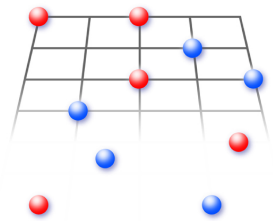


Double occupancy as a universal probe for antiferromagnetic correlations and entropy in cold fermions on optical lattices

Nils Blümer and Elena Gorelik
Gutenberg University Mainz, Germany

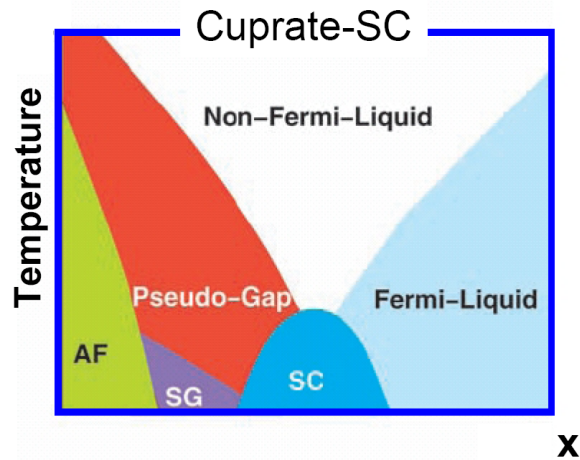


Transregional Collaborative Research Centre SFB / TRR 49
Condensed matter systems with variable many-body interactions
Frankfurt / Kaiserslautern / Mainz

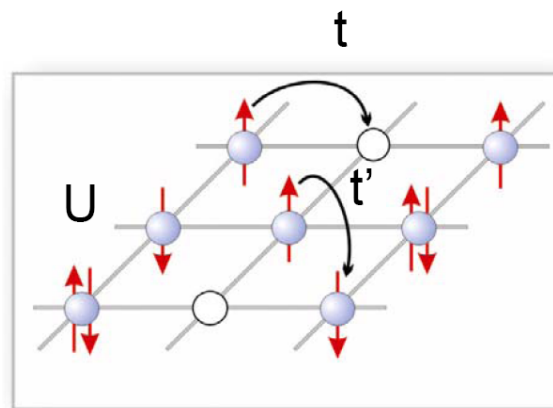
JOHANNES
GUTENBERG
UNIVERSITÄT
MAINZ

Motivation: Ultracold lattice fermions as quantum simulators?

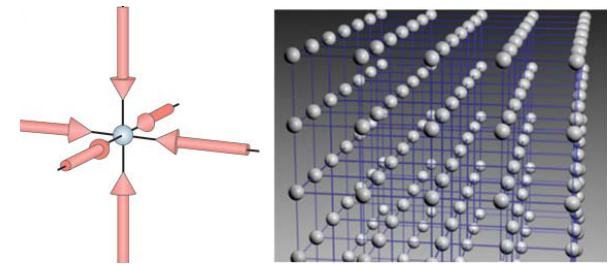
Correlated materials



Fermionic Hubbard model



Ultracold fermions on optical lattices

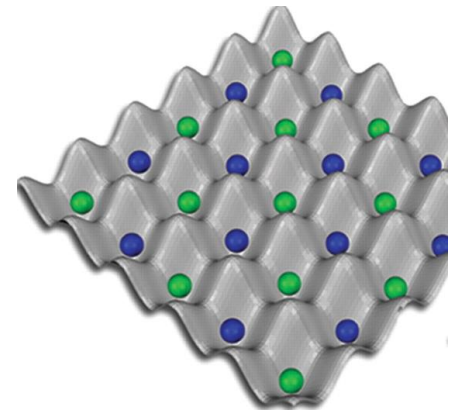


Recent breakthrough: paramagnetic Mott transition in 2-flavor mixtures
[Schneider et al, Science **322**, 1520 (2008), Jördens et al., Nature **455**, 204 (2008)]

Remaining challenge: antiferromagnetism (staggered order)

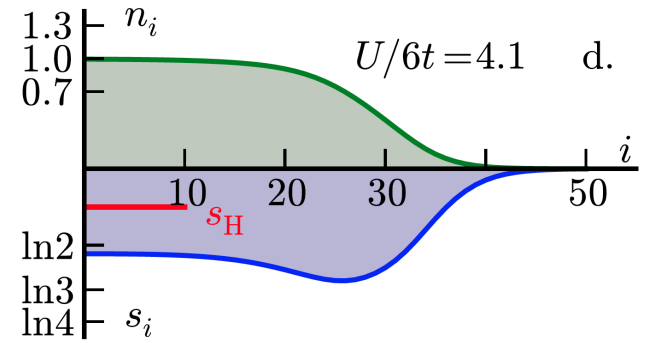
Problems:

- (i) difficult to reach sufficiently low temperatures/entropies
- (ii) detection of AF order is not straightforward
- (iii) inhomogeneity, time scale for global (spin) equilibrium



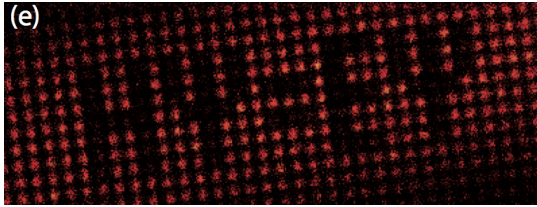
Questions for this talk

- How to detect AF order/correlations?
- Which entropy range is needed?
- How reliable are DMFT and LDA?
- General impact of dimensionality?

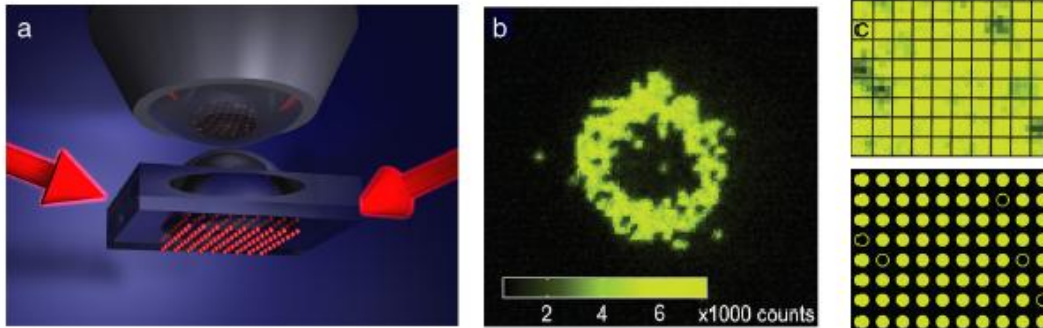


[Jördens et al., PRL **104**, 180401 (2010)]

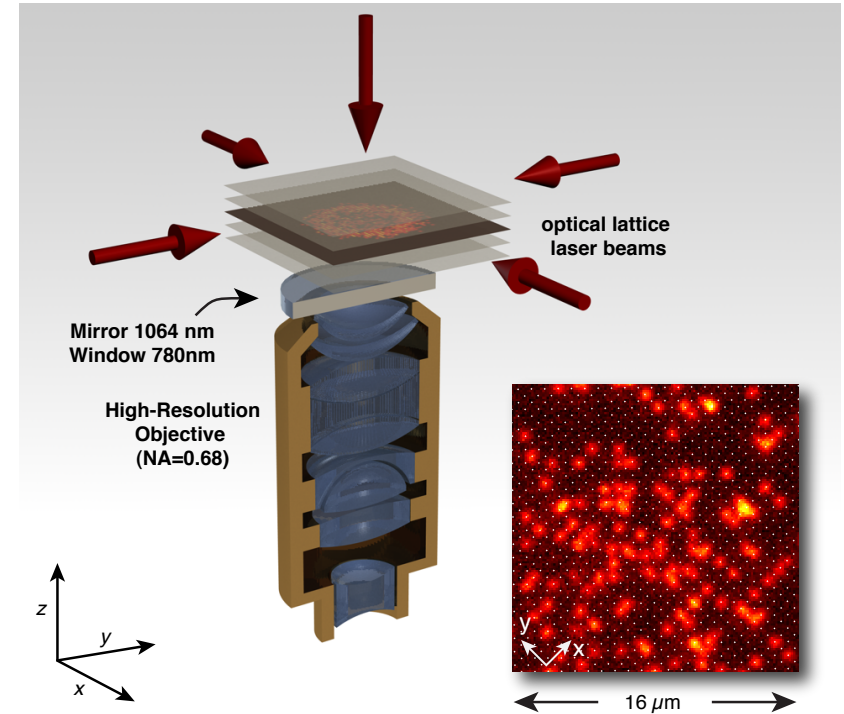
Experimental advantage of 2 dimensions:
single-site resolution (for bosons)



[Würtz et al., PRL **103**, 080404 (2009)]



[Bakr et al., Science **329**, 547 (2010)]



[Sherson et al., Nature **467**, 68 (2010)]

Outline

Motivation: Ultracold lattice fermions as quantum simulators?

Methods: (Real-space) DMFT, slab approximation, QMC

[N. Blümer and E. V. Gorelik, *Computer Physics Communications* **118**, 115 (2011)]

Results: Néel transition of lattice fermions in a harmonic trap

[E. V. Gorelik, I. Titvinidze, W. Hofstetter, M. Snoek, N. Blümer, *PRL* **105**, 065301 (2010)]

Effect of nonlocal correlations? Comparisons with direct QMC + BA

[E. V. Gorelik, T. Paiva, R. Scalettar, A. Klümper, N. Blümer, *arXiv:1105.3356*]

Summary and outlook

Collaborators

Postdoc



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Univ. Mainz



Walter Hofstetter
Univ. Frankfurt



Irakli Titvinidze
Univ. Hamburg



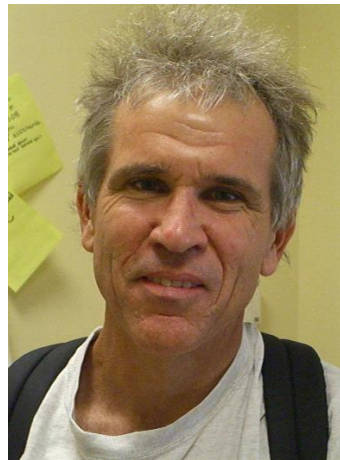
Michiel Snoek
Univ. Amsterdam



Andreas Klümper
Univ. Wuppertal



Thereza Paiva
Rio de Janeiro



Richard Scalettar
UC Davis

Real-space DMFT collaboration

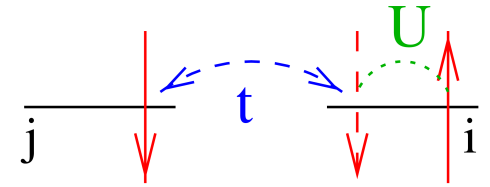
[Gorelik et al., PRL **105**, 065301 (2010)]

AF in finite dimensions

[Gorelik et al., arXiv:1105.3356]

Methods: Approaches for Hubbard-type models

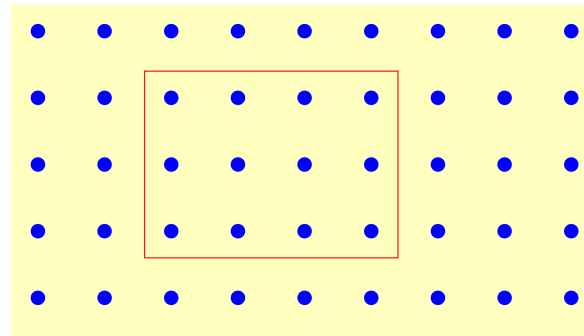
$$\hat{H} = \sum_{(i,j),\sigma} t_{ij} (\hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \text{h.c.}) + U \sum_i \hat{n}_{i\uparrow} \hat{n}_{i\downarrow}$$



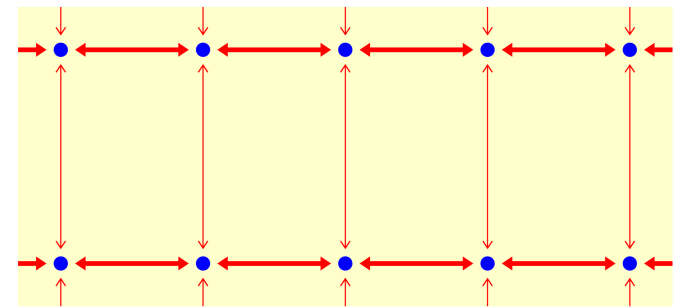
Perturbation theory

- $U \rightarrow 0$: Hartree-Fock
2nd order PT, . . .
- $t/U \rightarrow 0$ (for $n = 1$)
 \rightsquigarrow Heisenberg model

finite clusters: ED, QMC



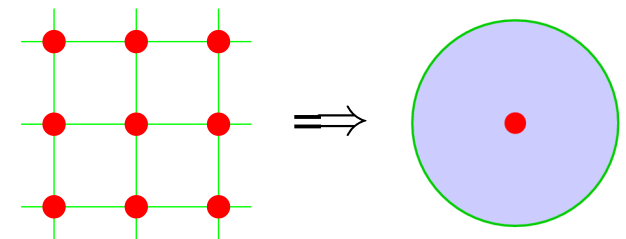
$d \rightarrow 1$: Bethe ansatz, DMRG



Dynamical mean-field theory (DMFT): local self-energy $\Sigma(\mathbf{k}, \omega) \equiv \Sigma(\omega)$

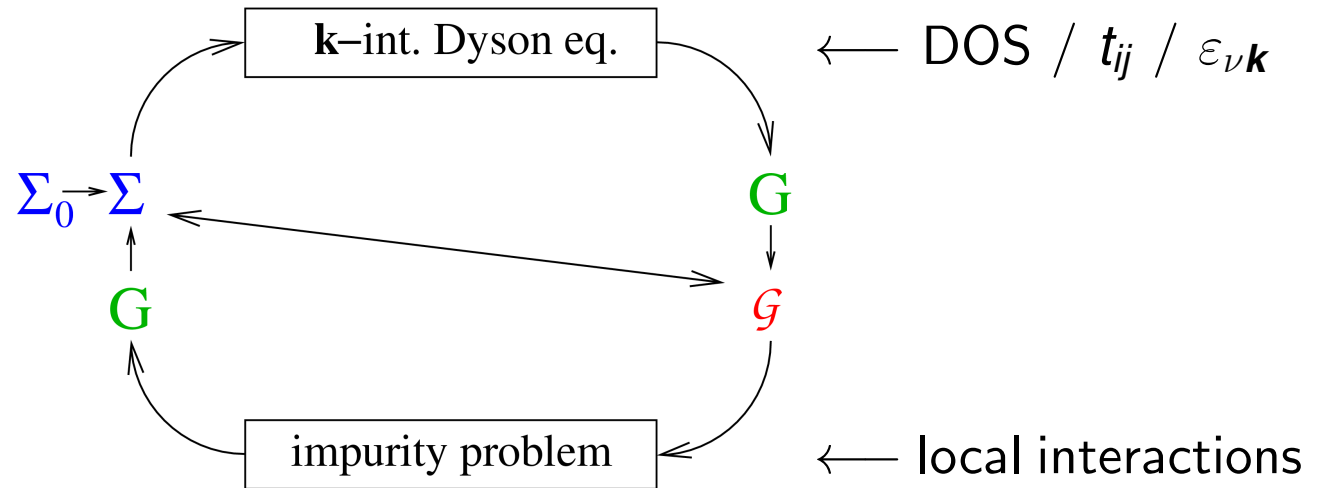
[Metzner, Vollhardt, PRL (1989), Georges, Kotliar, PRL (1992), Jarrell, PRL (1992)]

- + non-perturbative \rightsquigarrow valid at MIT
- + dynamical on-site correlations preserved
- + in thermodynamic limit
- +/- exact for coordination $Z \rightarrow \infty$



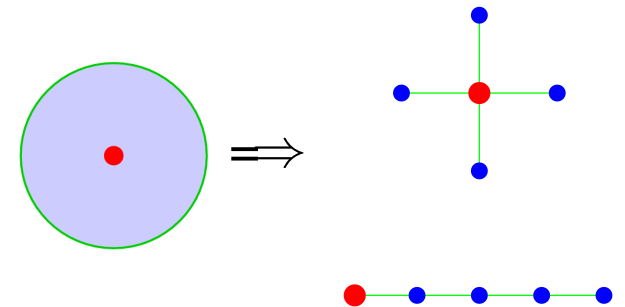
Iterative solution of DMFT equations

0. Initialize self-energy
1. Solve Dyson equation
2. Solve **single impurity Anderson model (SIAM)**



Impurity solver:

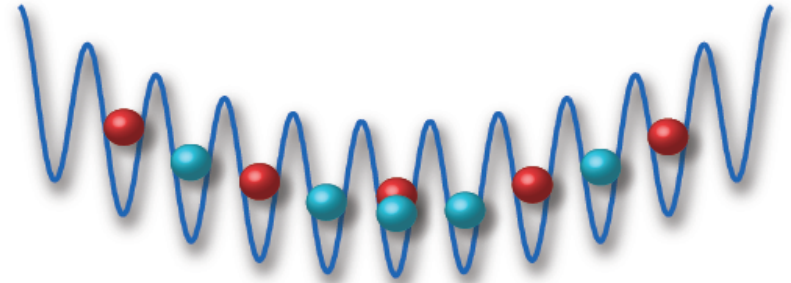
- Iterative perturbation theory (IPT; not controlled)
- Quantum Monte Carlo (QMC)
- Exact diagonalization (ED; large finite-size errors)
- Numerical renormalization group (NRG; 1-2 bands)
- Density matrix renormalization group (DMRG)



Real-space DMFT: use local self-energy in inhomogeneous system

Include **trapping potential**, e.g.: $V_i = V r_i^2$

$$H = - \sum_{(ij),\sigma} t_{ij} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_{i=1}^N n_{i\uparrow} n_{i\downarrow} + \sum_{i,\sigma} V_i n_{i\sigma}$$



\rightsquigarrow N single-site **impurities**, coupled by **real-space lattice Dyson equation**:

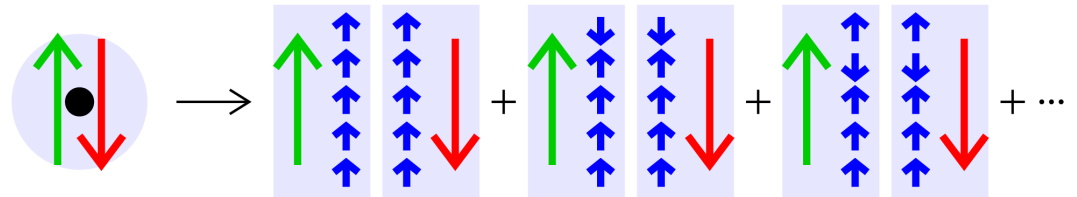
$$\left[G_\sigma(i\omega_n) \right]_{ij}^{-1} = (\mu_\sigma + i\omega_n) \delta_{ij} - t_{ij} - (V_i + \Sigma_{i\sigma}(i\omega_n)) \delta_{ij}$$

[Snoek et al., NJP (2008), Helmes et al., PRL (2008)]

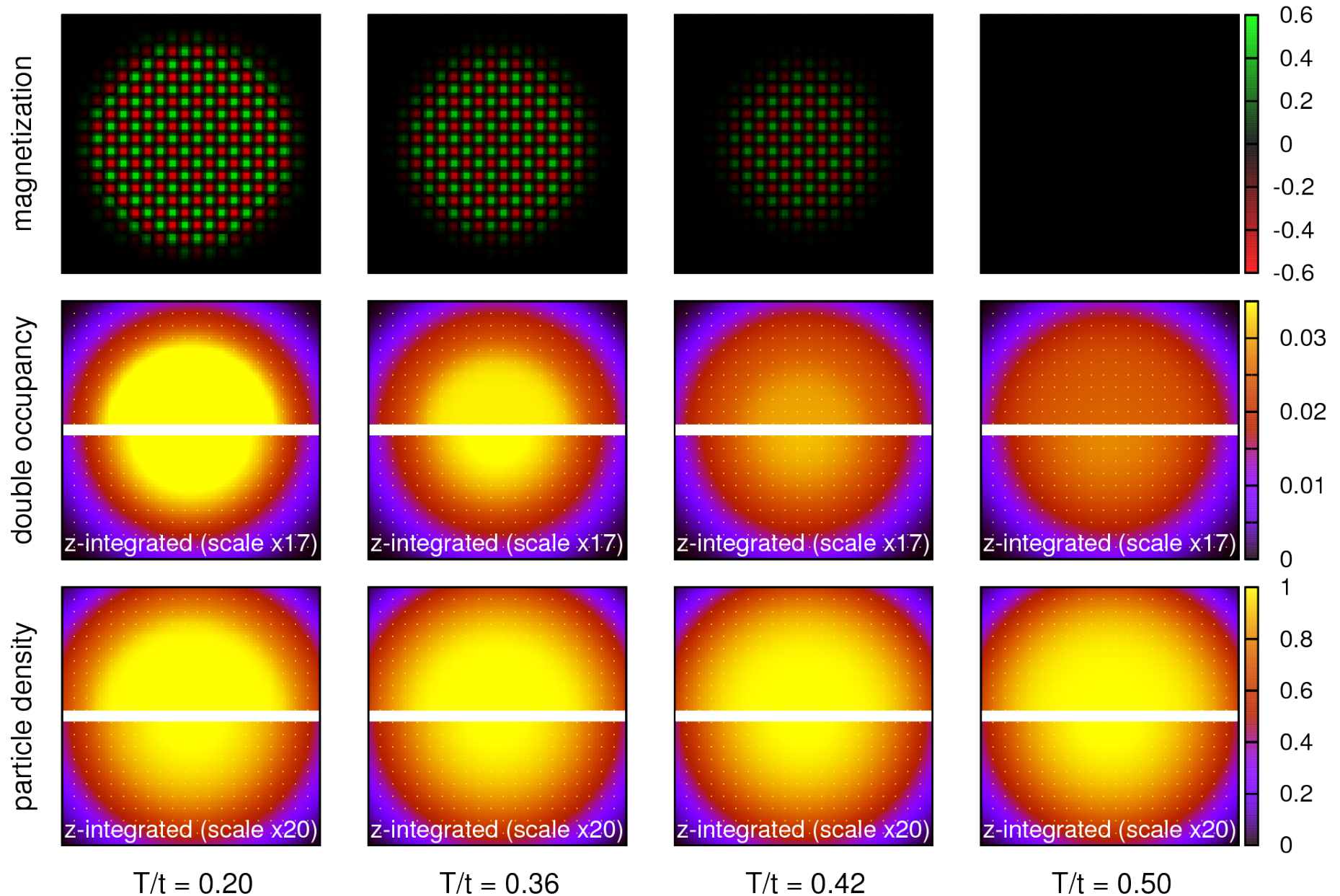
Note: impurity problems are **site-parallel**,
lattice Dyson equation is **frequency-parallel**

Previous implementations: **RDMFT+NRG** (problematic at elevated T)

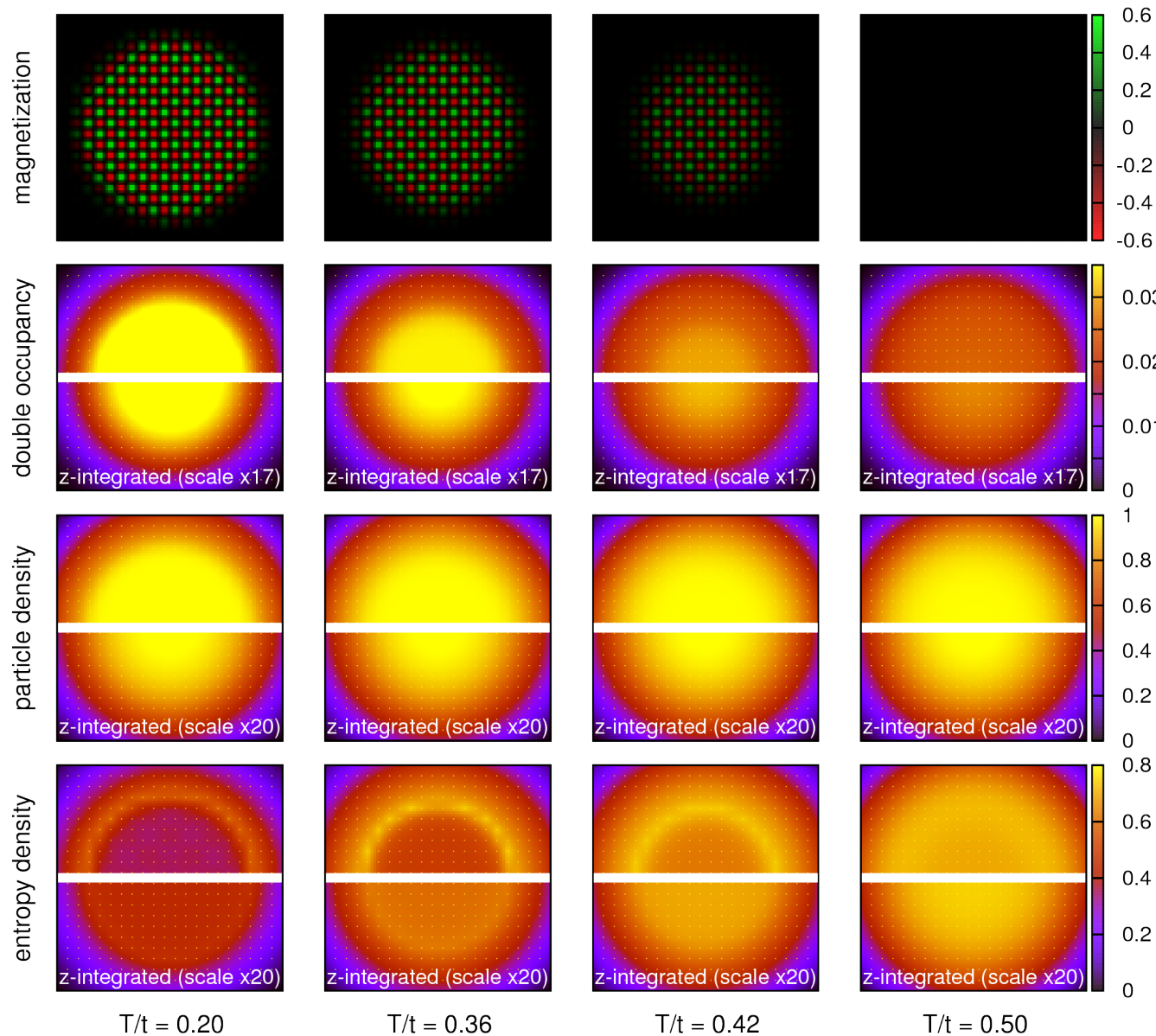
Here: **QMC** (cost $\propto T^{-3}$)
“**slab method**” + pbc
 \sim exact for $\mathcal{O}(10^5)$ atoms



Results: RDMFT-QMC (cubic lattice, $V = 0.05t$, $U = W = 12t$)



Results: RDMFT-QMC (cubic lattice, $V = 0.05t$, $U = W = 12t$)



AF core:

nearly fully polarized at
 $T = 0.20t$

vanishes at $T_N \approx 0.46t$

AF \leftrightarrow enhanced D !

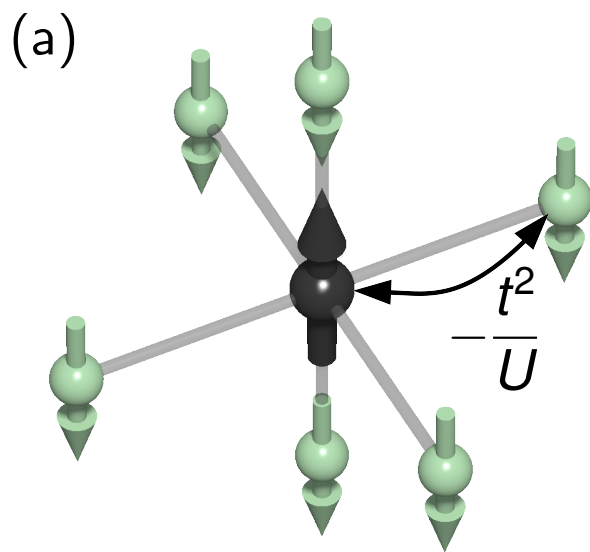
~ 6000 atoms
(naively $\sim 30^3 = 27000$
sites needed)

Entropy

$$S = \int_{-\infty}^0 d\mu' \frac{dN}{dT}$$

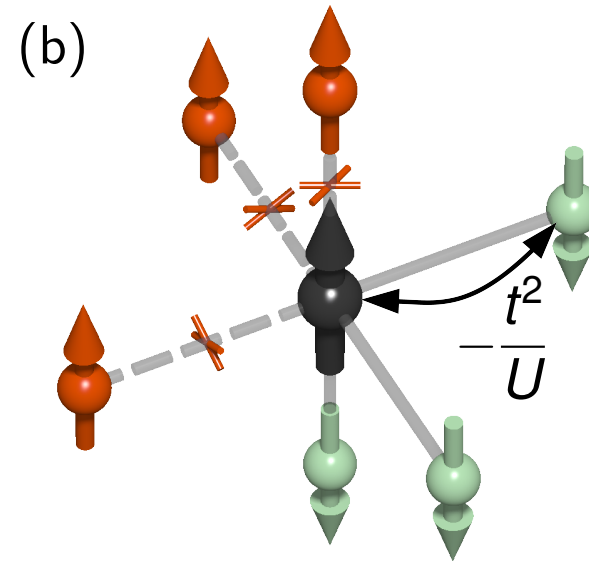
Enhanced double occupancy: a signature of AF order

Illustration of mechanism for enhanced double occupancy (at strong coupling):



AF state: hopping
to all $Z = 6$ neighbors

$$E_{\text{AF}} = -\frac{Zt^2}{U}$$



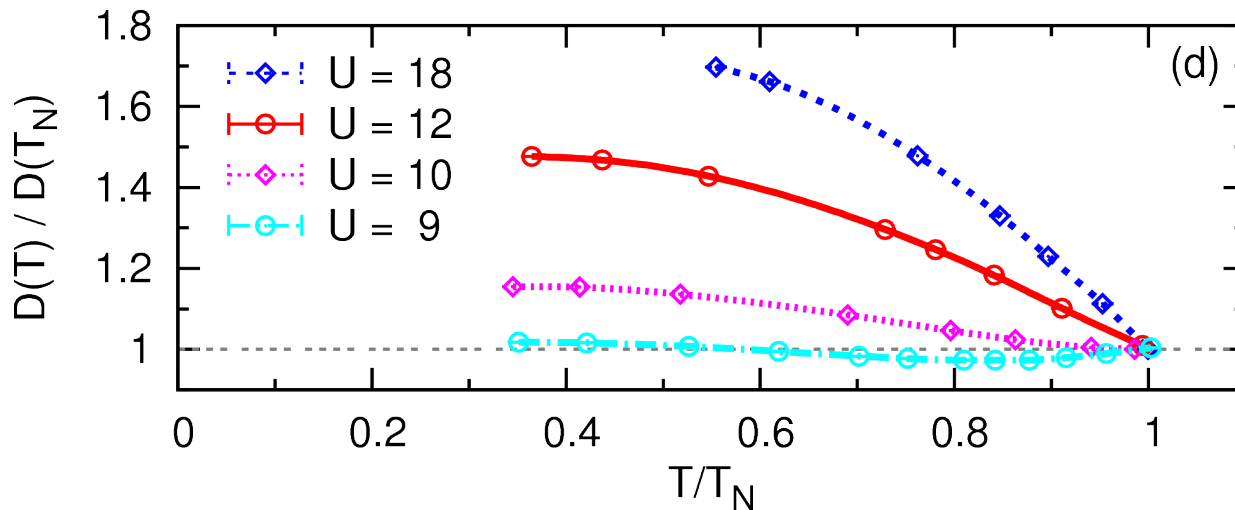
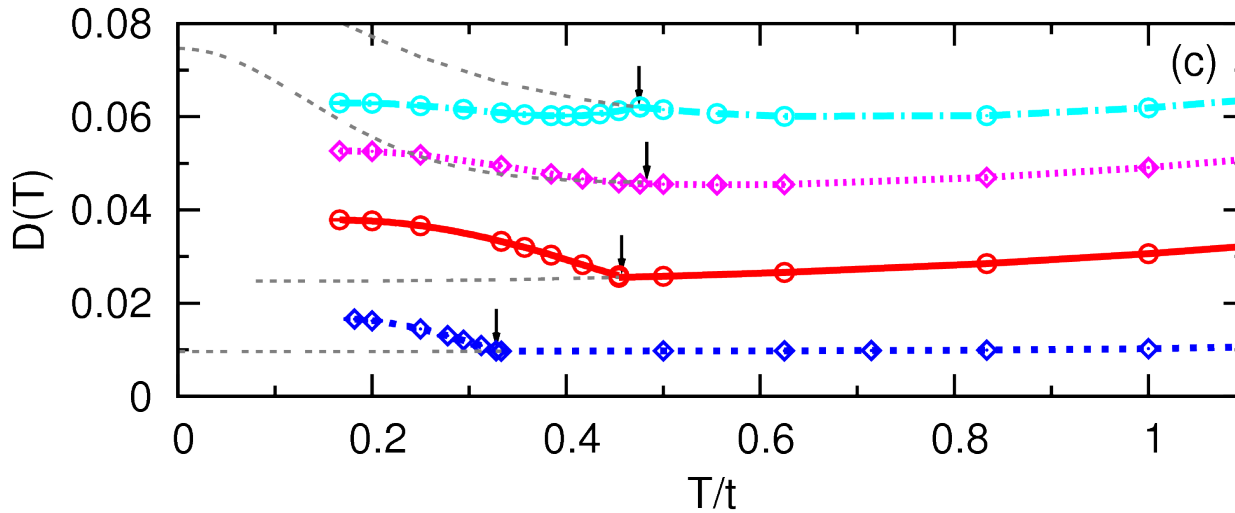
Para/nonmagnetic state:
1/2 of neighbors Pauli forbidden

$$E_{\text{p}} = -\frac{Zt^2}{2U}$$

By $D = dE/dU$ (at $T = 0$), the argument implies $D_{\text{AF}}/D_{\text{p}} \xrightarrow{U \rightarrow \infty} 2$.

Exact relation for all d [Takahashi, 1977]: $E_0 = -\frac{Zt^2}{2U} (1 - \langle \sigma_i \cdot \sigma_j \rangle) + \mathcal{O}\left(\frac{t^4}{U^3}\right)$

DMFT-QMC estimates of D at half filling



AF \Rightarrow

enhanced D at $U \gtrsim 10t$

arrows: Néel temperatures

thin lines: metastable paramagnetic phase.

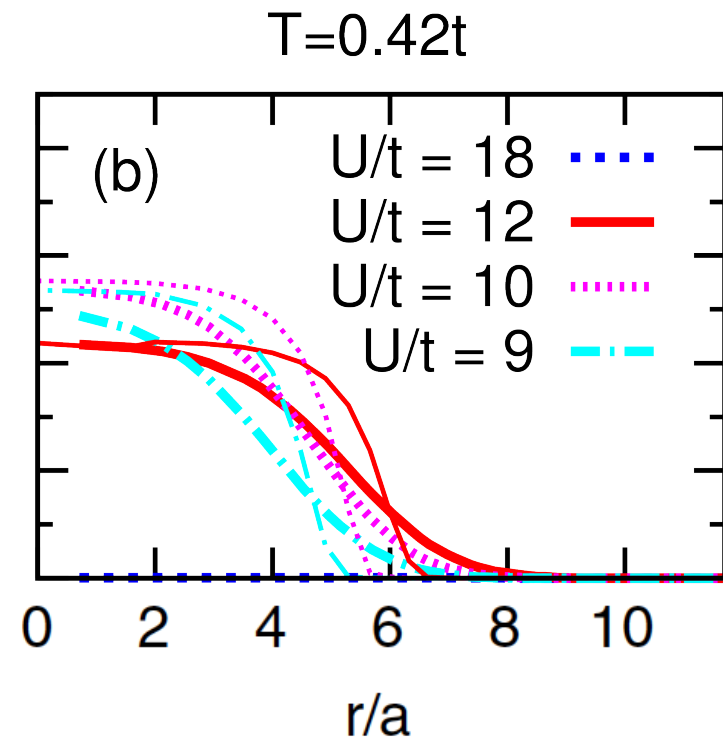
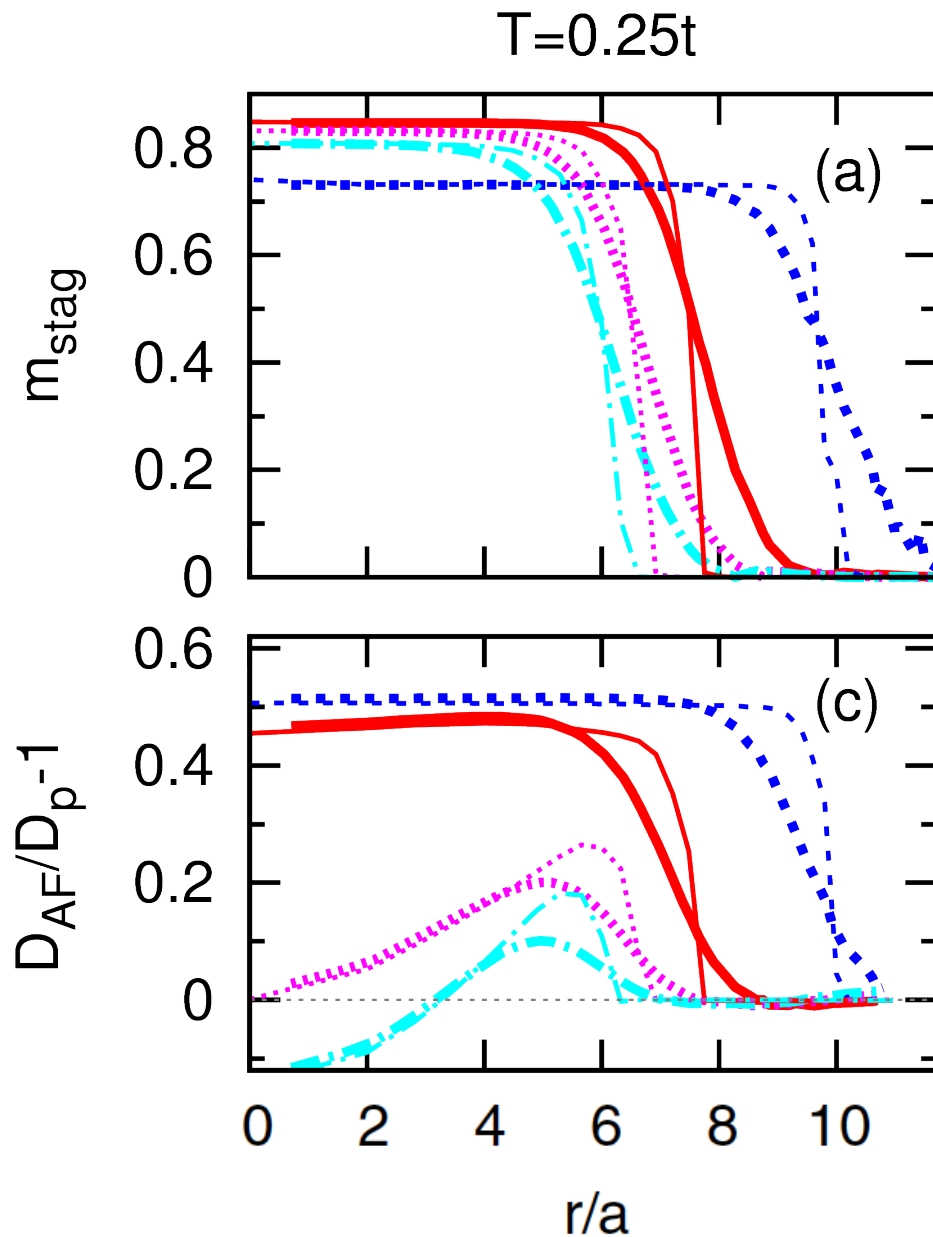
Data scaled to values of critical point:

relative enhancement

$$D/D(T_N) \xrightarrow{U \rightarrow \infty} 2$$

Note: AF kills Pomeranchuk cooling [Werner, Parcollet, Georges, Hassan, PRL (2005)]!

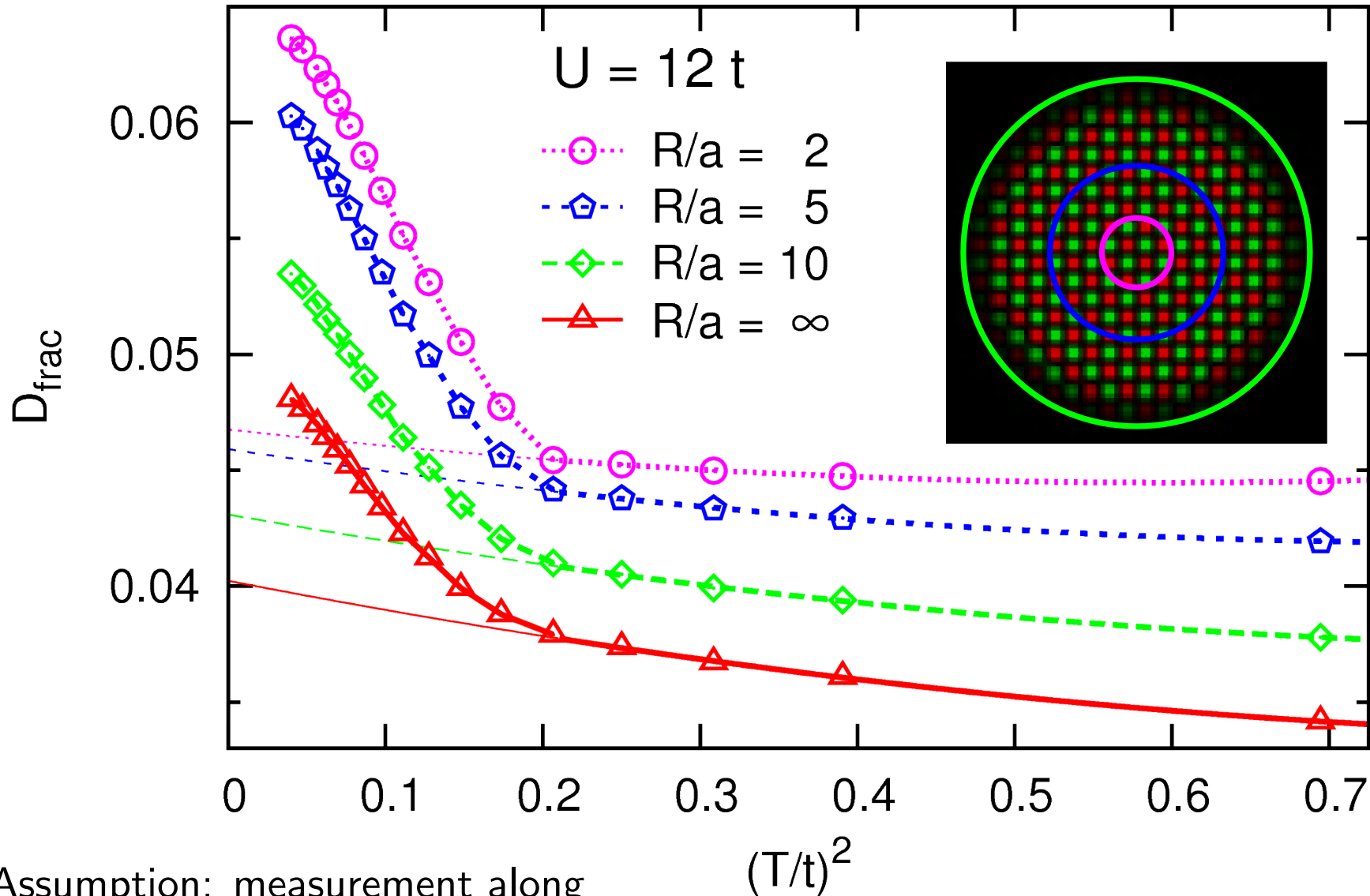
Radial dependence of m_{stag} and D : RDMFT calculations ($V = 0.05t$)



Strong proximity effects
beyond LDA (thin lines)

significant enhancement of D
only at strong coupling

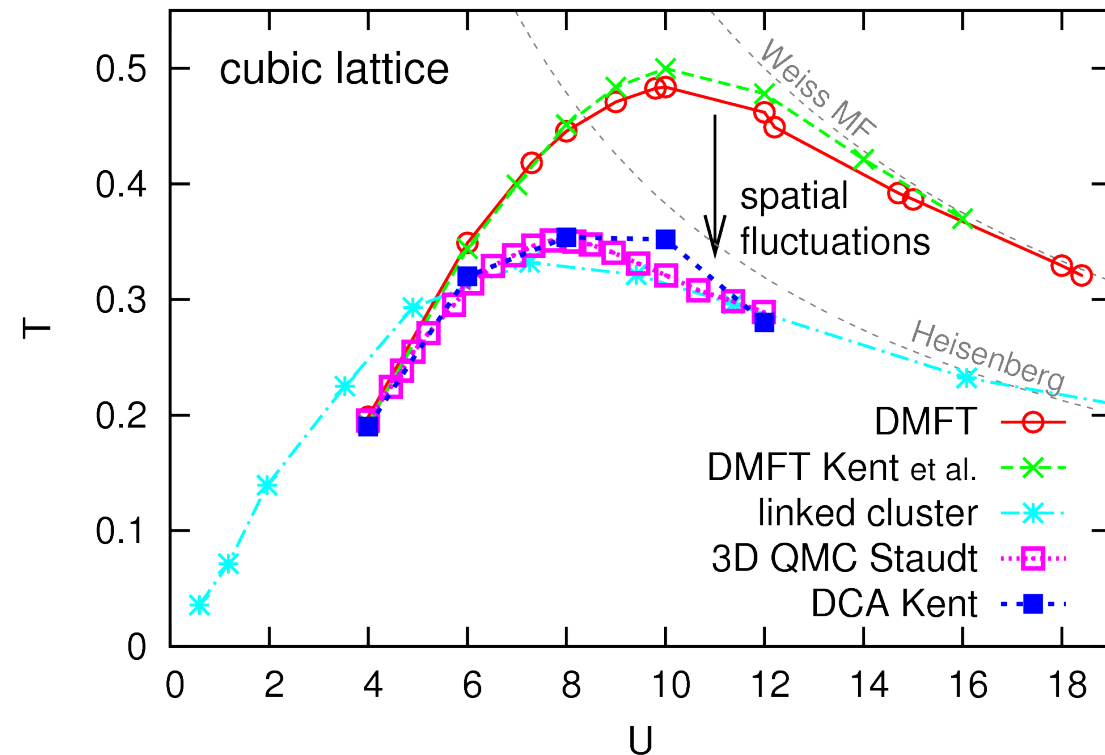
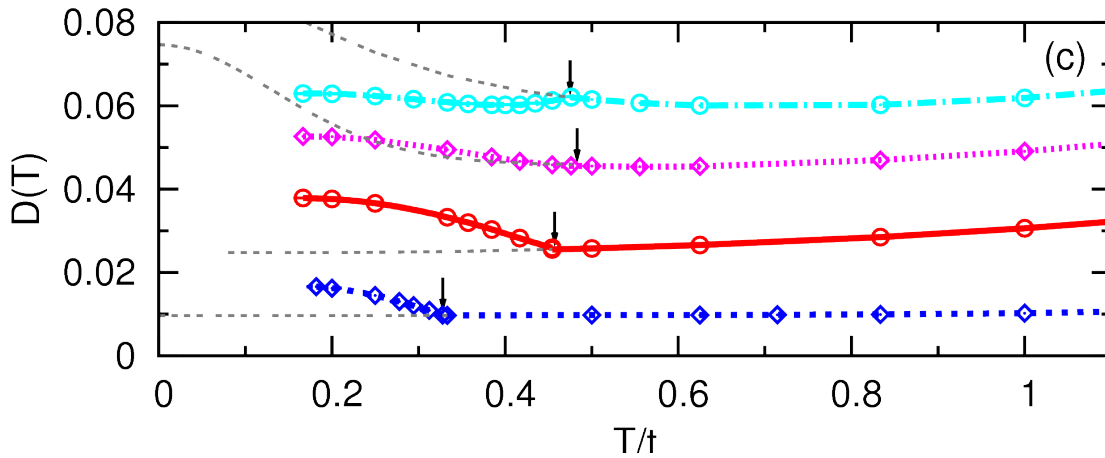
Néel transition visible in integrated quantities? Yes!



Assumption: measurement along beam with gaussian profile

but: effects of nonlocal correlations?

Modification of DMFT predictions by spatial fluctuations in 3d: how?

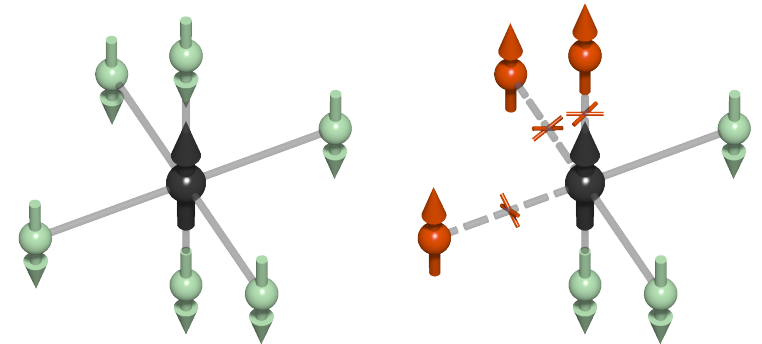


Unavoidable change: kinks cannot remain at $T = T_N^{\text{DMFT}} > T_N!$

Constraints:

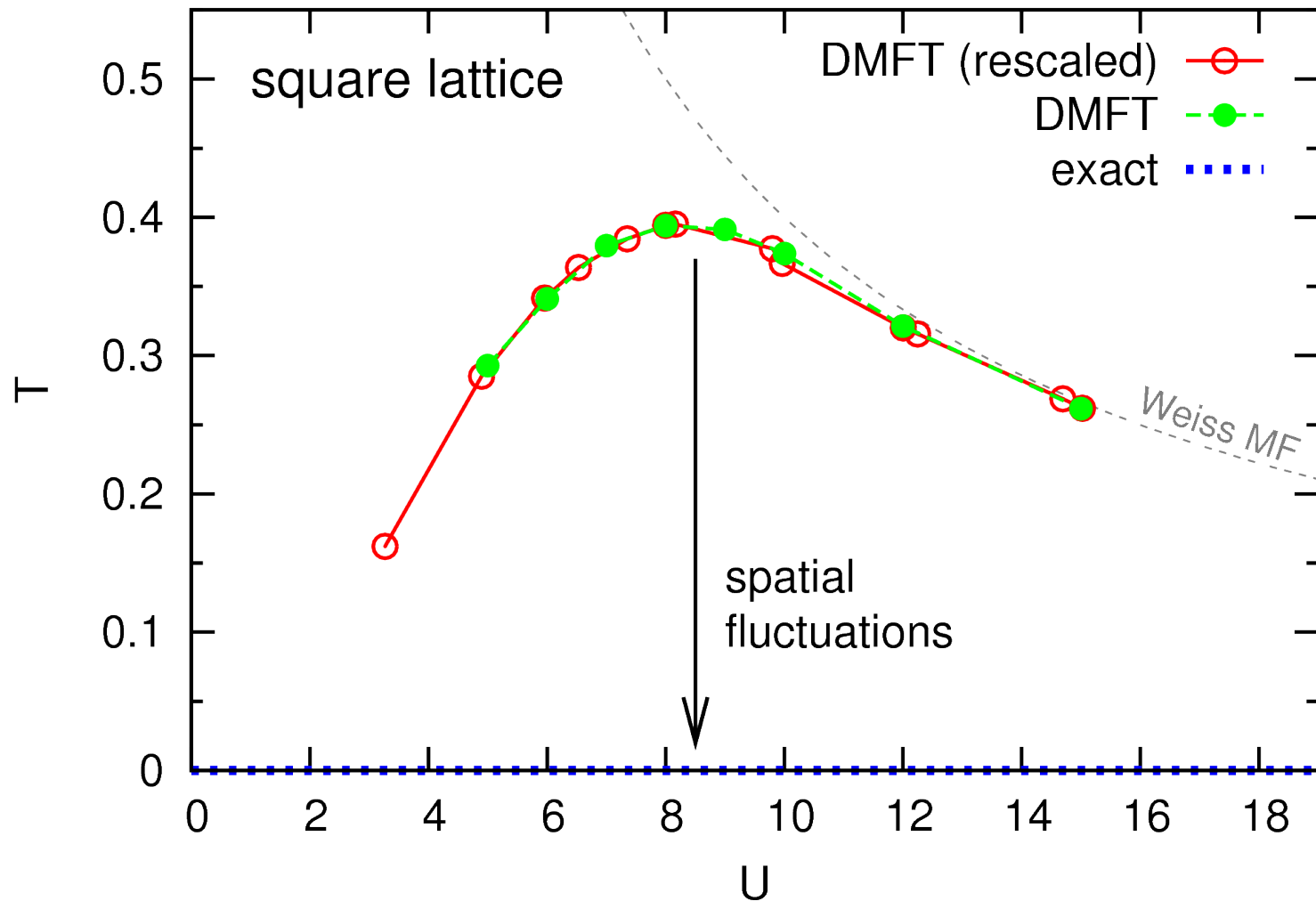
- DMFT results for $D(T)$ agree with high- T expansion at $T \gg T_N$ [Jördens et al., PRL (2010)]

- argument for $D_{\text{AF}}/D_{\text{para}} \xrightarrow{U \rightarrow \infty} 2$ is not explicitly d -dependent



and independ. of long-range order

Situation “worse” in 2d: no antiferromagnetism at finite T !



Will any enhancement of D at low T remain? At which temperature scale?

How large are the DMFT errors in $D(T)$ for $T \gtrsim T_N^{\text{DMFT}}$?

Fermions in 2D Optical Lattices: Temperature and Entropy Scales for Observing Antiferromagnetism and Superfluidity

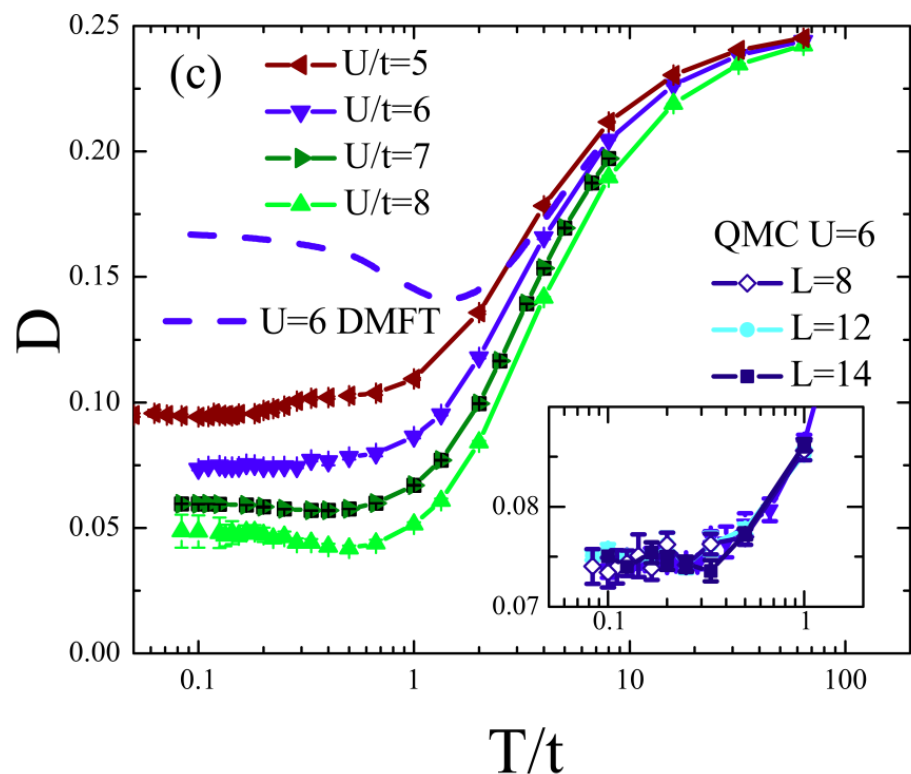
Thereza Paiva,¹ Richard Scalettar,² Mohit Randeria,³ and Nandini Trivedi³

¹*Instituto de Física, Universidade Federal do Rio de Janeiro Cx.P. 68.528, 21941-972 Rio de Janeiro RJ, Brazil*

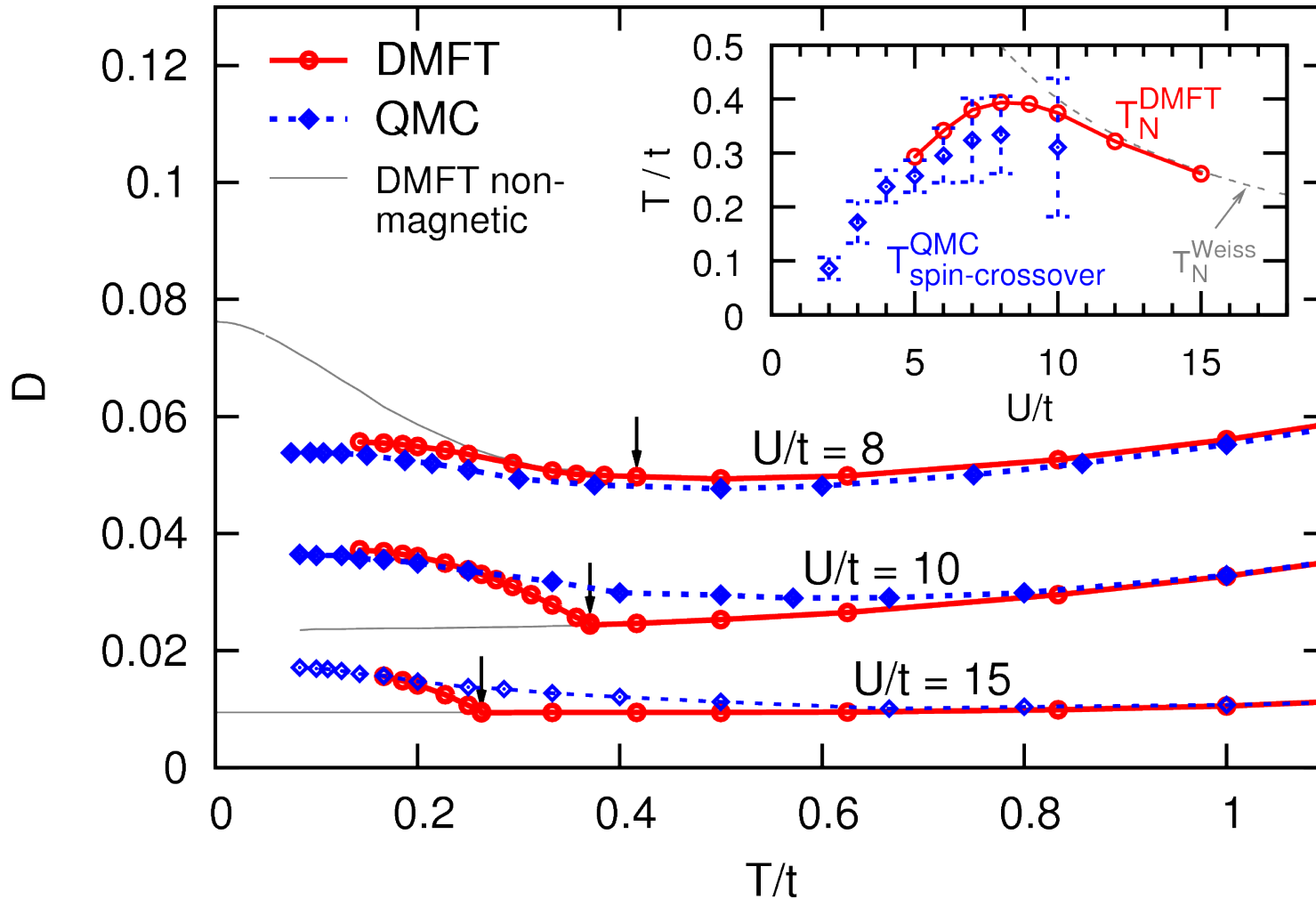
²*Department of Physics, University of California, Davis, California 95616, USA*

³*Department of Physics, The Ohio State University, Columbus, Ohio 43210, USA*

(Received 18 June 2009; published 11 February 2010)



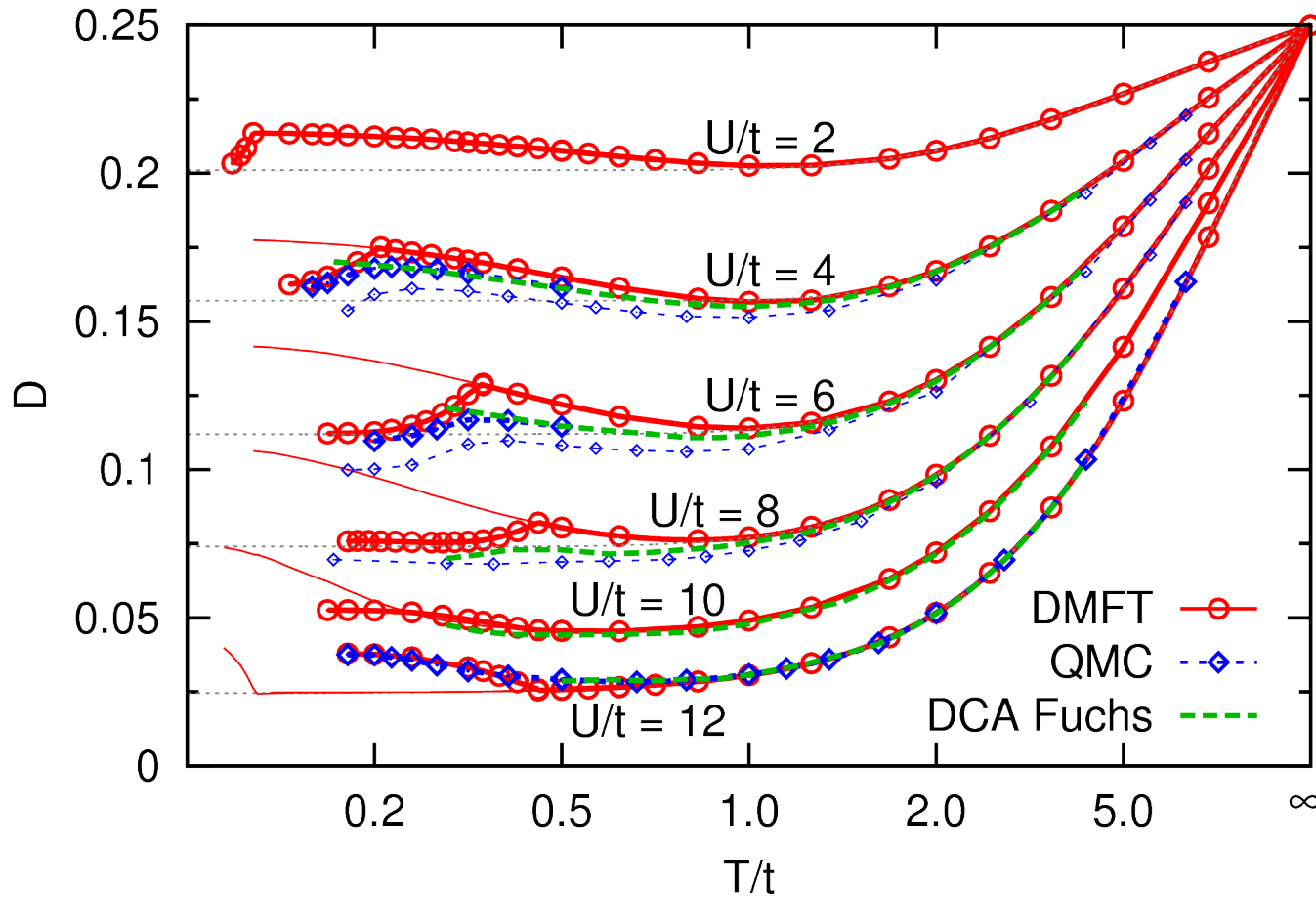
Comparison DMFT – direct QMC for the 2d square lattice ($n = 1$)



Excellent agreement at large T and low T , rounding off at $T \approx T_N^{\text{DMFT}}$

T_N^{DMFT} is relevant temperature scale for AF correlations!

Comparison DMFT – direct QMC for the 3d cubic lattice ($n = 1$)

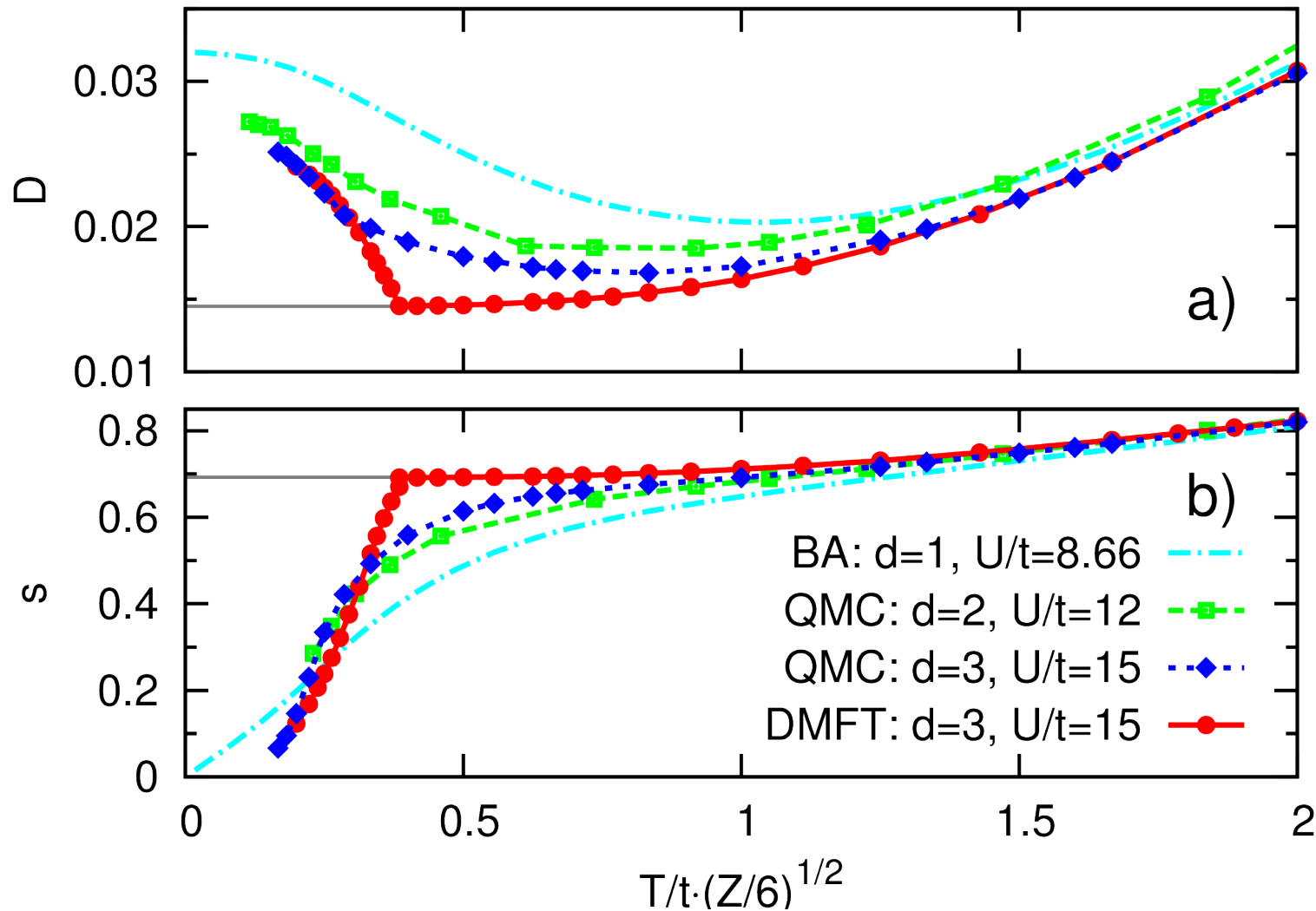


Excellent general agreement, even at small U (reduction of D by AF)

Dynamical cluster approximation (DCA) misses AF signatures in D

Typical QMC discretization errors (thin lines) larger than DMFT deviations

Dimensional comparison at $n = 1$: DMFT ($d = \infty$) vs. $d = 3, 2, 1$



Regular dimensional convergence (only for proper scaling $\propto \sqrt{Z}t$)

But: minimum in $D(T)$ shifts with d – offset by $s(T)$?

Thermodynamics of the three-dimensional Hubbard model: Implications for cooling cold atomic gases in optical lattices

Lorenzo De Leo,¹ Jean-Sébastien Bernier,¹ Corinna Kollath,^{1,2} Antoine Georges,^{1,3} and Vito W. Scarola⁴

¹*Centre de Physique Théorique, Ecole Polytechnique, CNRS, 91128 Palaiseau, France*

²*Département de Physique Théorique, Université de Genève, 24 quai Ernest Ansermet, CH-1211 Genève 4, Switzerland*

³*Collège de France, 11 place Marcelin Berthelot, 75005 Paris, France*

⁴*Department of Physics, Virginia Tech, Blacksburg, Virginia 24061, USA*

(Received 15 September 2010; published 10 February 2011)

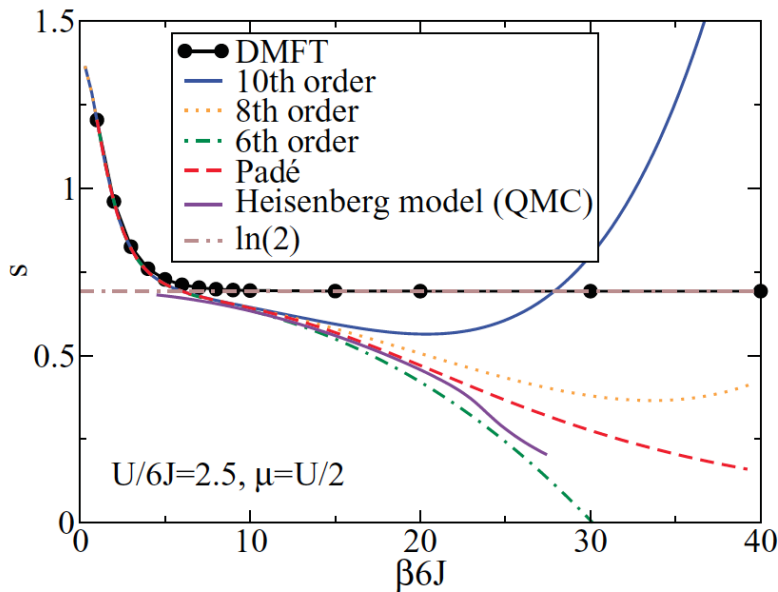
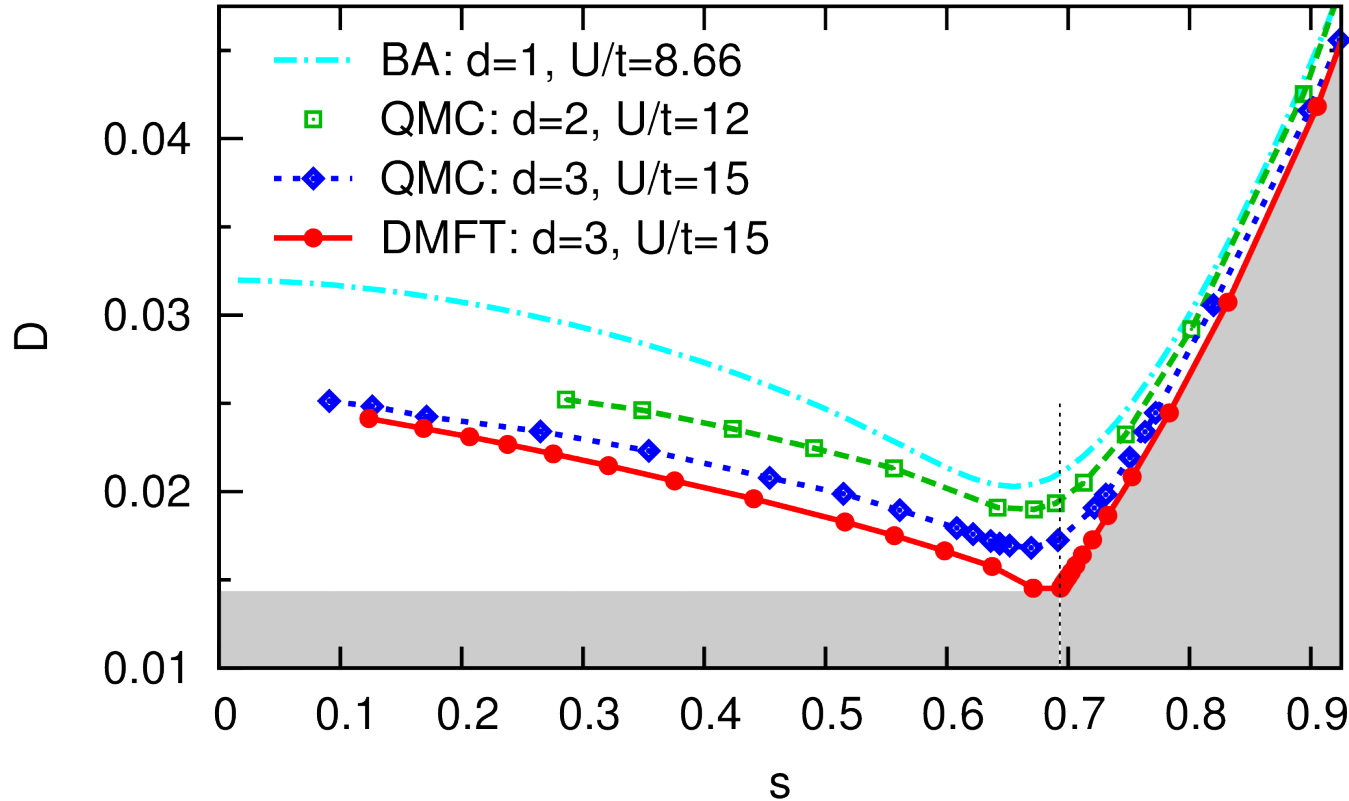


FIG. 14. (Color online) Entropy per particle in a system at half filling and intermediate interaction strength $U/6J = 2.5$ obtained by series expansion, DMFT, and QMC (for the Heisenberg model) [26].

Double occupancy as a universal measure of AF correlations + entropy



For all dimensions: $\min\{D(s)\}$ at $s \approx \log 2$

No significant features in $d = 3$ at Néel transition

Effect is larger for lower dimensions:

$$\langle \sigma_i \cdot \sigma_j \rangle_0 = \begin{cases} -1.00 & DMFT \\ -1.20 & (d = 3) \\ -1.34 & (d = 2) \\ -1.77 & (d = 1) \end{cases}$$

$$D_0 = (1 - \langle \sigma_i \cdot \sigma_j \rangle) Zt^2/(2U^2) + \mathcal{O}(t^4/U^4)$$

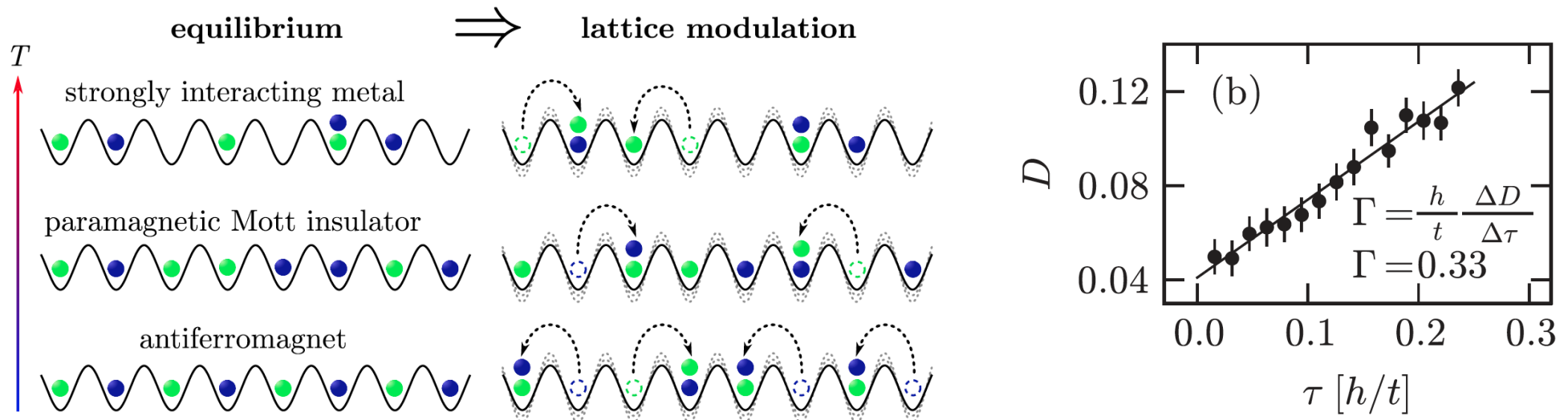
Probing Nearest-Neighbor Correlations of Ultracold Fermions in an Optical Lattice

Daniel Greif, Leticia Tarruell,* Thomas Uehlinger, Robert Jördens, and Tilman Esslinger

Institute for Quantum Electronics, ETH Zurich, 8093 Zurich, Switzerland

(Received 2 December 2010; revised manuscript received 1 February 2011; published 5 April 2011)

We demonstrate a probe for nearest-neighbor correlations of fermionic quantum gases in optical lattices. It gives access to spin and density configurations of adjacent sites and relies on **creating additional doubly occupied sites by perturbative lattice modulation**. The measured correlations for different lattice temperatures are in good agreement with an *ab initio* calculation without any fitting parameters. This probe opens new prospects for studying the approach to magnetically ordered phases.



$$\text{Correlator } \mathcal{P}_{i,i+1} = \sum_{\sigma} \langle n_{i,\sigma} (1 - n_{i,\bar{\sigma}}) n_{i+1,\bar{\sigma}} (1 - n_{i+1,\sigma}) \rangle \xrightarrow{n \rightarrow 1} (1 - \langle \sigma_i^z \sigma_{i+1}^z \rangle) / 2$$

MF?

Summary

QMC based implementation of real-space DMFT

Efficient for cold-atom temperatures, flexible

$\mathcal{O}(10^5)$ particles within slab approximation (\sim GGA)

Real-space DMFT study of antiferromagnetism

AF correlations at finite T signaled by enhanced D

Proximity effects important – LDA deficient

[Gorelik, Titvinidze, Hofstetter, Snoek, Blümer, PRL (2010)]

DMFT surprisingly accurate in low dimensions

Double occupancy: universal probe of AF correlations and entropy

[E. V. Gorelik, T. Paiva, R. Scalettar, A. Klümper, N. Blümer, arXiv:1105.3356]

D quantifies frustration effects (square – triangular – cubic lattice)

