

Orbital-selective Mott transitions in the anisotropic 2-band Hubbard model at finite temperatures

Nils Blümer

Outline

Motivation: OSMTs in $\text{Ca}_{2-x}\text{Sr}_x\text{RuO}_4$

Complementing QMC with high-frequency expansions

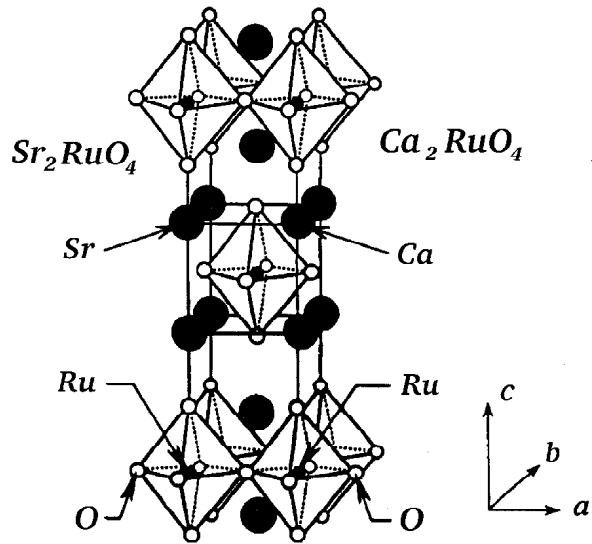
Mott transition in frustrated 1-band Hubbard model

High-precision ground state estimates from QMC

Orbital-selective Mott transitions in 2-band Hubbard model

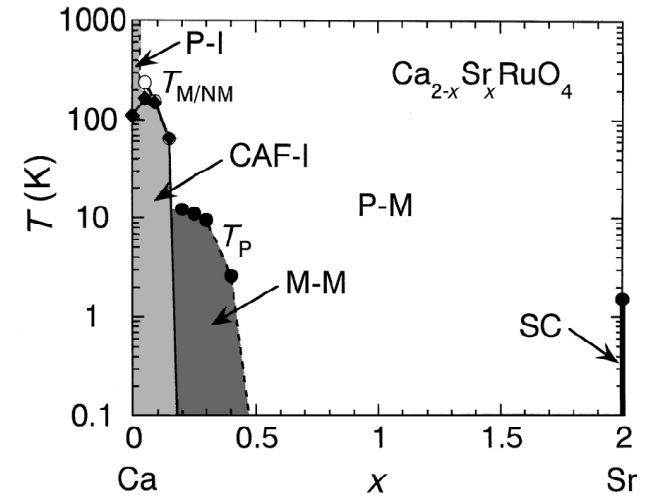
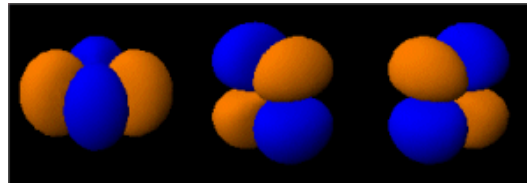
Summary

Motivation: OSMTs in $\text{Ca}_{2-x}\text{Sr}_x\text{RuO}_4$



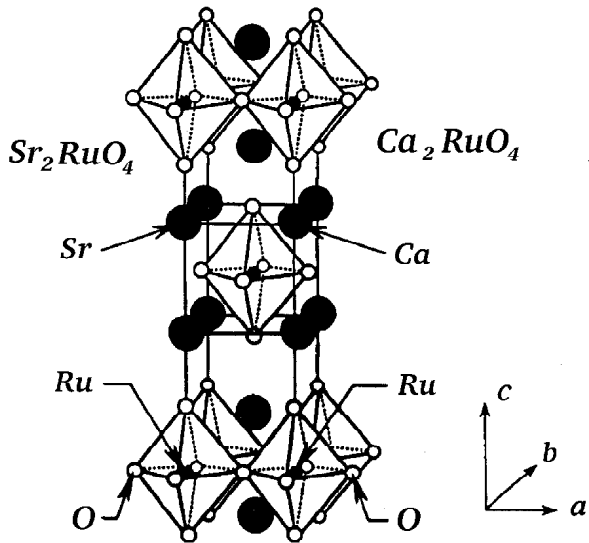
$\text{Ca}_{2-x}\text{Sr}_x\text{RuO}_4$:
isostructural to $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$

Sr_2RuO_4 : quasi-2d FL,
spin-triplet superconductor



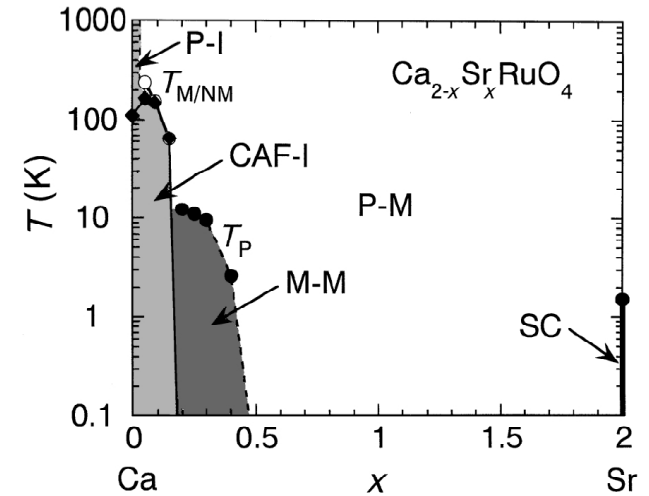
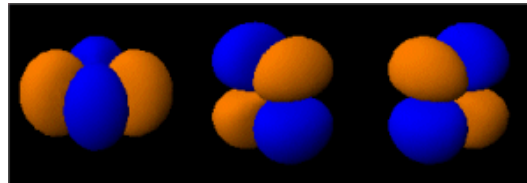
[Nakatsuji, Maeno, PRL **84**, 2666 (2000)]

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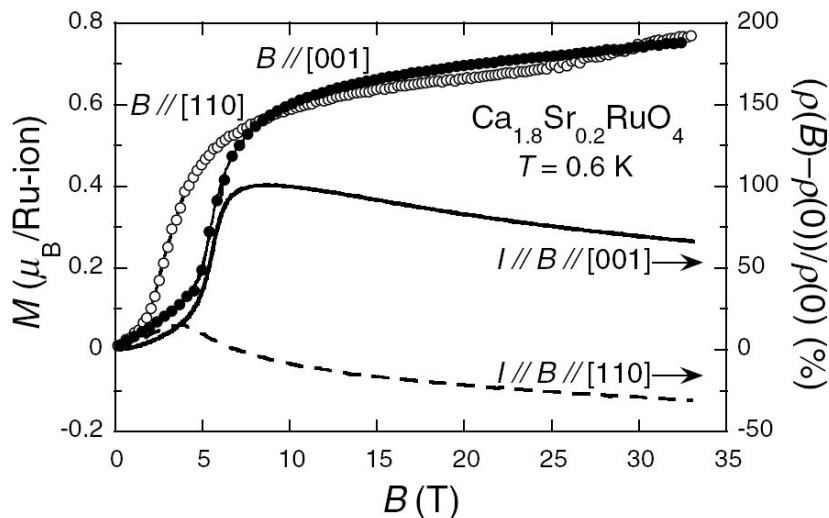


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[Nakatsuji, Maeno, PRL **84**, 2666 (2000)]



saturation moment, susceptibility

$\rightsquigarrow S = 1/2$ system for $x \gtrsim 0.2$ (not $S = 1$)

strongly anisotropic magnetoresistance

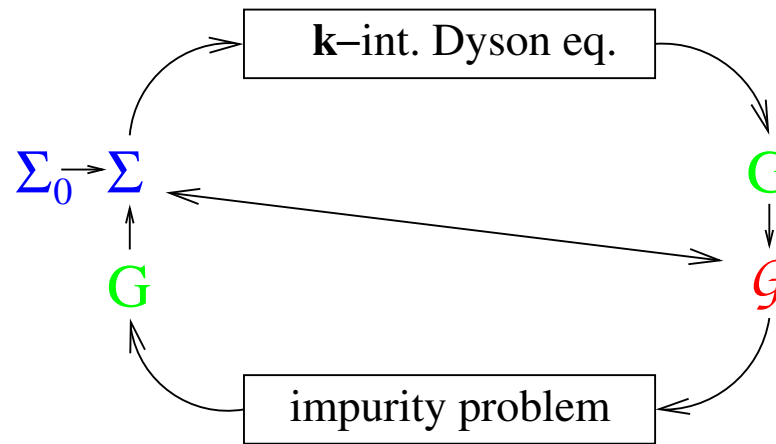
orbital-selective Mott metal-insulator transitions
for $x \approx 0.5$, $x \approx 0.2$?

high-precision methods needed!

[Nakatsuji *et al.*, PRL **90**, 137202 (2003)]

Complementing QMC with high-frequency expansions

Iterative solution of
DMFT equations



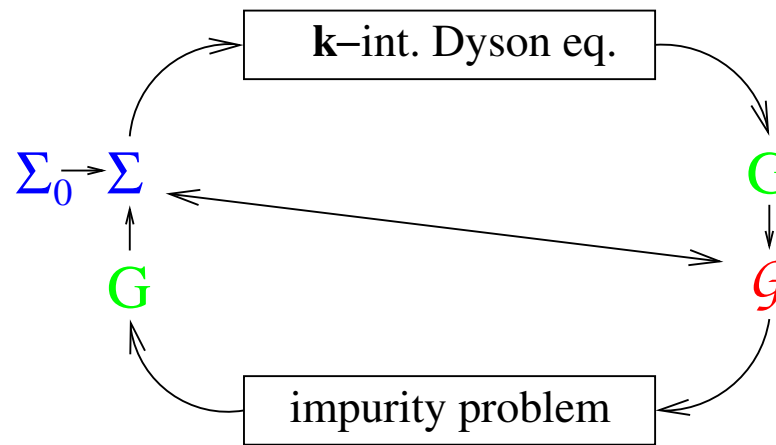
← DOS / t_{ij} / $\epsilon_{\nu k}$

Fourier transformations

← local interactions

Complementing QMC with high-frequency expansions

Iterative solution of DMFT equations

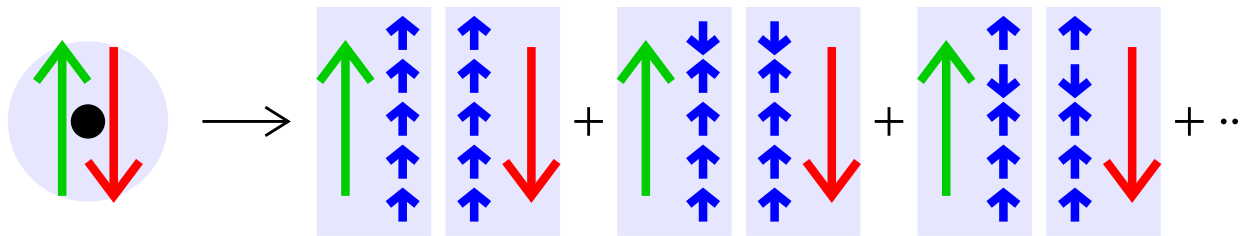


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Fourier transformations

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QMC: discretization $\beta = \Lambda \Delta\tau$, Trotter decoupling, discrete Hubbard-Stratonovich transformation



Wick theorem:

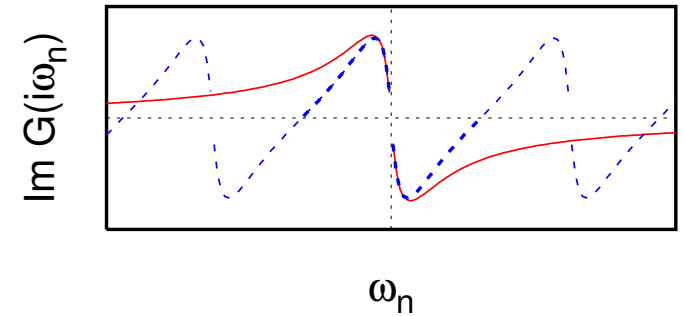
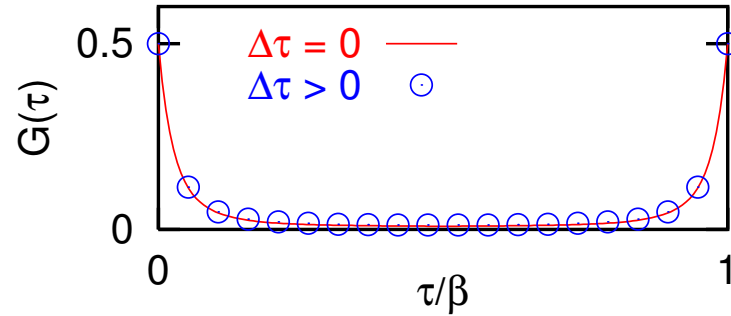
$$G = \frac{\sum M \det\{M\}}{\sum \det\{M\}}$$

Metropolis MC importance sampling over auxiliary Ising field, 2^Λ configurations, $50 \lesssim \Lambda \lesssim 400$

+ numerically exact - effort scales as T^{-3}

- no info for $\omega \gtrsim \omega_{\text{Nyquist}}$ \rightsquigarrow problems with Fourier transforms

naive discrete FT
 \rightsquigarrow oscillations
 (instead of
 $G(\omega) \xrightarrow{\omega \rightarrow \infty} 1/\omega$)



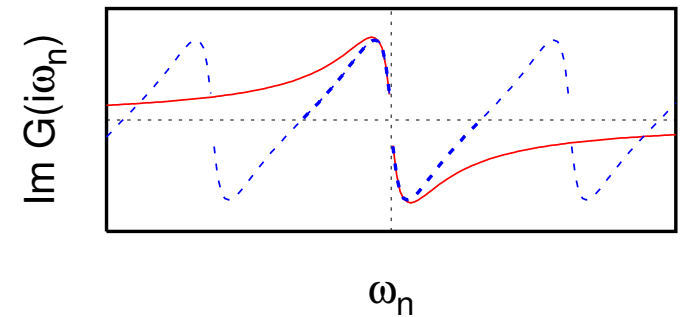
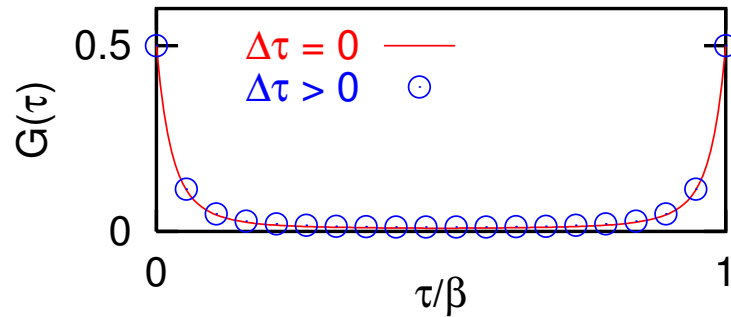
1st solution: correct unphysical behavior for $|\omega| \lesssim \omega_{\text{Nyquist}}$ by transformation [Ulmke]

2nd solution: interpolate $G_{\text{QMC}}(\tau)$
 by cubic splines [Jarrell, Krauth, . . .]

but: natural boundary conditions
 not appropriate for $G(\tau)$:

- adjust bc's [Oudovenko]
- spline-fit only difference w.r.t.
 reference problem:
 - IPT [Jarrell]

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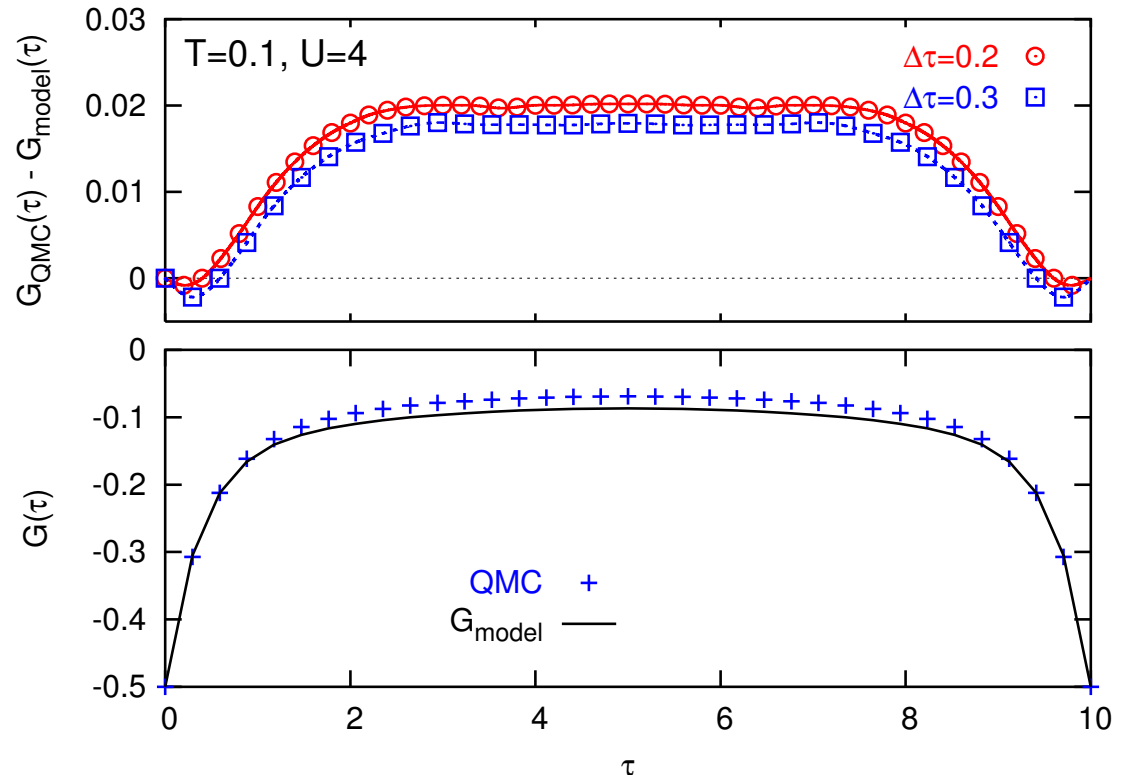


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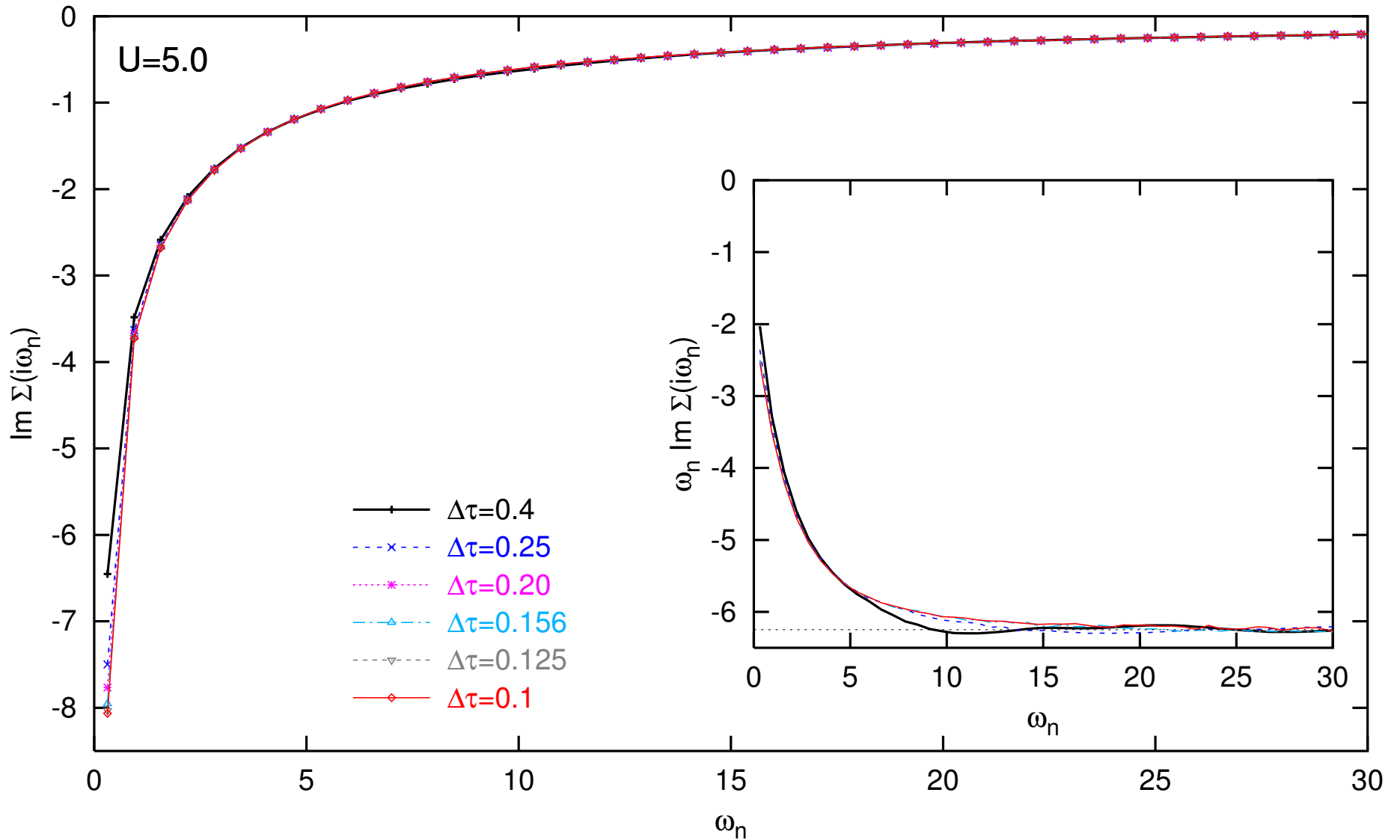
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- spline-fit only difference w.r.t. reference problem:
 - IPT [Jarrell]
 - high-frequency expansion for Σ + fit-param. [Knecht, NB]

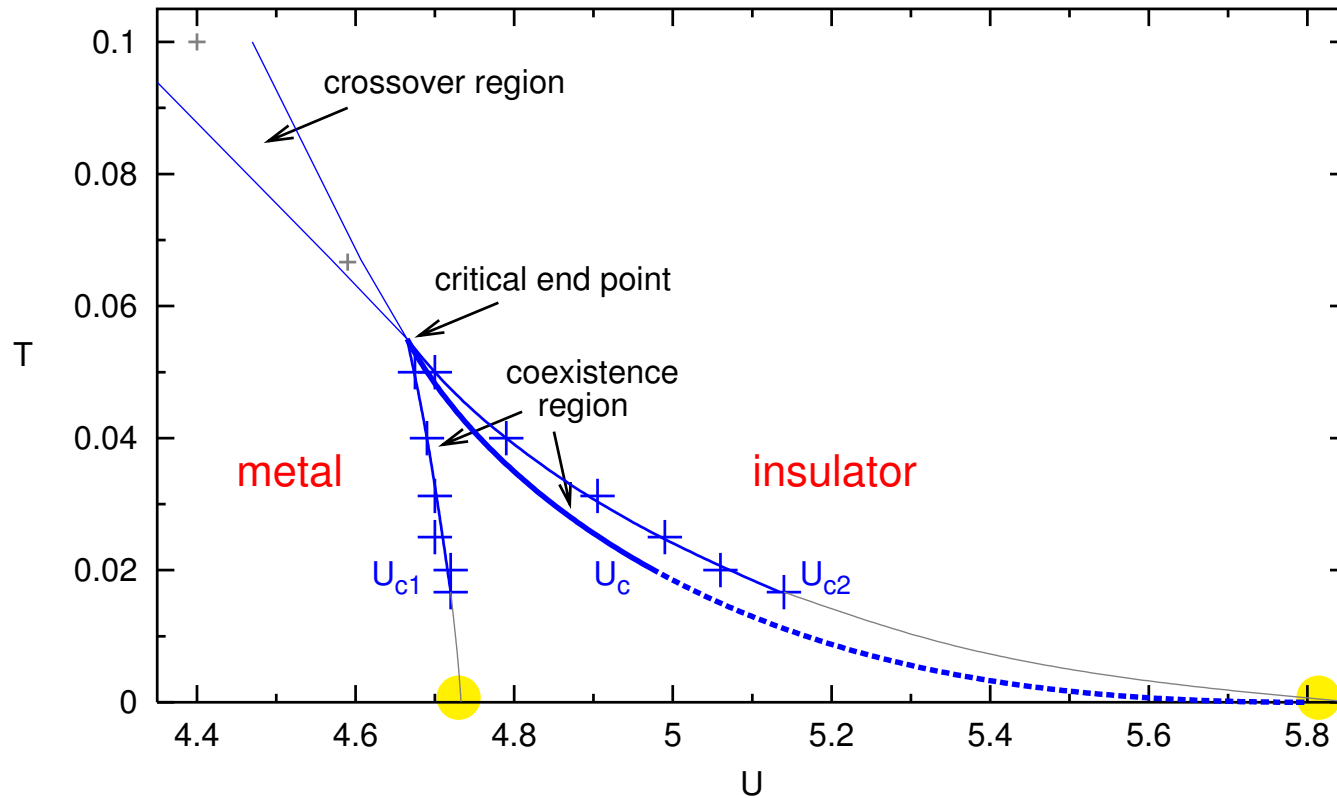


Sensitive test: self-energy $\Sigma(i\omega_n)$ for insulating phase ($T = 0.1, U = 5.0$)



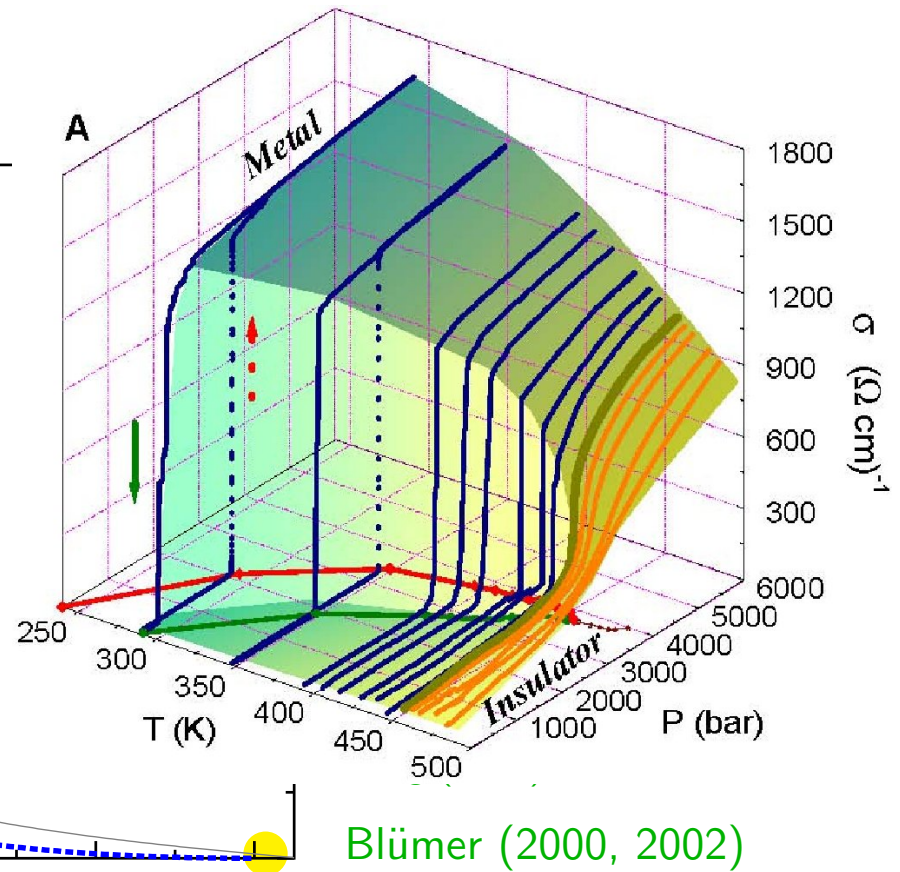
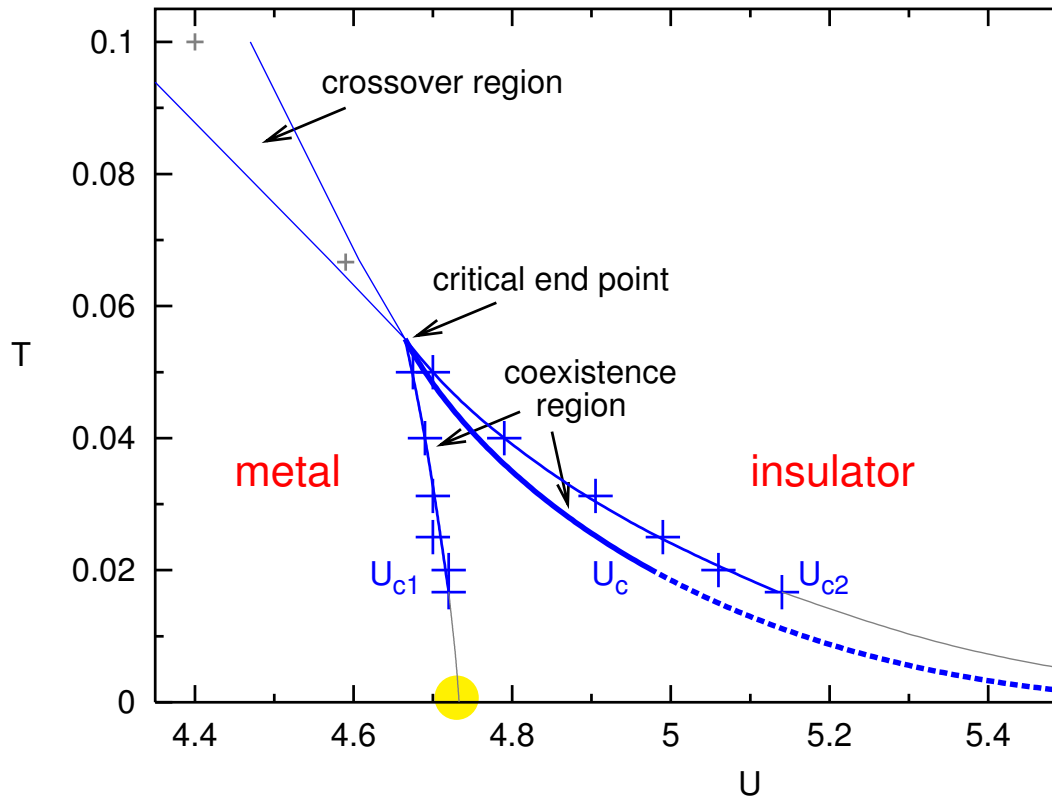
Rapid convergence at all frequencies for “QMC + $1/\omega$ ” DMFT solver

Mott transition in frustrated 1-band Hubbard model



- Georges and Krauth (1993)
- Rozenberg, Kotliar, Zhang (1994)
- Georges *et al.* (RMP, 1996)
- Schlipf *et al.* (1999)
- Rozenberg, Chitra, Kotliar (1999)
- Krauth (2000)
- Bulla (1999, 2001)
- Joo, Oudovenko (2001)
- Tong (2001)
- Blümer (2000, 2002)

Mott transition in frustrated



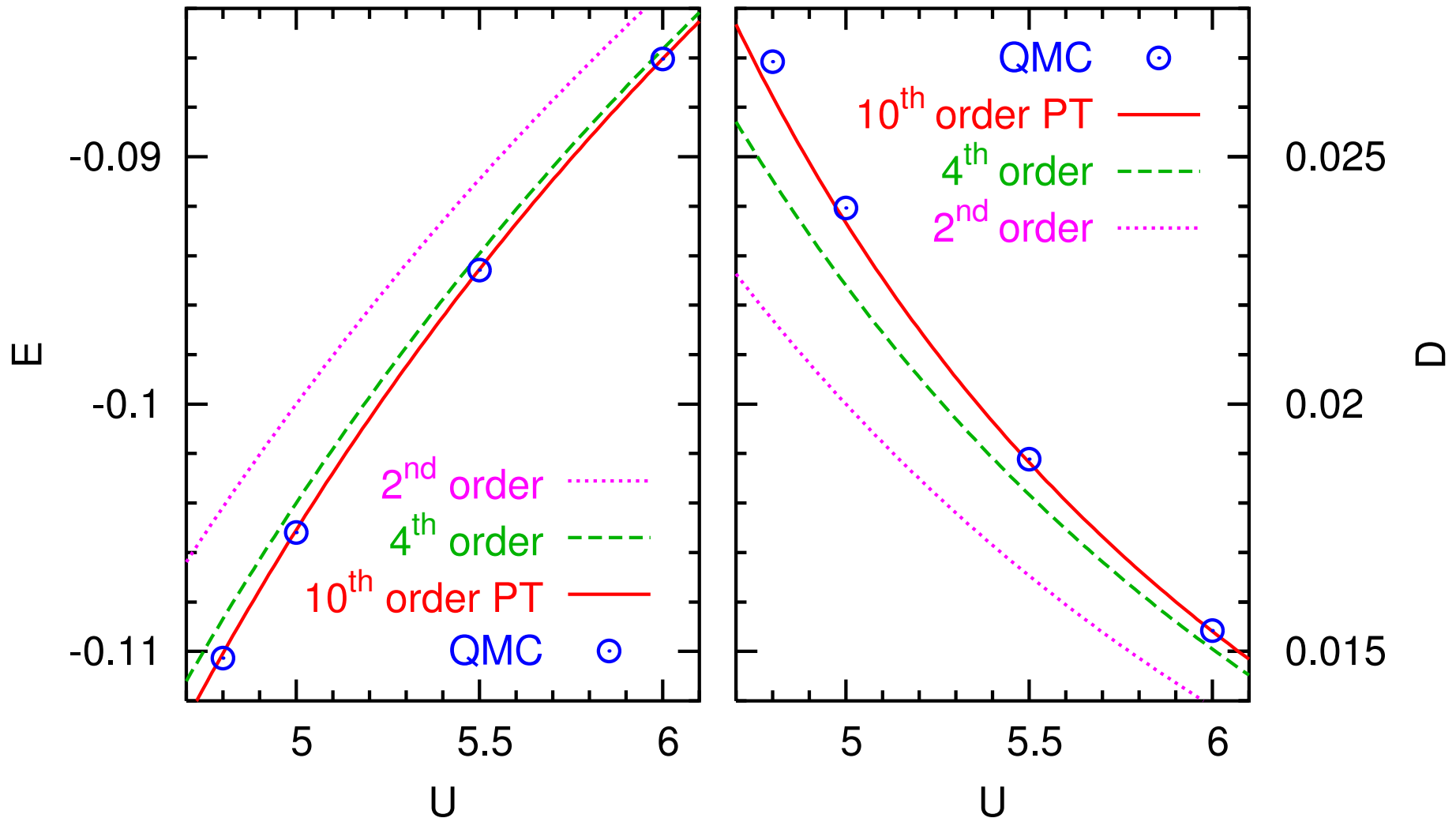
1st order line from
$$\frac{dU_c(T)}{dT} = f(T, U_c(T)); \quad f(T, U) := \frac{\Delta E(T, U)}{T \Delta D(T, U)}$$

low- T asymptotics from
$$U_c(T) = U_c^0 - \sqrt{\frac{2S_0 T}{a}} + \frac{\gamma_0}{4S_0} T + \mathcal{O}(T^{3/2})$$

High-precision energetics needed, even for $T \rightarrow 0$

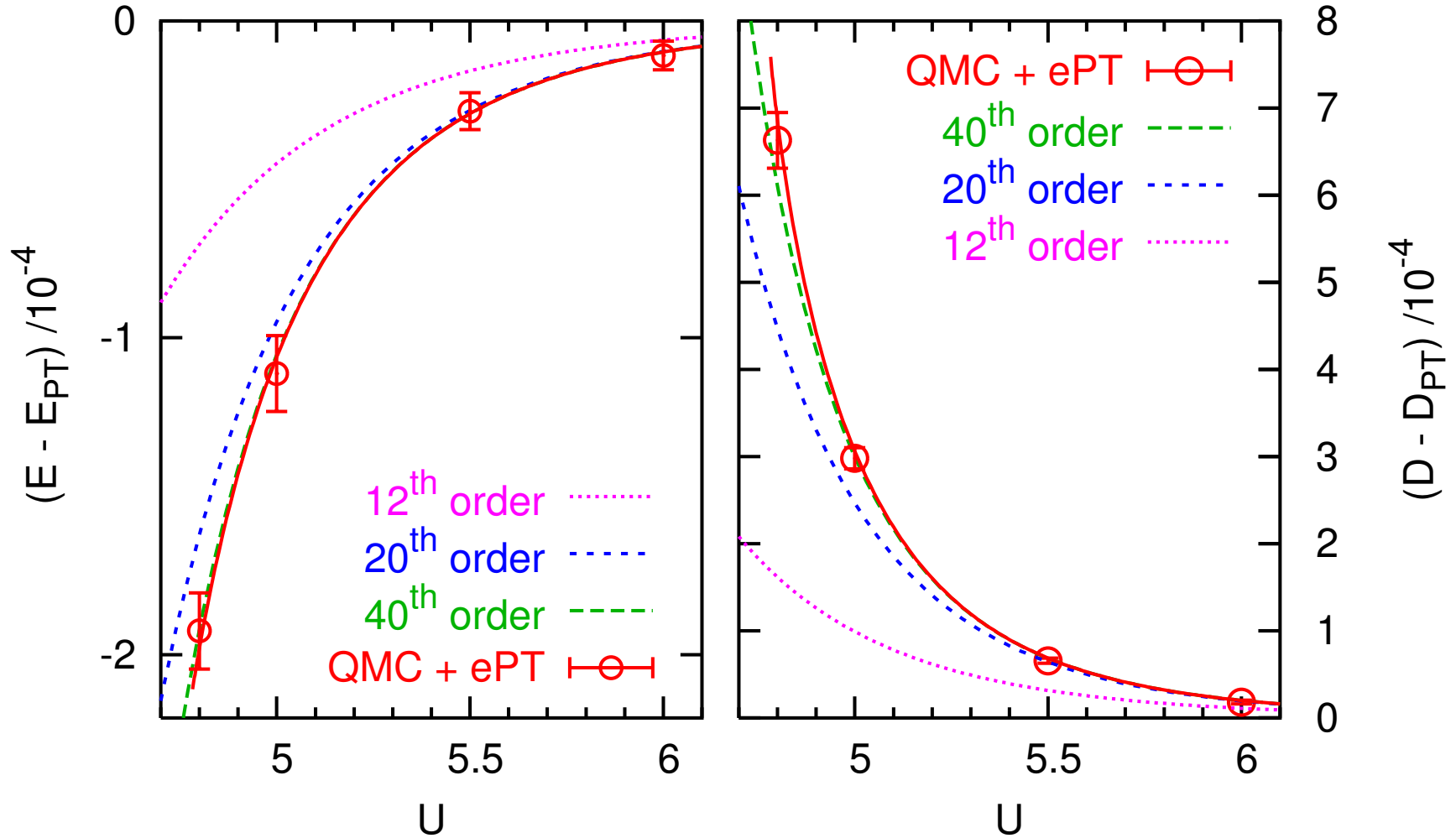
High-precision ground state estimates from QMC

QMC (+ high-frequency expansion) vs. strong-coupling PT for insulating phase



Excellent agreement at $U = 6.0$, deviations below.

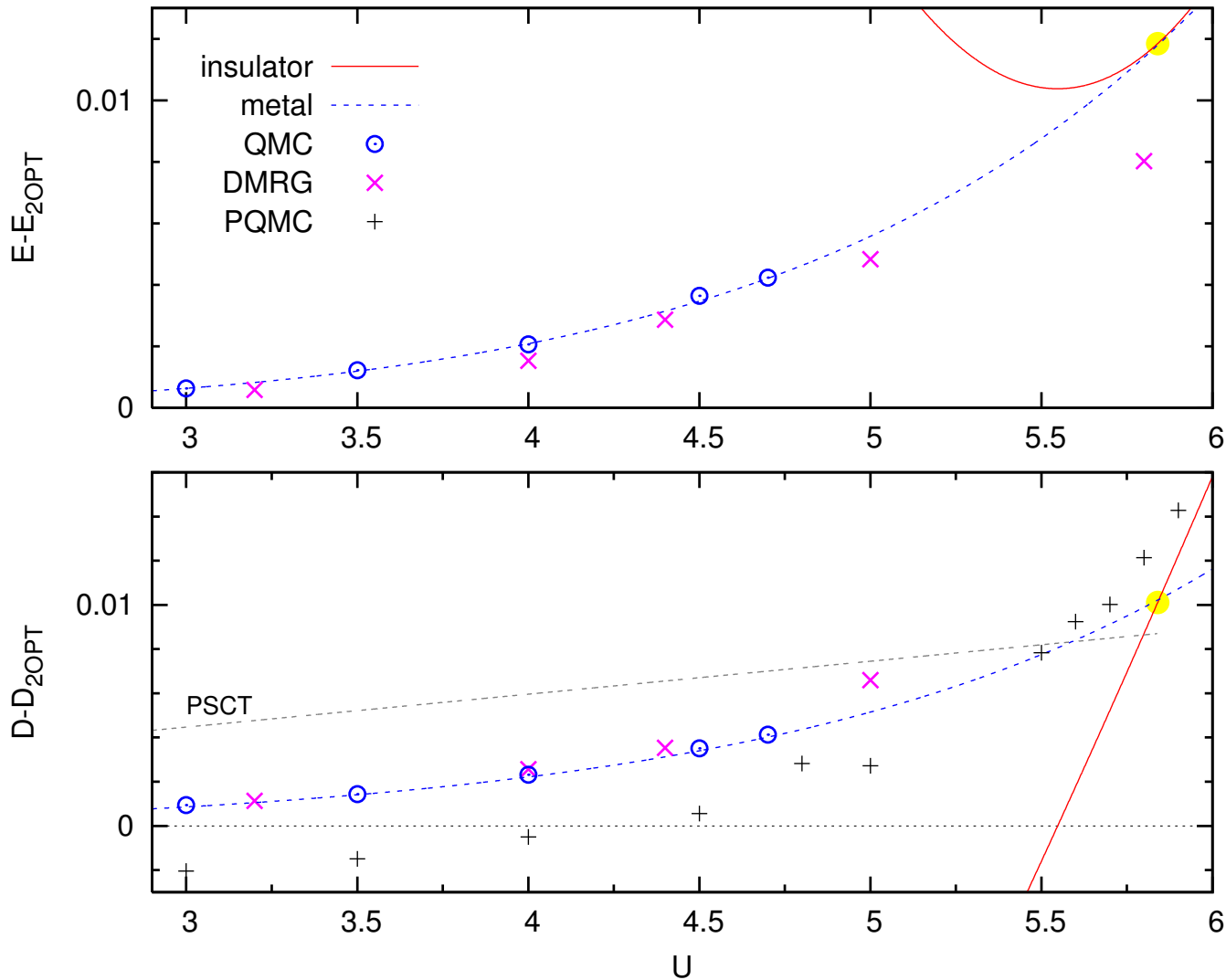
Higher resolution plots: differences w.r.t. 10th order PT



ePT: extrapolation of PT to infinite order [NB, Kalinowski, Phys. Rev B **71**, 195102 (2005)]
 \rightsquigarrow critical interaction U_{c1} , critical exponents, benchmark ($\Delta E \lesssim 10^{-6}$)

QMC algorithm has passed only available authoritative (1-band) test!

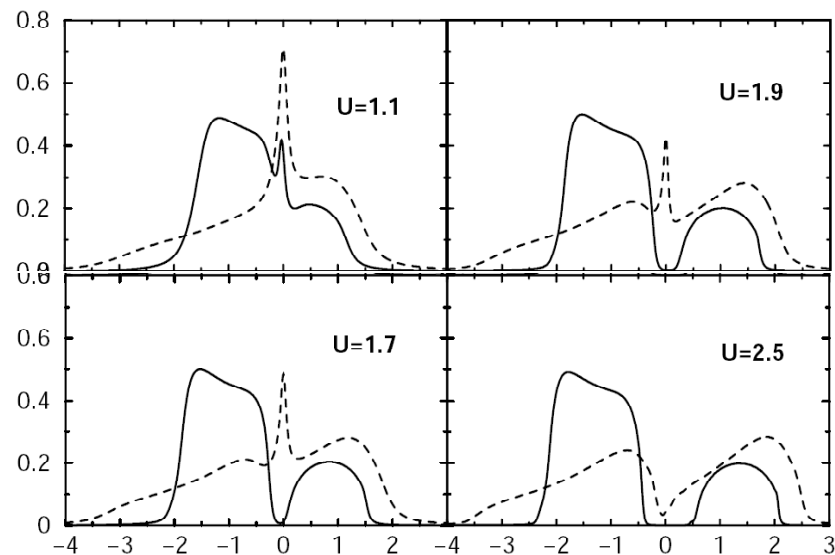
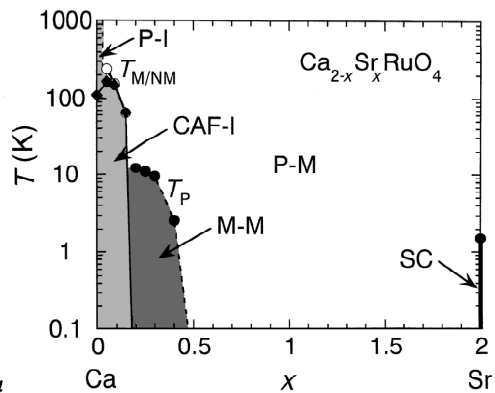
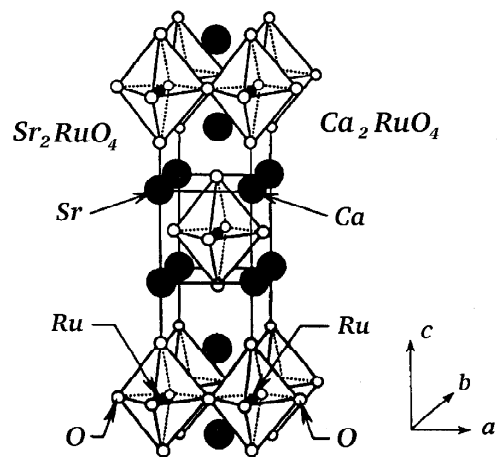
Metallic phase: differences w.r.t. 2nd order weak-coupling PT for E and D



QMC-fit consistent with ePT for insulator; larger deviations of PSCT, DMRG, and PQMC.

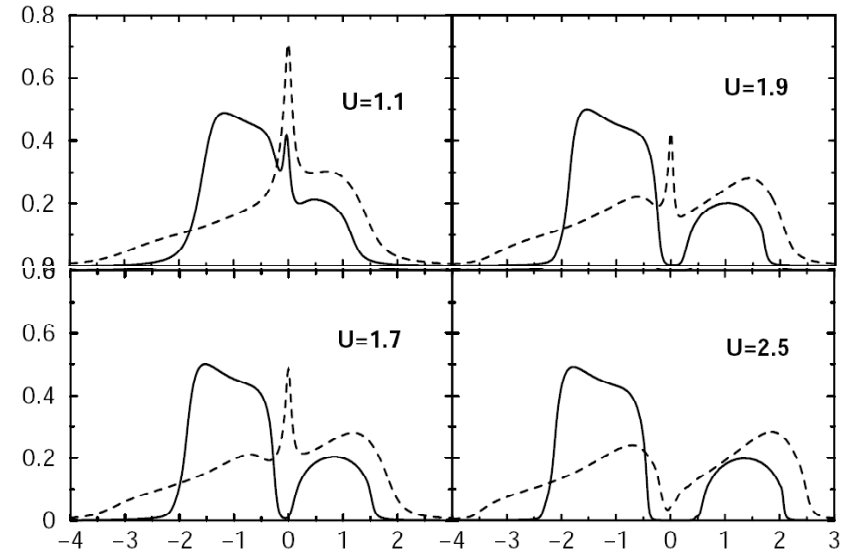
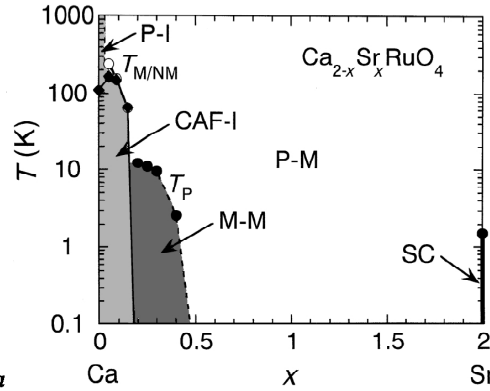
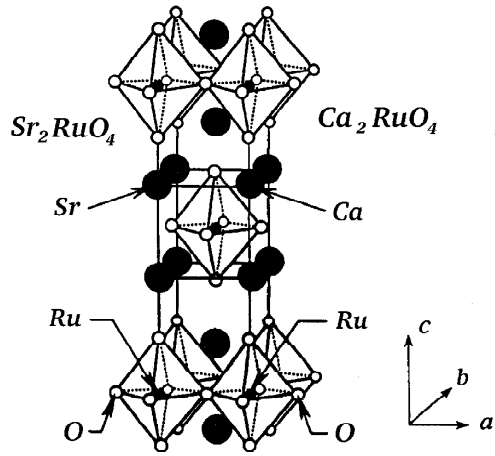
4th order PT coefficient corrected by QMC ($-62 \rightarrow +5$)

Orbital-selective Mott transitions in 2-band Hubbard model



[LDA+DMFT(NCA): Anisimov *et al.* (2002)]

Orbital-selective Mott transitions in 2-band Hubbard model



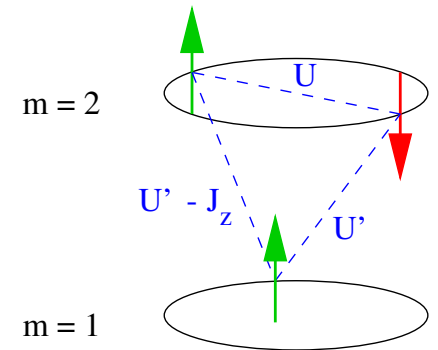
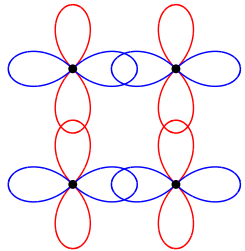
[LDA+DMFT(NCA): Anisimov *et al.* (2002)]

Minimal model: 2-band Hubbard model with orbital-dependent hopping

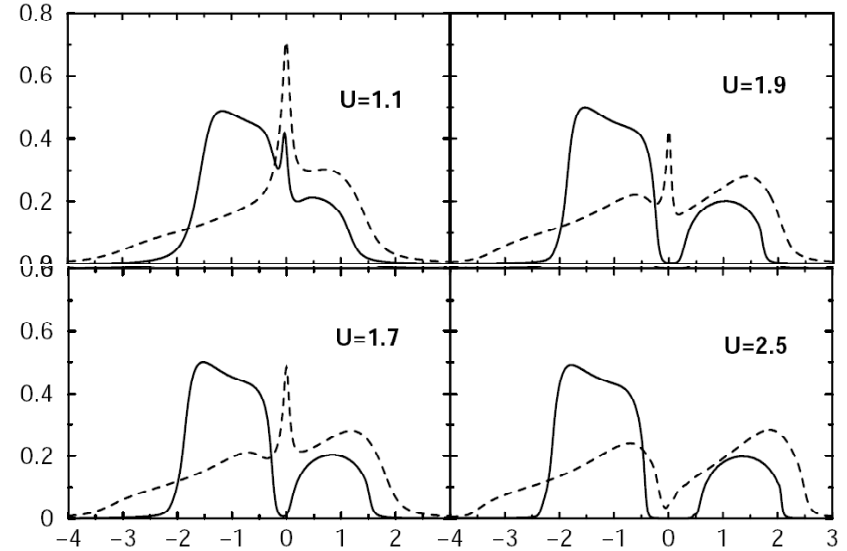
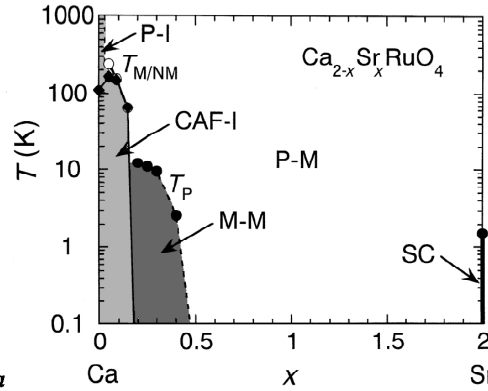
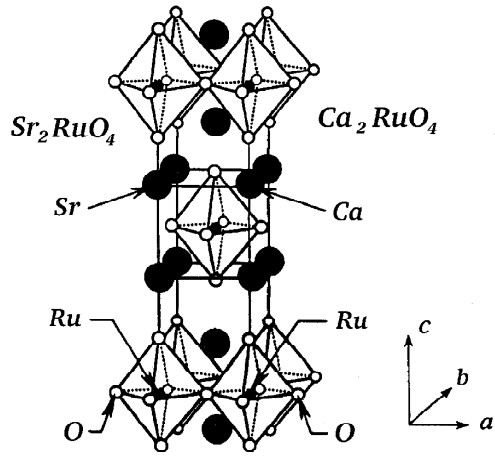
$$H = \sum_{m=1}^2 \left[- \sum_{\langle ij \rangle \sigma} t_m c_{im\sigma}^\dagger c_{jm\sigma} + U \sum_i n_{im\uparrow} n_{im\downarrow} \right]$$

$$+ \sum_{i\sigma\sigma'} (U' - \delta_{\sigma\sigma'} J_z) n_{i1\sigma} n_{i2\sigma'}$$

$$+ \frac{1}{2} J_\perp \sum_{i\sigma} \left[c_{i1\sigma}^\dagger \left(c_{i2\bar{\sigma}}^\dagger c_{i1\bar{\sigma}} + c_{i1\bar{\sigma}}^\dagger c_{i2\bar{\sigma}} \right) c_{i2\sigma} + \text{h.c.} \right]$$



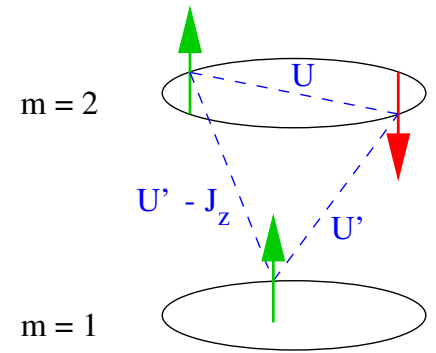
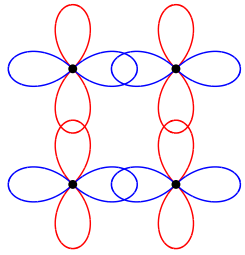
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 \end{aligned}$$



For Bethe DOS, $t_2 = 2 t_1$: two 1st order MITs for $U' = J_z = J_\perp = 0$ (trivial)

two distinct MITs for $J_z = J_\perp = U/4$ [Koga *et al.*, PRL **92**, 216402 (2004)]

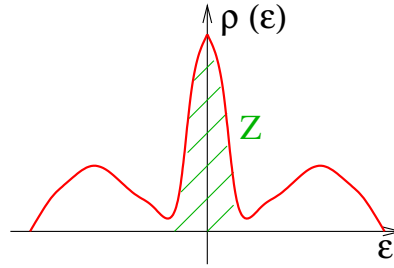
single MIT or **two OSMTs** for $J_\perp = 0, J_z = U/4$?

Earlier DMFT-QMC results: single Mott transition in J_z model ($J_\perp = 0$)

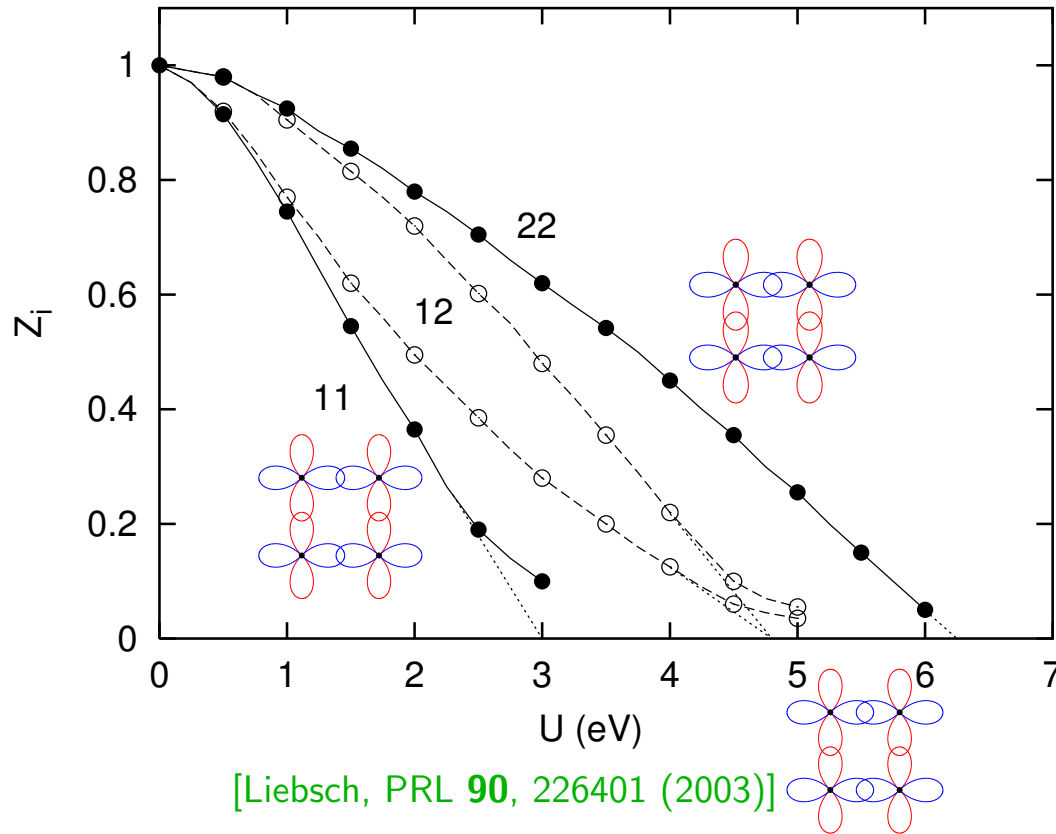
quasiparticle weight $Z_i = m/m_i^*$

discrete QMC estimate:

$$Z_i \approx [1 - \text{Im}\Sigma(i\pi T)/\pi T]^{-1}$$



$$J_z = 0.2, U' = U - 0.4$$

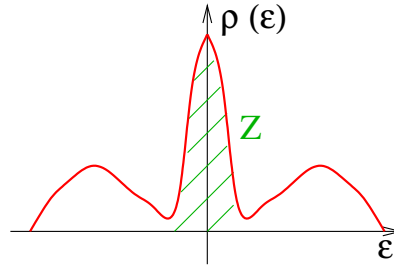


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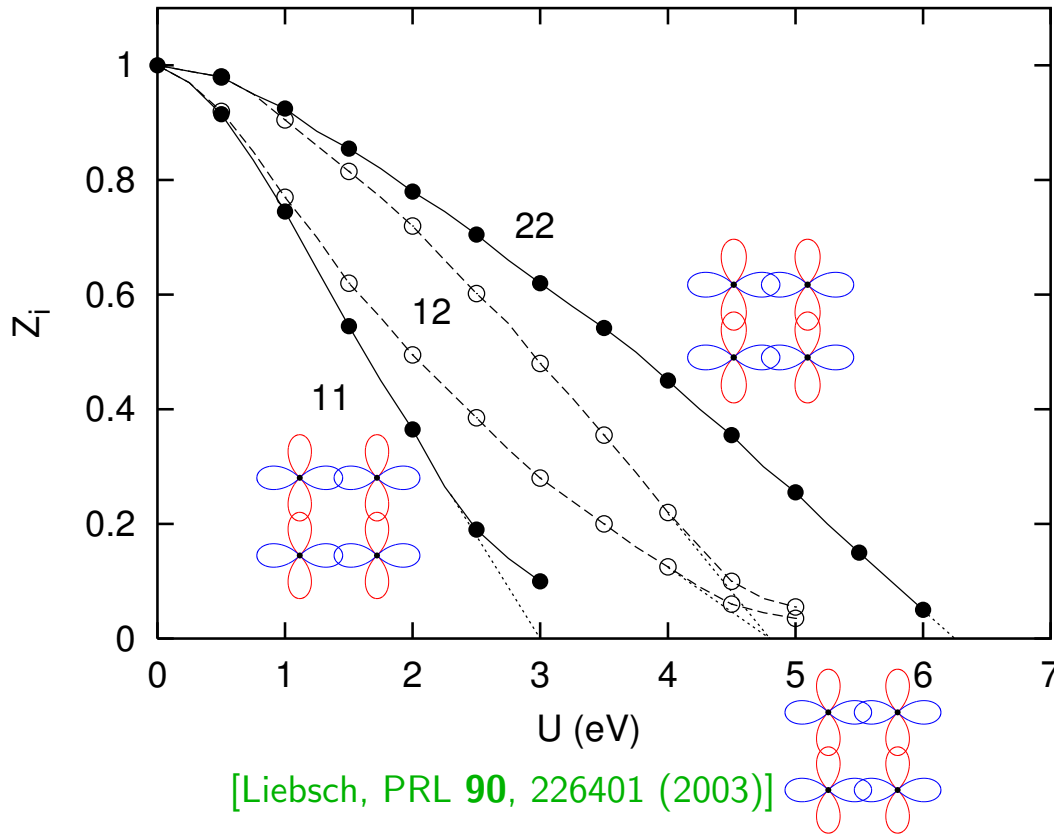
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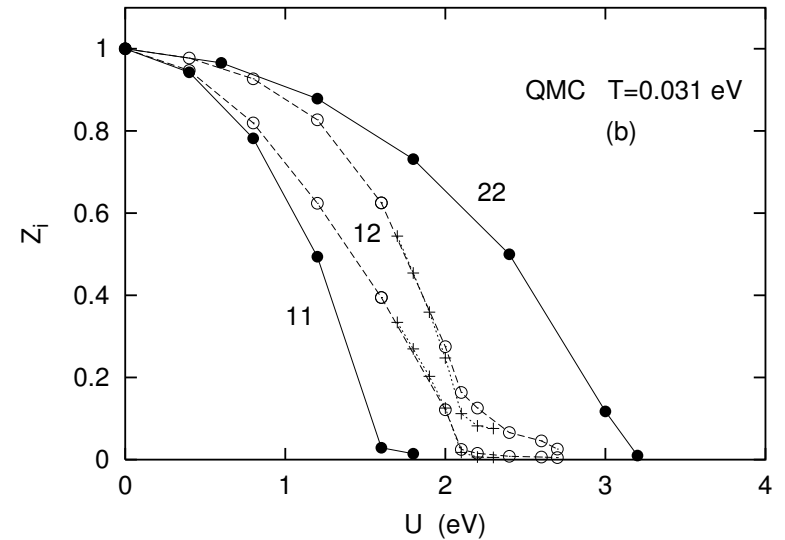
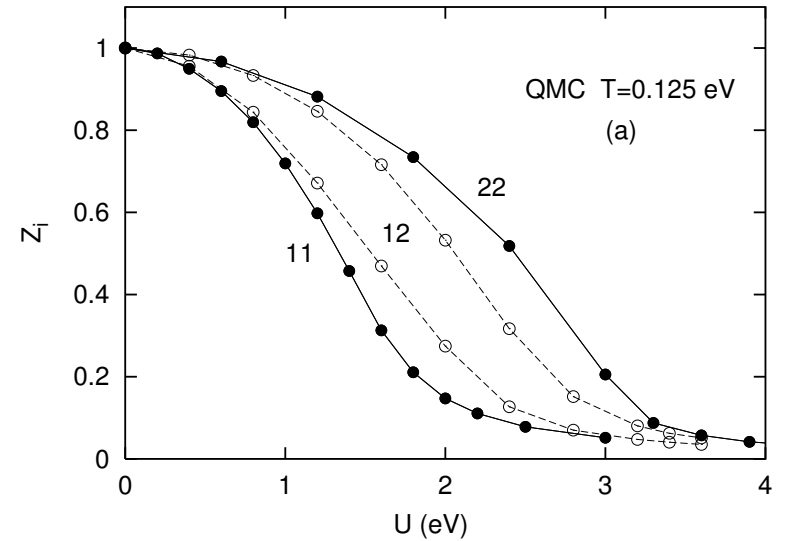


$$J_z = U/4, U' = U/2$$

$$J_z = 0.2, U' = U - 0.4$$

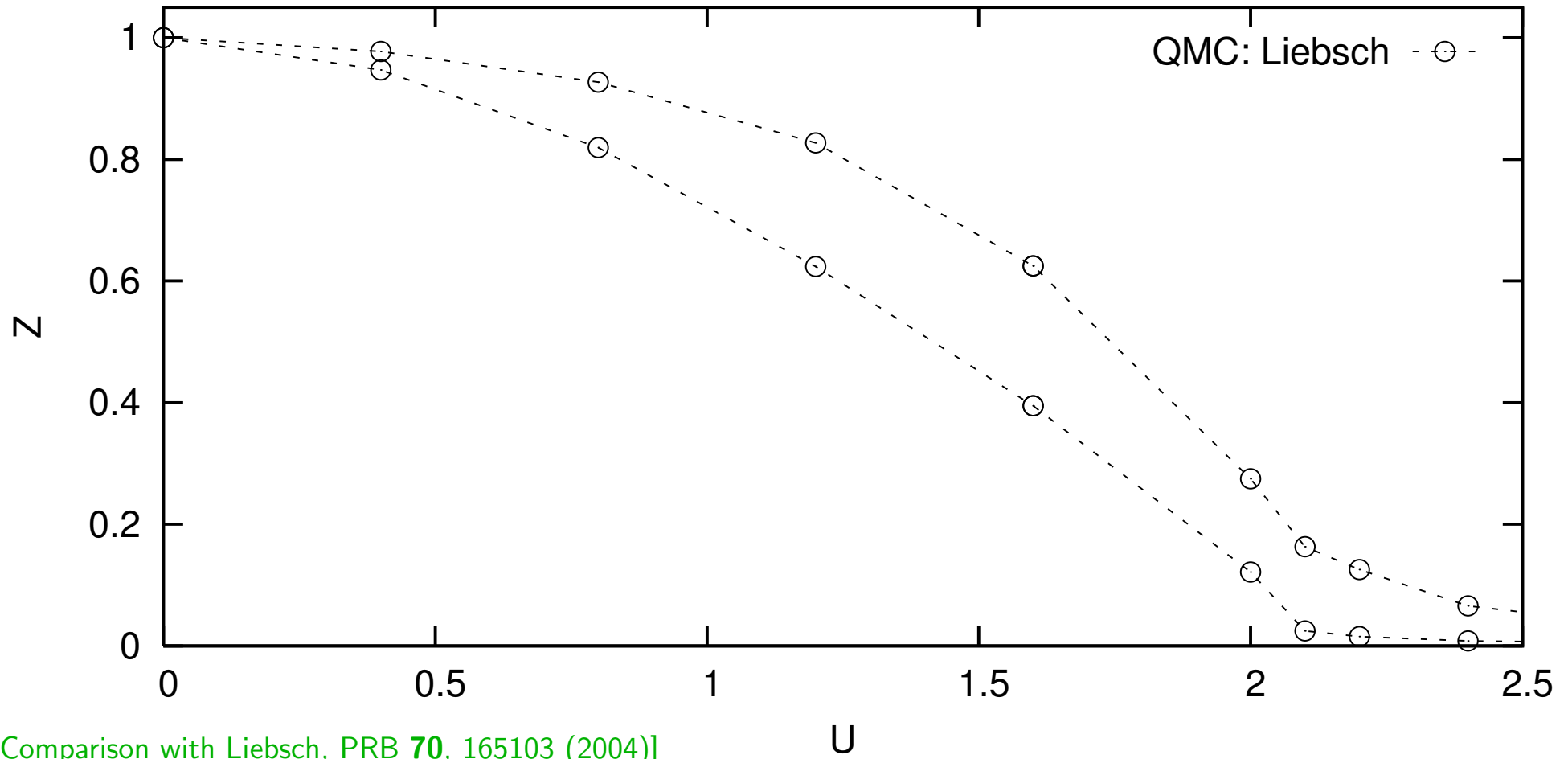


[Liebsch, PRL **90**, 226401 (2003)]



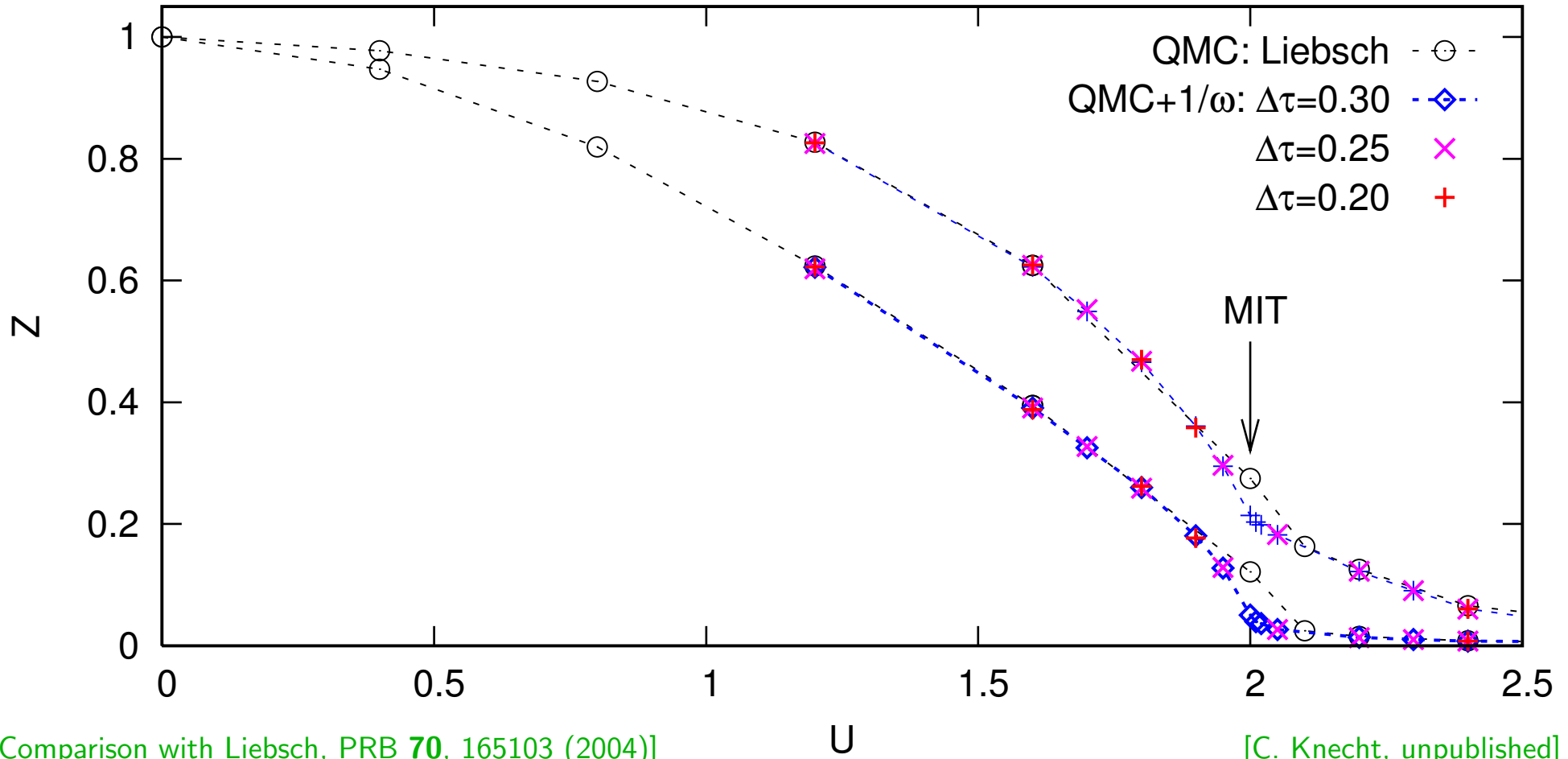
[Liebsch, PRB **70**, 165103 (2004)]

Test for multiband-QMC: quasiparticle weights $Z = m/m^*$ in 2-band model



[Comparison with Liebsch, PRB **70**, 165103 (2004)]

Test for multiband-QMC: quasiparticle weights $Z = m/m^*$ in 2-band model

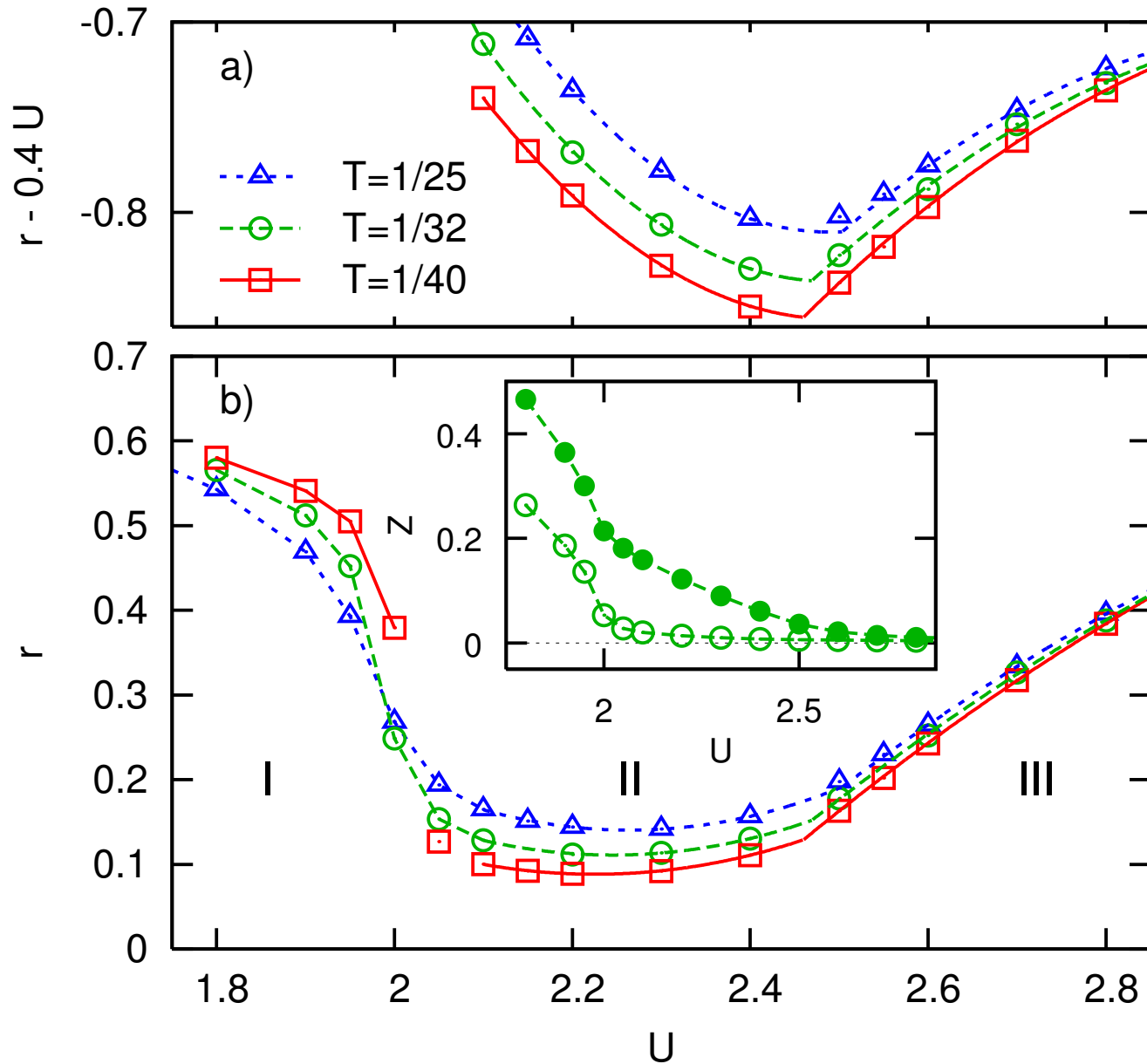


Very small dependence on discretization $\Delta\tau$ (here for $T = 1/32$).

Conclusion in 3/2005: New algorithm clearly exposes (single) metal-insulator transition (MIT)

But: wide band still "quite metallic" for $U > 2.0$ - 2nd transition?

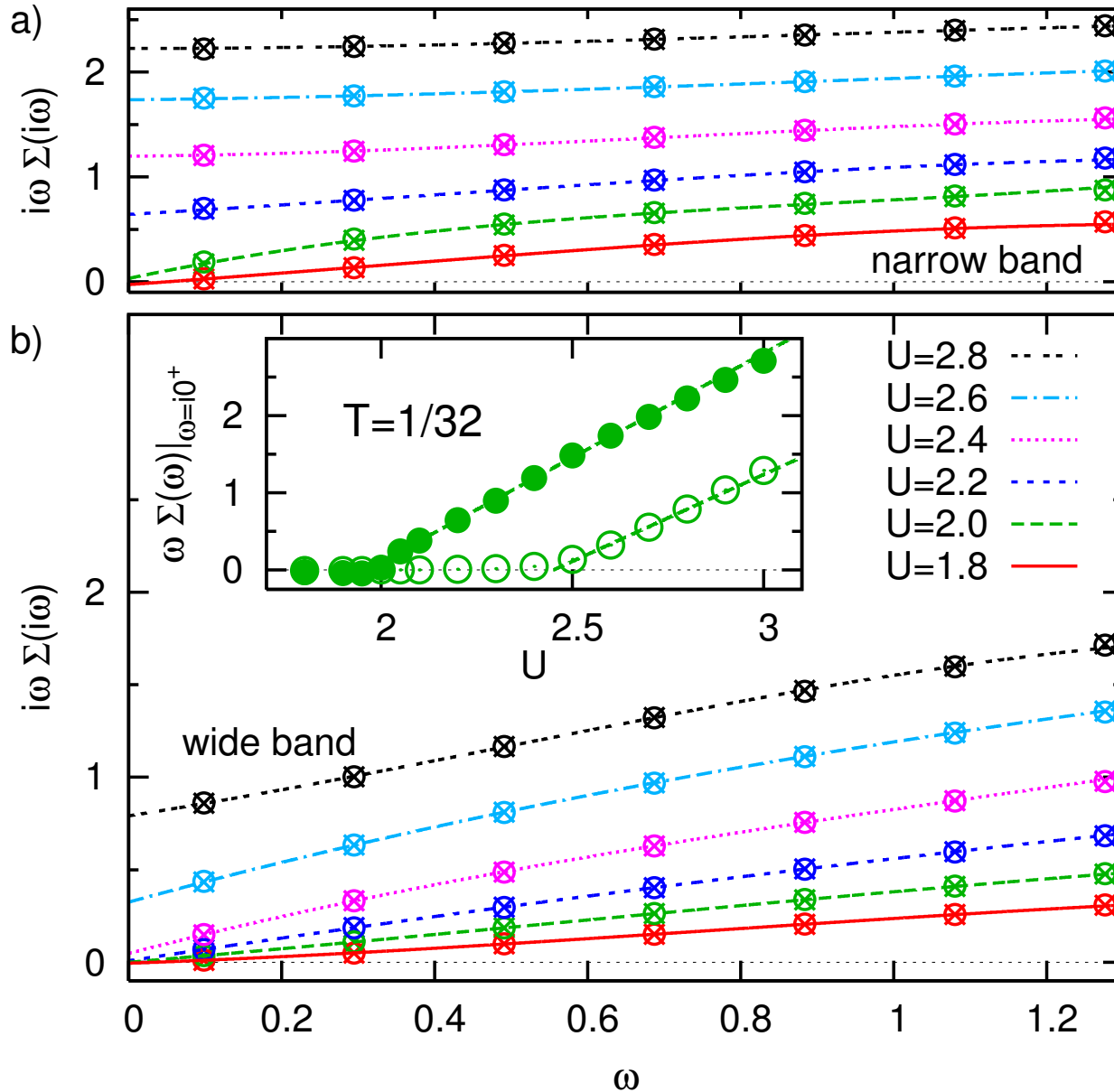
Ratio of quasiparticle weights $r = Z_{\text{narrow}}/Z_{\text{wide}}$



3 regions of different character

kinks indicate 2nd transition at $U \approx 2.5$

Low-frequency analysis of self-energy



for regular self-energy:

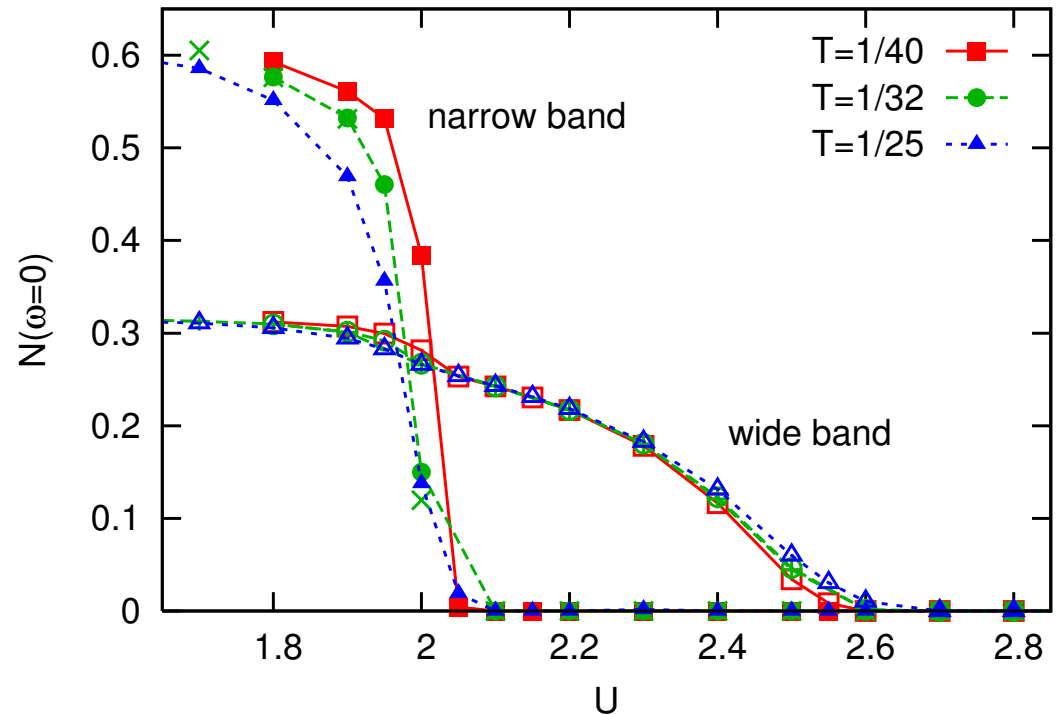
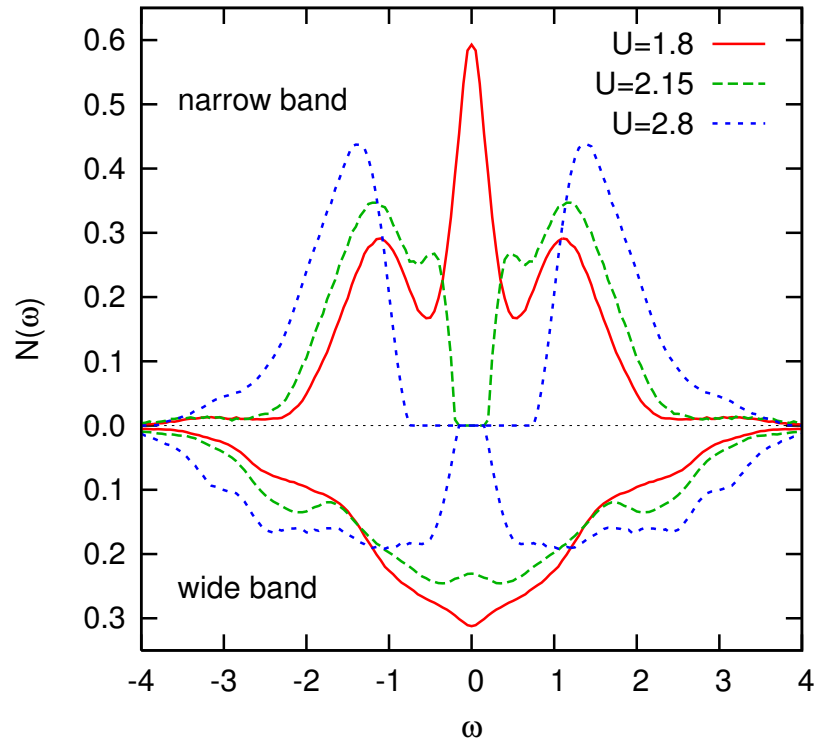
$$\omega \Sigma(\omega) \xrightarrow{\omega \rightarrow 0} 0$$

singularities (\sim gap) for

$U \gtrsim 2$ in narrow band

$U \gtrsim 2.5$ in wide band

Spectral function (interacting DOS)



Clear indications for second singularity

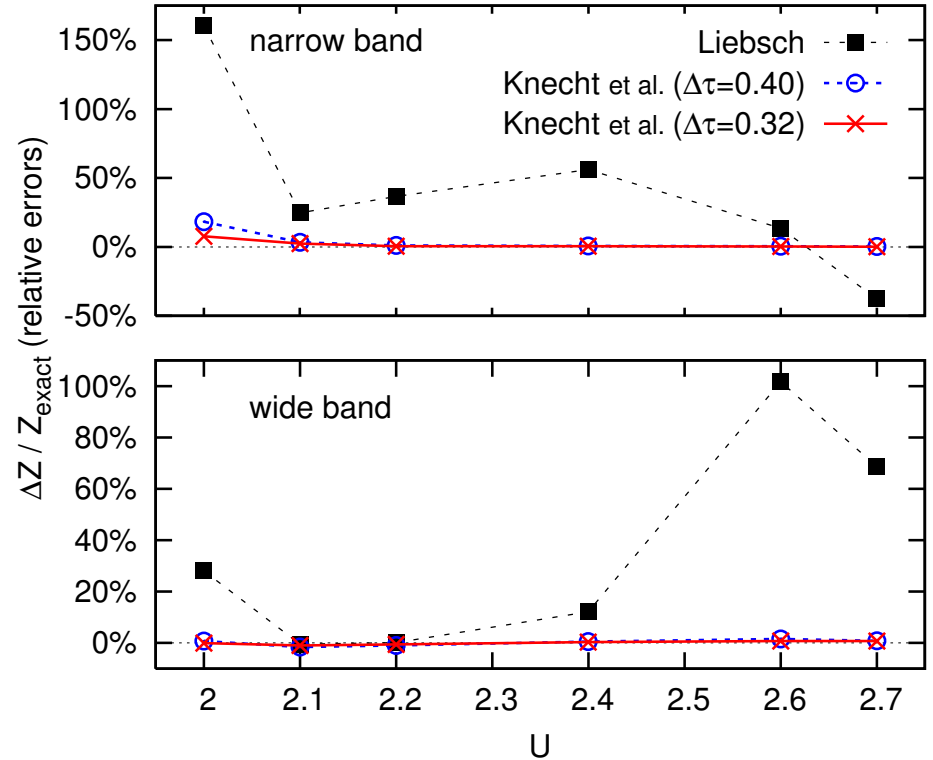
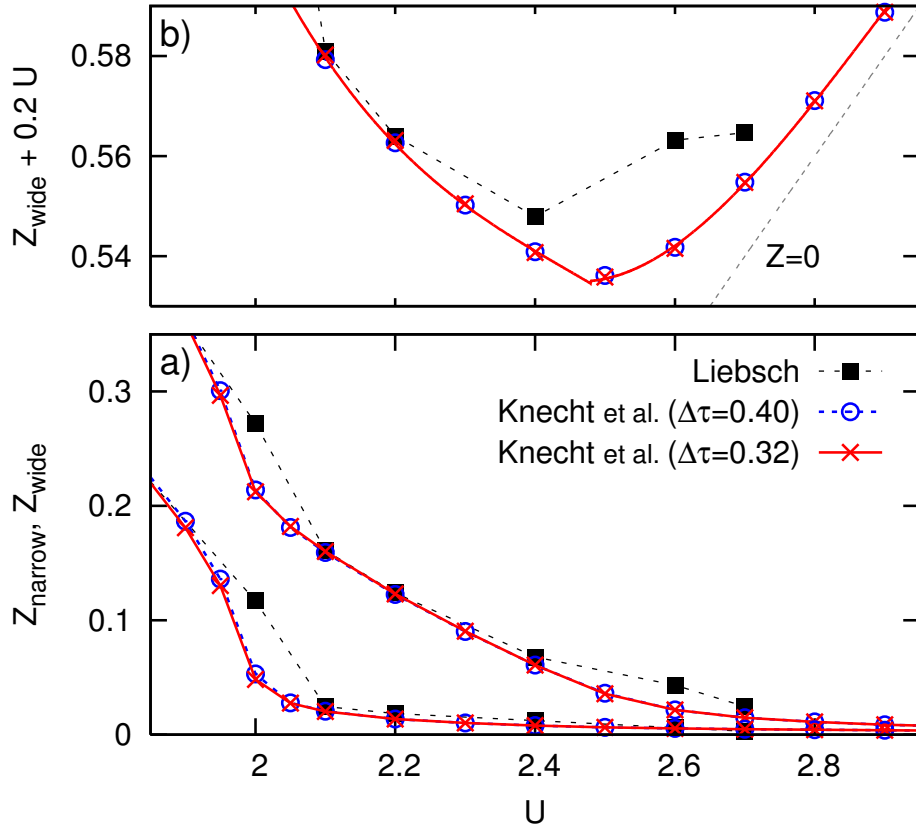
Wide band remains metallic at $U \approx 2.0$

↪ **two orbital-selective Mott transitions** [Knecht, NB, van Dongen, cond-mat/0505106, to appear in PRB RC]

same conclusions from slave-spin approximation [de' Medici, Georges, Biermann, cond-mat/0503764]

Comparison at $T = 1/32$ with Liebsch, PRB **70**, 165103 (2004)

triggered by Comment [\[Liebsch, cond-mat/0506138\]](#) on our preprint

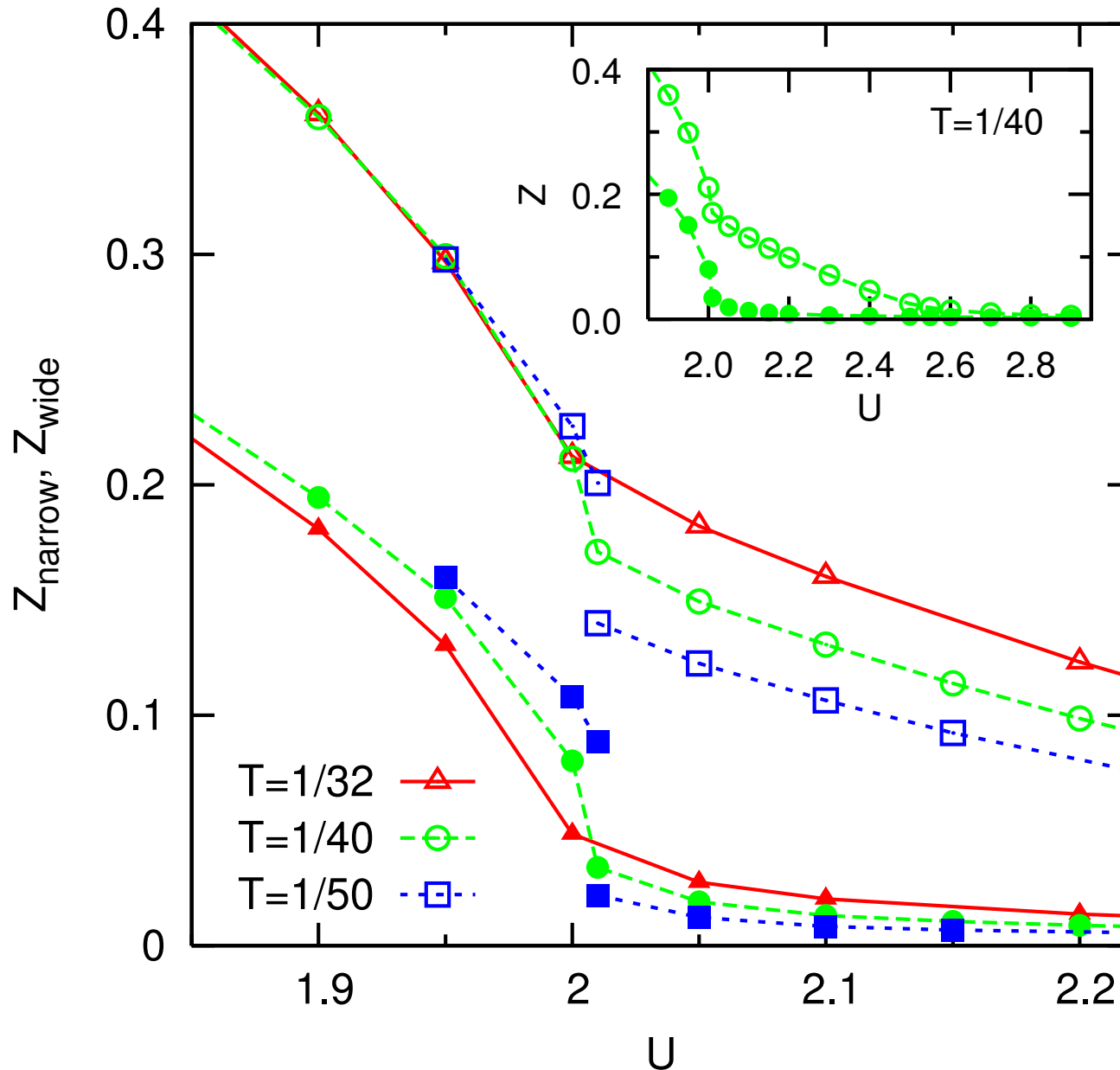


Numerical noise in Liebsch's QMC data obscures second transition

Liebsch's relative errors $> 100\%$ at both transitions [our error: $\mathcal{O}(1\%)$]

[\[Knecht, NB, van Dongen, cond-mat/0506450\]](#)

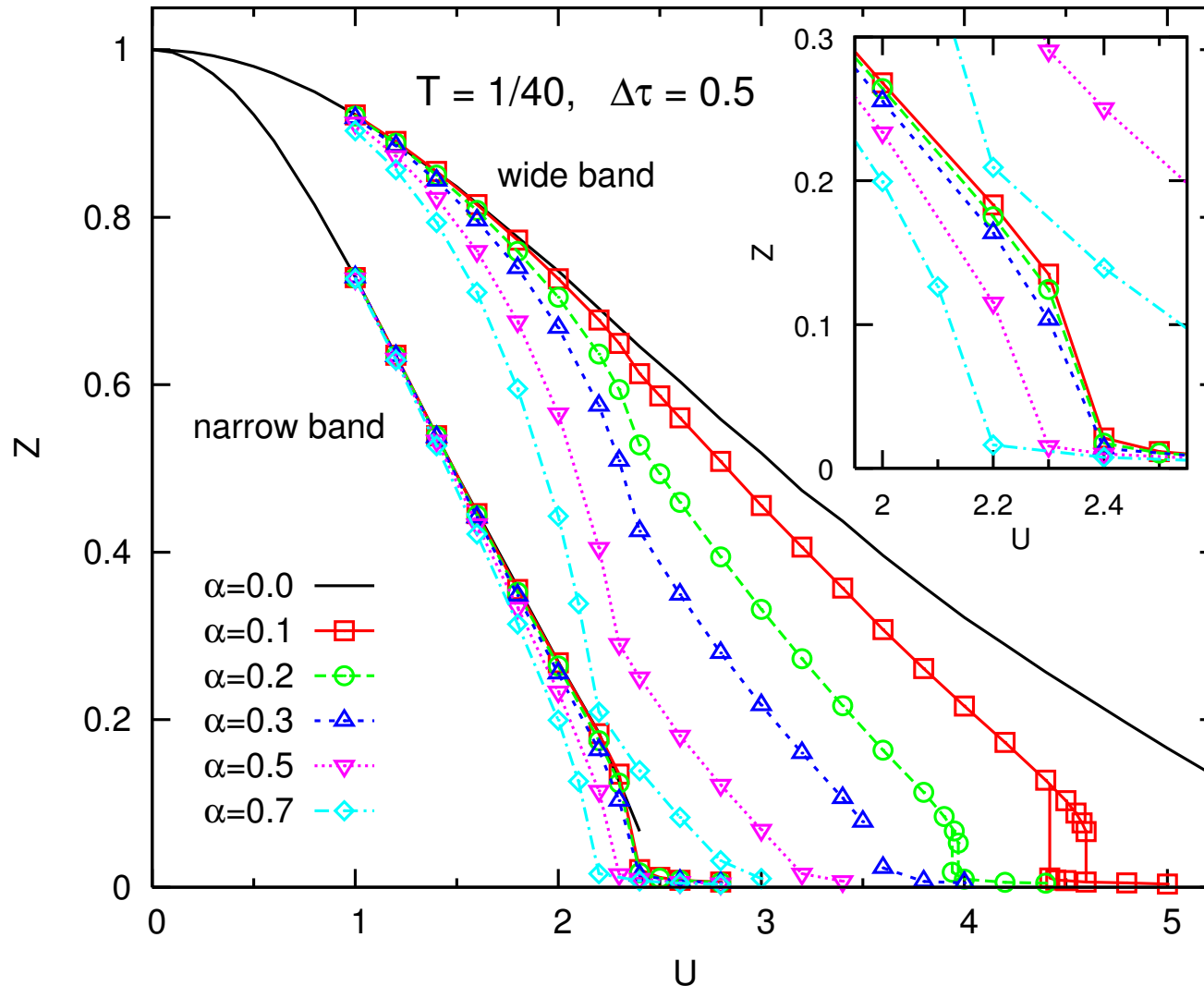
Determination of critical temperature for narrow-band transition



Coexistence only for
 $T \lesssim 0.02 = 1/50$
[van Dongen, Knecht, NB,
cond-mat/0507682]

Systematic study: effect of inter-orbital coupling (preliminary results)

$$H = \sum_{m=1}^2 \left[- \sum_{\langle ij \rangle \sigma} t_m c_{im\sigma}^\dagger c_{jm\sigma} + U \sum_i n_{im\uparrow} n_{im\downarrow} \right] + \alpha \sum_{i\sigma\sigma'} (U/2 - \delta_{\sigma\sigma'} U/4) n_{i1\sigma} n_{i2\sigma'}$$



Both orbital-selective Mott transitions remain first order (at least) for small α

Summary

Improved DMFT-QMC scheme using high-frequency expansion of $\Sigma(\omega)$

Mott transition in frustrated 1-band Hubbard model

High-precision ground state estimates from QMC

Critical exponents from (infinite-order) ePT

Orbital-selective Mott transition in 2-band Hubbard model

J_z model: minimal model for OSMTs in $\text{Ca}_{2-x}\text{Sr}_x\text{RuO}_4$

narrow-band transition 1st order for $T \lesssim 0.02$

wide-band transition 1st order for small J_z, U'

Thanks to: D. Vollhardt, E. Kalinowski, C. Knecht, P. van Dongen

NIC Jülich, DFG (BI775/1)