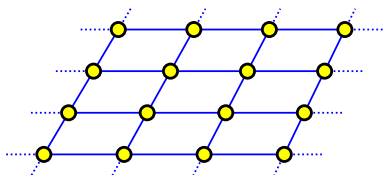
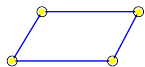
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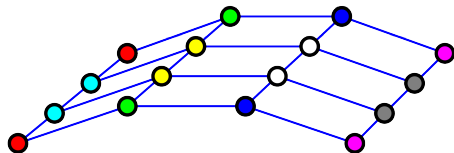
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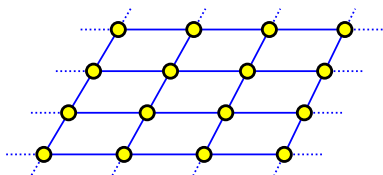
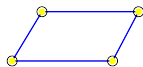
Consider 4x4 lattice with PBC's

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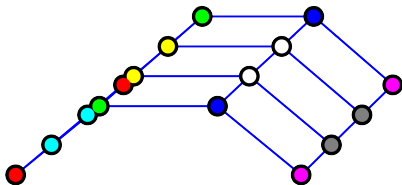
Hint: 2x2 lattice

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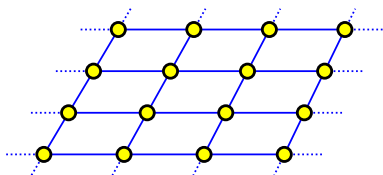
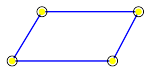
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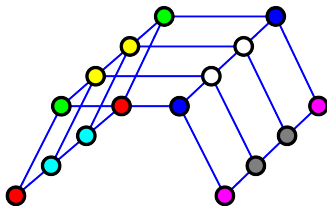
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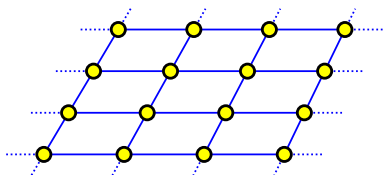
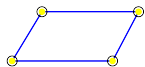
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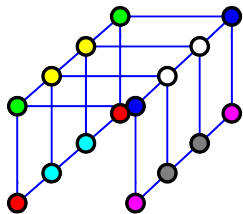
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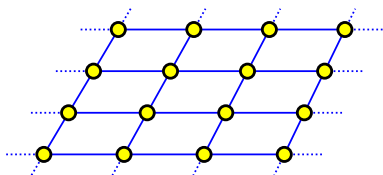
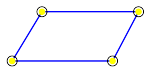
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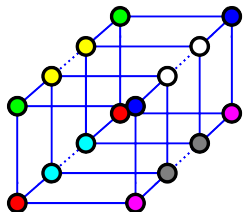
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- 1 transform to tube
- 2 reinsert “horizontal” PBC's
- 3 split in two  $2 \times 2 \times 2$  cubes (one inverted)
- 4 consider “front-back” PBC's  $\rightsquigarrow$  like colors linked - 4<sup>th</sup> dimension

## Dimensionality of finite-size systems: lattice case

Message so far: dimensionality of small lattice systems is ambiguous

More general claim: fermion can “feel” dimensionality only by “exploring”  
large spatial regions (low entropy  $\rightsquigarrow$  coherence)

So which lattice properties are always important?

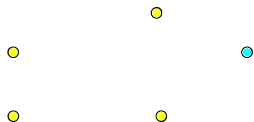
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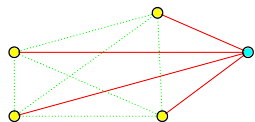
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Relevant measure: 2<sup>nd</sup> moment of local DOS

$$\langle \epsilon^2 \rangle_i \equiv \int_{-\infty}^{\infty} d\epsilon \epsilon^2 \rho_i(\epsilon) = \sum_j t_{ij}^2$$

uniform case  $\rightsquigarrow$  energy scale  $\sqrt{\langle \epsilon^2 \rangle} = \sqrt{Z}t$

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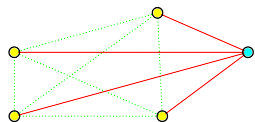
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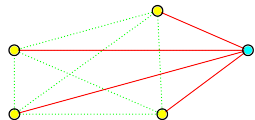
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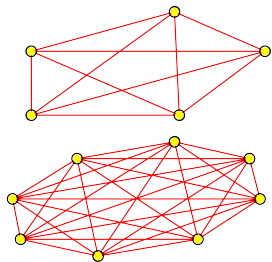
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3) Is lattice bipartite (hopping from “A” site to “B” site)?  $\rightsquigarrow$  AF order

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Example: **fully connected lattice** ( $Z = N - 1$ , non-bipartite,  $d$  undefined)

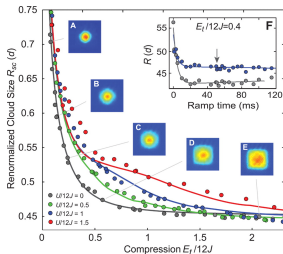
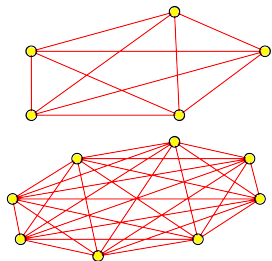
$\rightsquigarrow$  paramagnetic DMFT physics



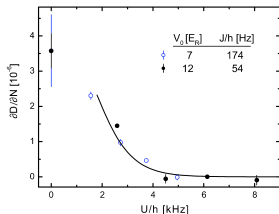
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[Schneider et al, Science (2008)]

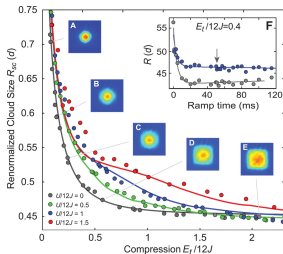
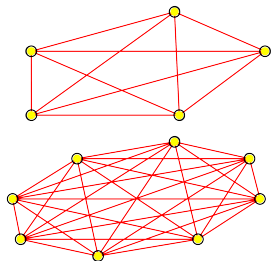


[Jördens et al., Nature (2008)]

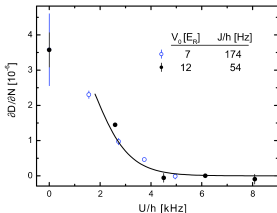
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Example: **fully connected lattice** ( $Z = N - 1$ , non-bipartite,  $d$  undefined)

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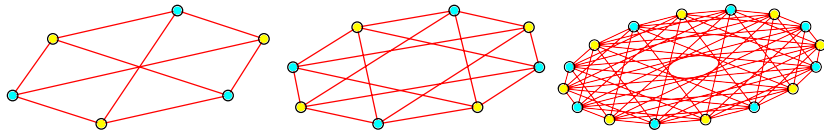


[Schneider et al, Science (2008)]



[Jördens et al., Nature (2008)]

Variation: **bipartite maximally connected lattice**  $\rightsquigarrow$  AF DMFT physics



# Universal probes for antiferromagnetic correlations and entropy in cold fermions on optical lattices

E. Gorelik, D. Rost, T. Paiva, R. Scalettar, A. Klümper, and N. Blümer

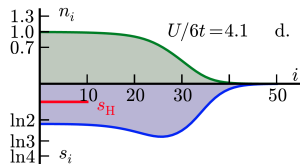
Phys. Rev. A **85**, 061602(R) (2012)

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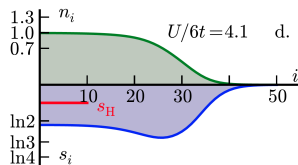


[Jördens et al., PRL **104**, 180401 (2010)]

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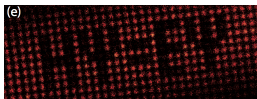
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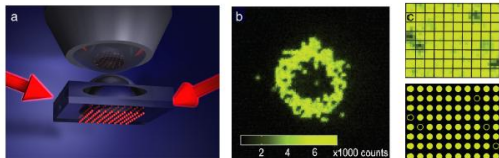
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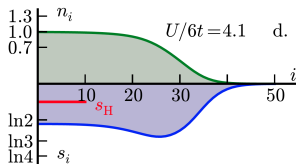
Experimental advantage of 2 dimensions:  
single-site resolution



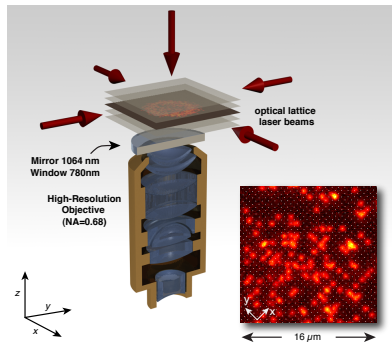
[Würtz et al., PRL 103, 080404 (2009)]



[Bakr et al., Science 329, 547 (2010)]



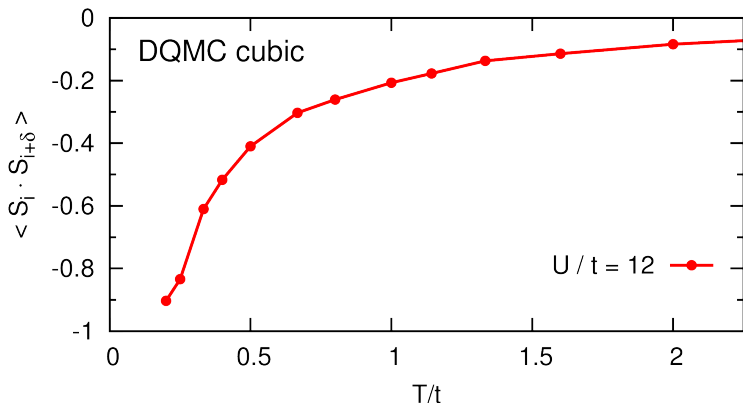
[Jördens et al., PRL 104, 180401 (2010)]



[Sherson et al., Nature 467, 68 (2010)]

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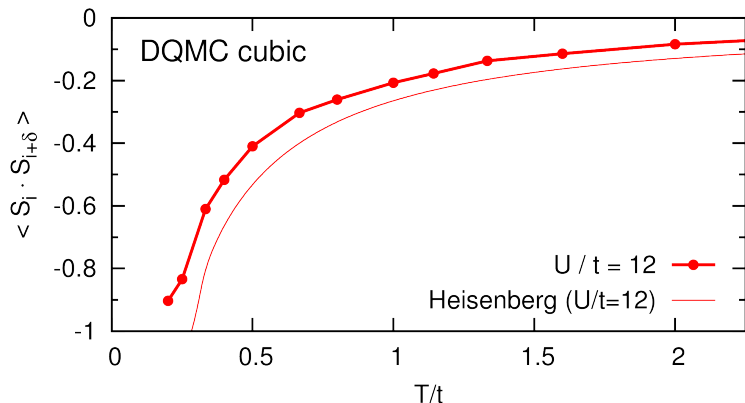
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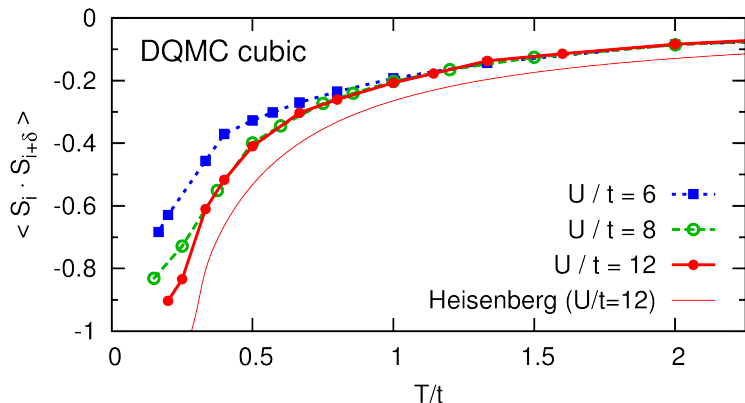
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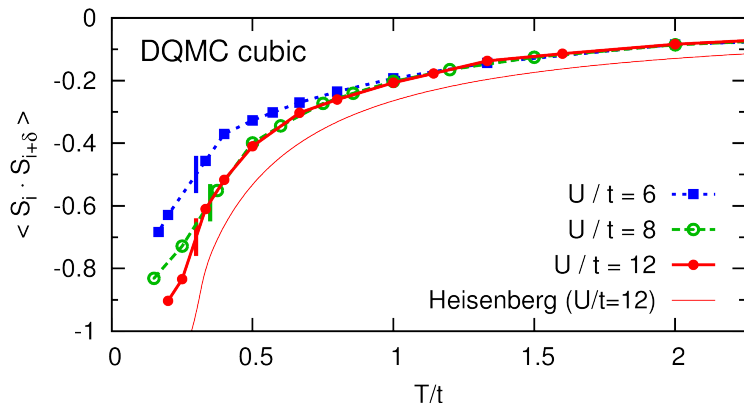
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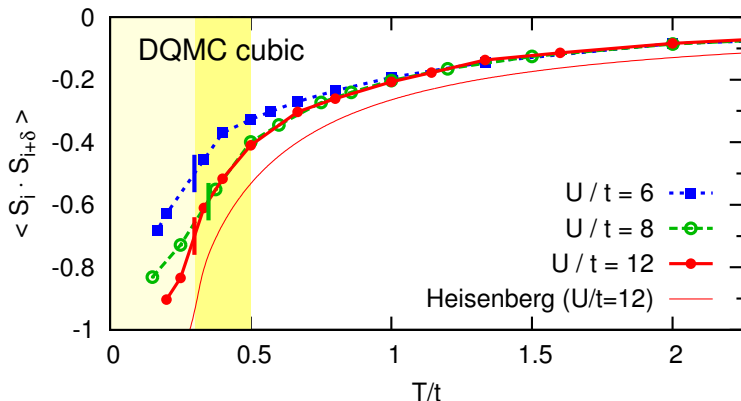
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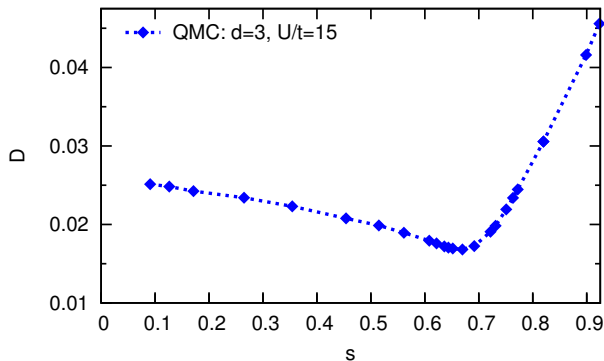
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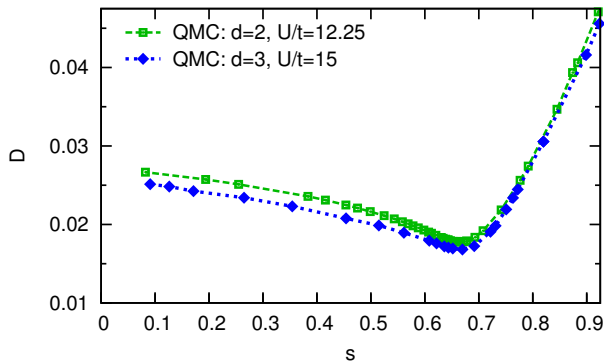
Interesting spin physics above  $T_N$ , not visible in NN correlations

# Double occupancy as a universal measure of AF correlations + entropy



Exact determinantal QMC:  
AF enhances  $D$  at low- $s$

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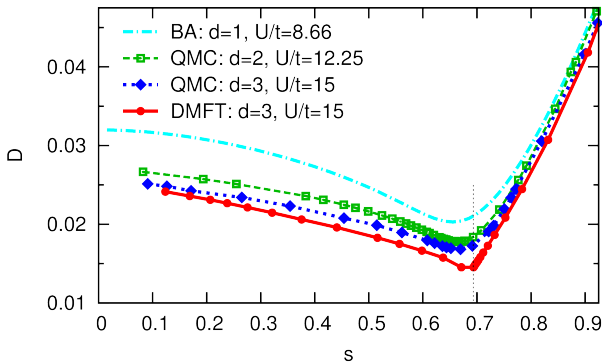
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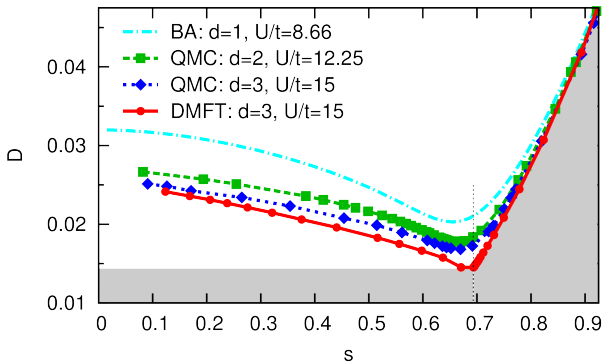
Universal minimum at

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$\ln(2)/2 \lesssim s \lesssim \ln(2)$  !

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AF enhancement of  $D$  is larger

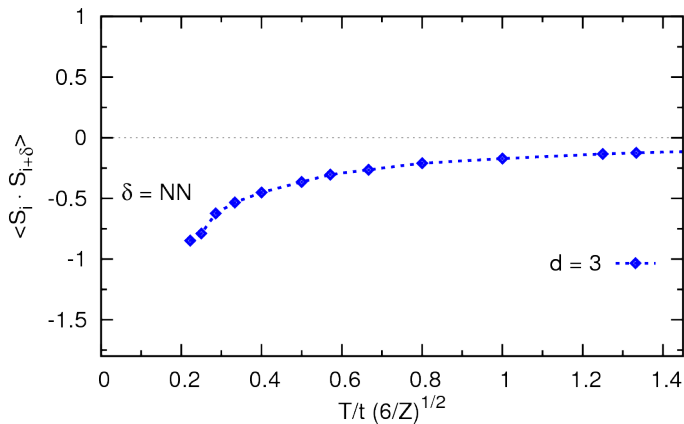
in lower dimensions:

$$D_0 = (1 - \langle \sigma_i \cdot \sigma_j \rangle) Z \frac{t^2}{2U^2} + \mathcal{O}(t^4/U^4)$$

$$\langle \sigma_i \cdot \sigma_j \rangle_0 = \begin{cases} -1.00 & DMFT \\ -1.20 & (d = 3) \\ -1.34 & (d = 2) \\ -1.77 & (d = 1) \end{cases}$$

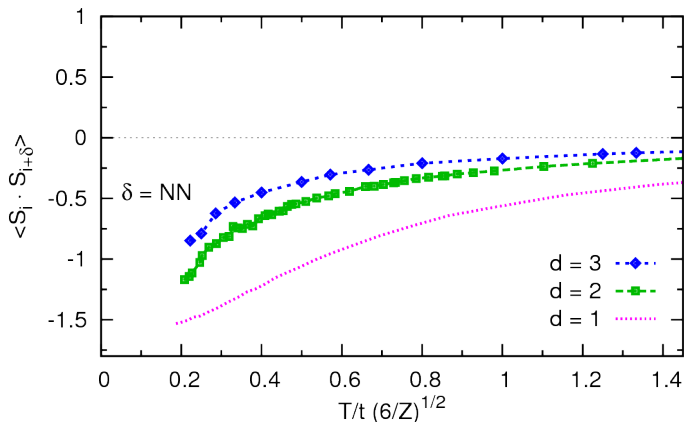
## Spin correlation functions: what range is needed?

NN spin correlations: high- $T$  tails



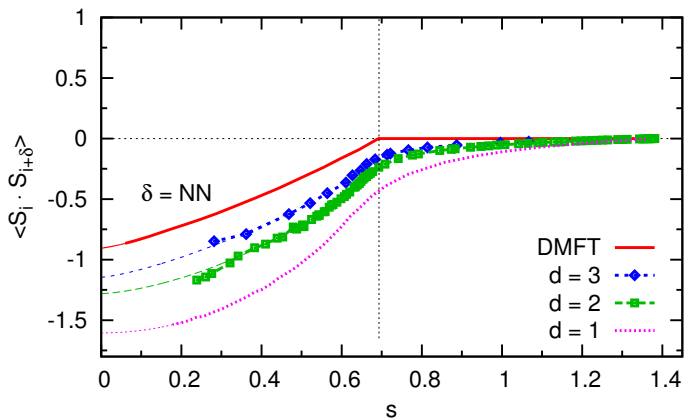
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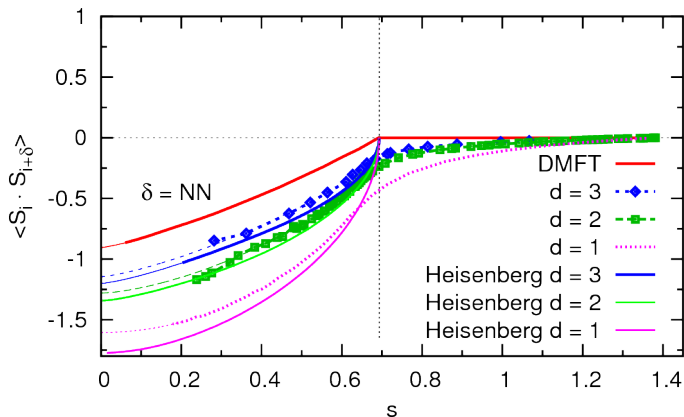


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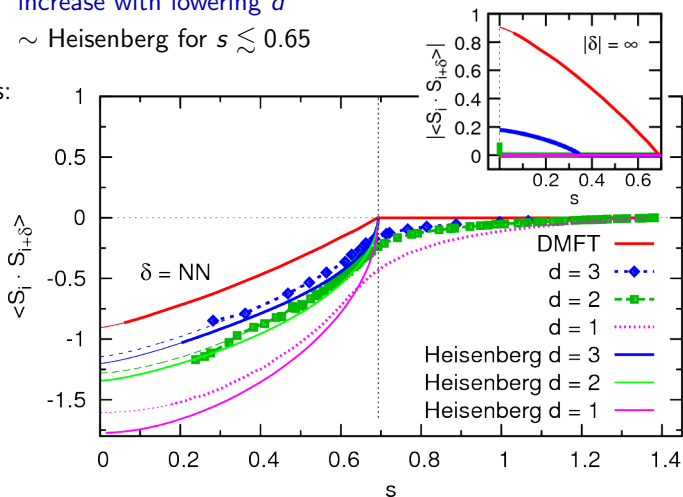
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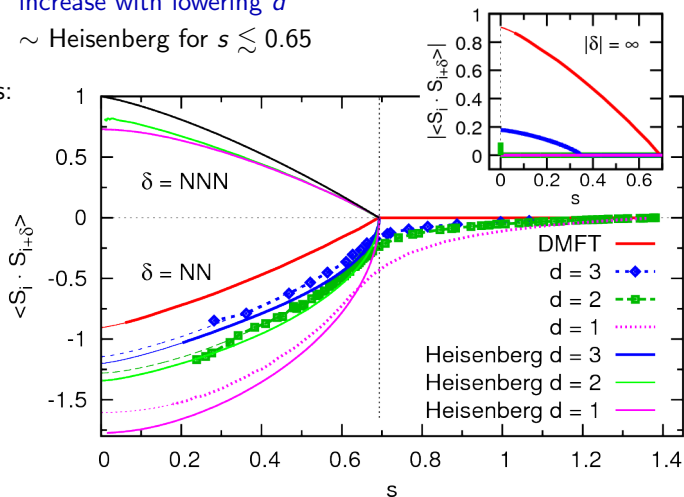
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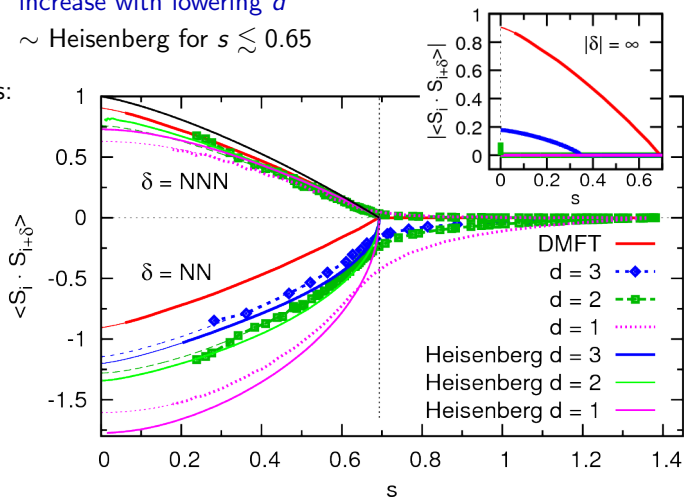
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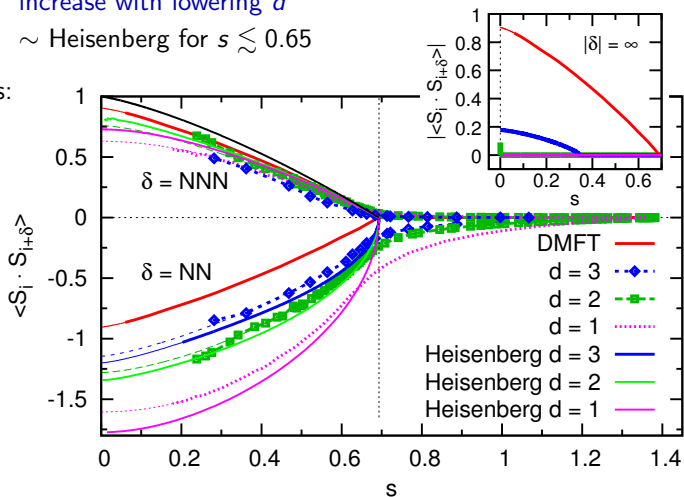
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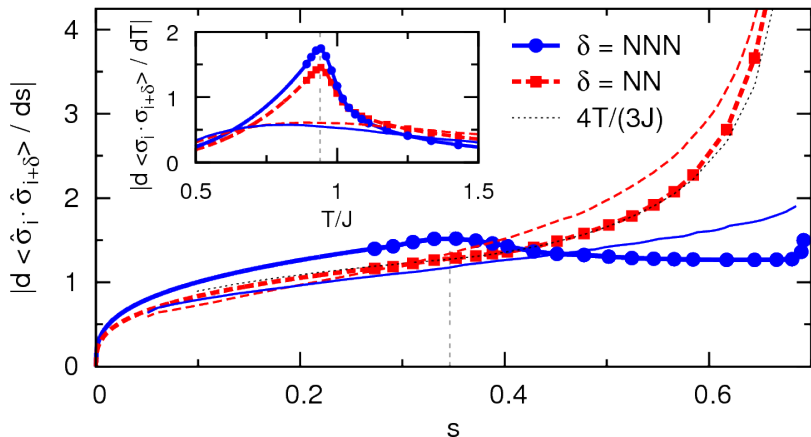
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NNN spin correlation function signals: Heisenberg regime, low entropy

# Spin correlation functions versus $s$ : why no signature of $T_N$ ?

Heisenberg model:  $d = 3$  (symbols),  $d = 2$  (thin lines)



Use of entropy as thermal variable can hide important physics!

# Discriminating antiferromagnetic signatures in ultracold fermions by tunable geometric frustration

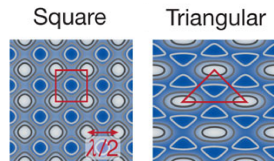
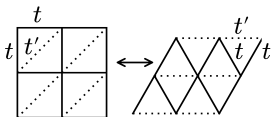
C.-C. Chang, R. T. Scalettar, E. V. Gorelik, and N. Blümer

[arXiv:1306.4687](https://arxiv.org/abs/1306.4687), to appear in *Phys. Rev. B*

# Tunable geometric frustration

Interpolate between square  
and triangular lattice:

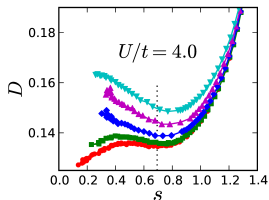
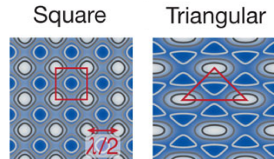
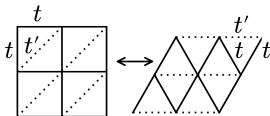
vary  $\alpha = t'/t$



# Tunable geometric frustration: weak coupling

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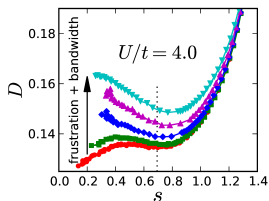
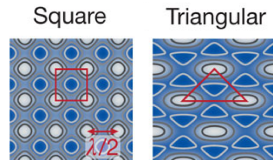
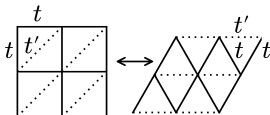


- $\alpha=0.0$
- $\alpha=0.2$
- ◆  $\alpha=0.4$
- ▲  $\alpha=0.6$
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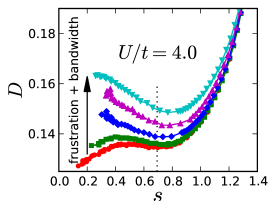
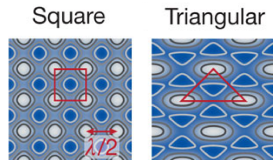
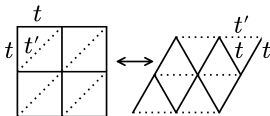
significant changes of  $D$  at all  $s$   
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Eliminate bandwidth effects:

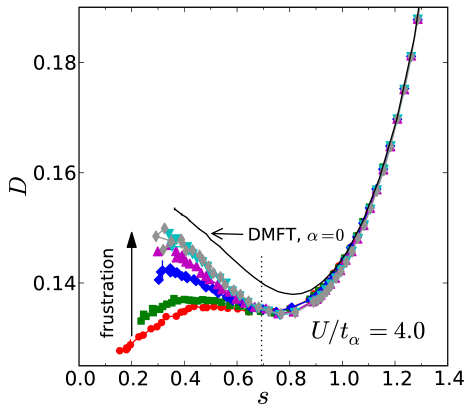
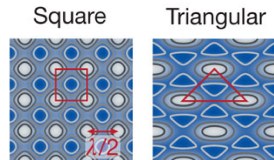
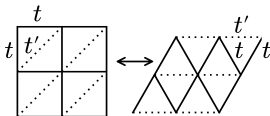
scale with  $t_\alpha = t\sqrt{Z_\alpha/Z_0}$ ;  $Z_\alpha = 4 + 2\alpha^2$

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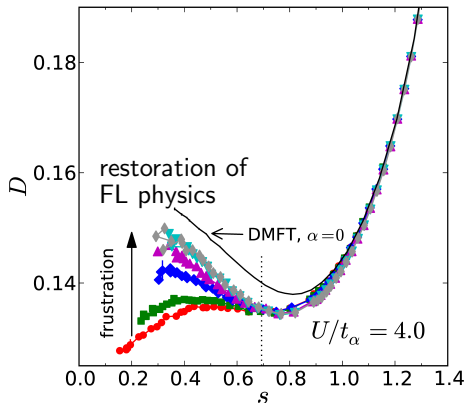
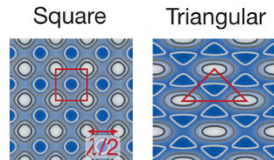
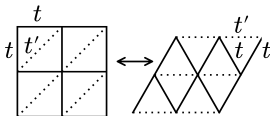


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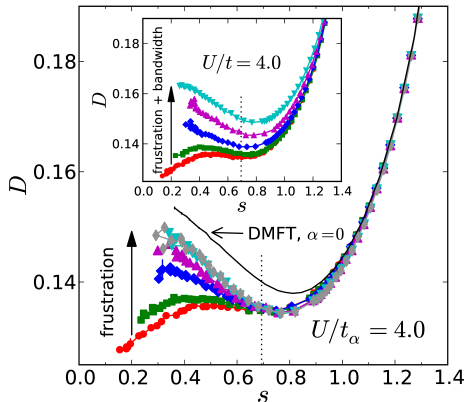
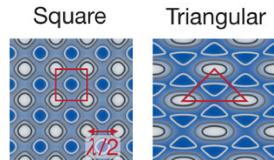
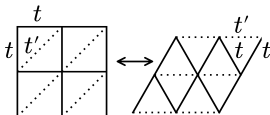


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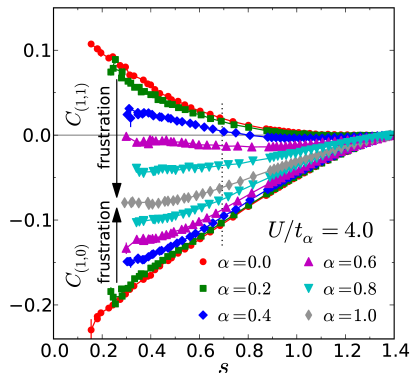
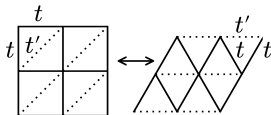
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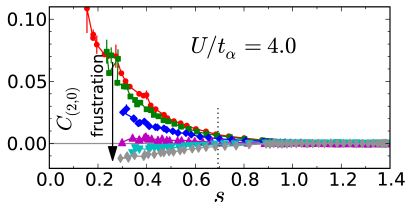
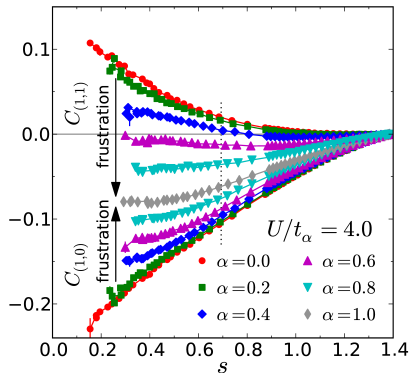
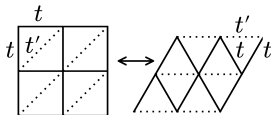
## Short range correlation functions



- robust evolution of  $C_{(1,0)}$  and  $C_{(1,1)}$  with  $\alpha$
- at  $\alpha = 1$   $C_{(1,0)} \equiv C_{(1,1)}$

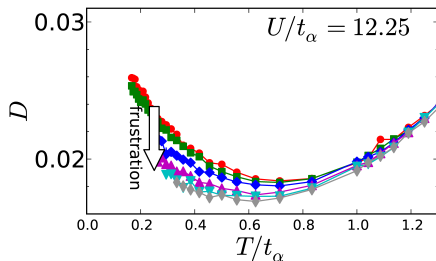
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- sharp signature in  $C_{(2,0)}$

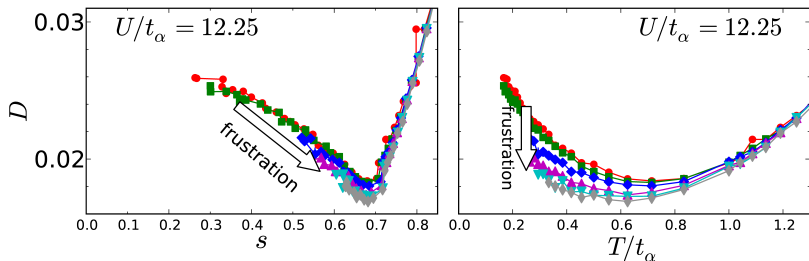
## Tunable geometric frustration: intermediate and strong coupling



### Strong coupling:

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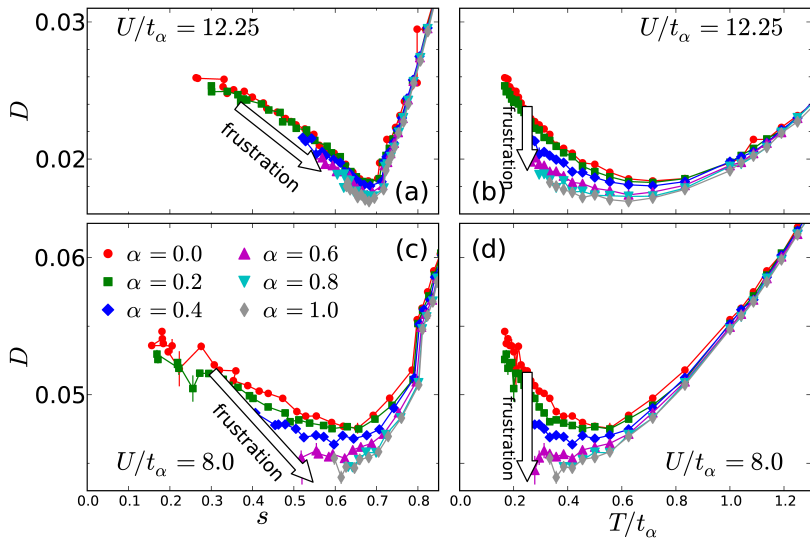
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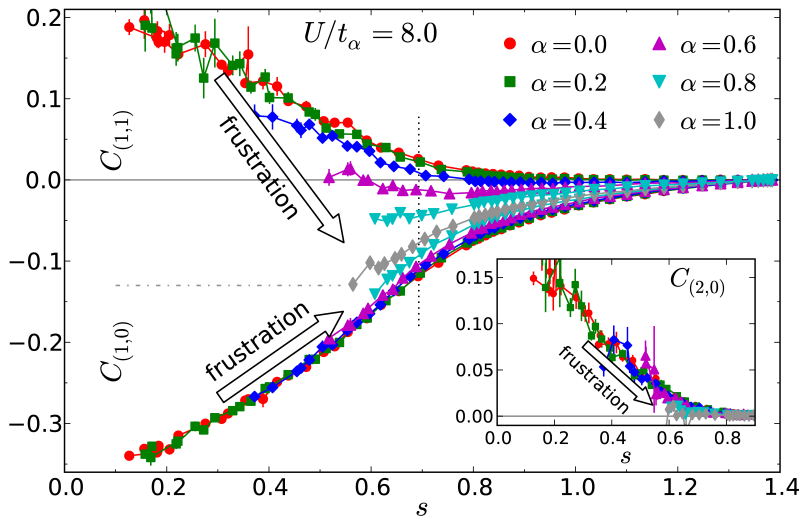
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- cancellation of  $\alpha$ -effect in  $D(s)$

# Tunable geometric frustration: intermediate and strong coupling



# Tunable geometric frustration: intermediate coupling

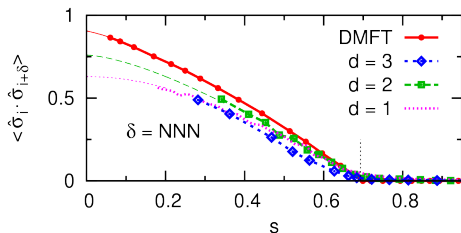
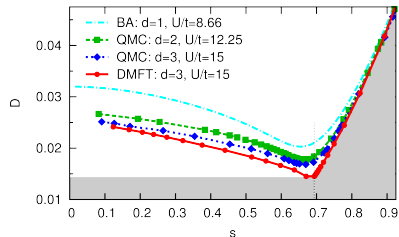




## Summary

NN AF correlations  $\leftrightarrow$  "finite-range antiferromagnetism"  $\leftrightarrow$  LRO

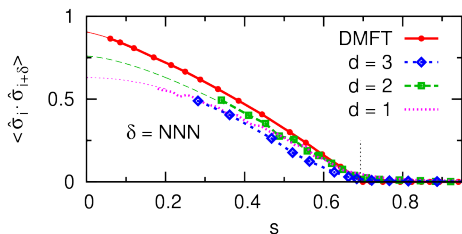
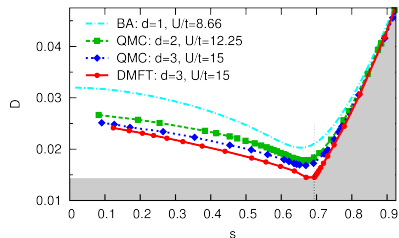
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Tunable frustration:  
discriminate magnetic from  
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[arXiv:1306.4687]

